

1 Sway Model for the Lateral Torsional Buckling Analysis of 2 Wooden Twin-Beam-Deck Systems

3 By Yang Du ^a, Magdi Mohareb ^b, and Ghasan Doudak ^c

4 ^a Department of Civil Engineering, University of Ottawa, Ottawa, ON, Canada, K1W6N5. Email:
5 ydu082@uottawa.ca

6 ^b Department of Civil Engineering, University of Ottawa, Ottawa, ON, Canada, K1W6N5 (corresponding
7 author). Email: mmohareb@uottawa.ca

8 ^c Department of Civil Engineering, University of Ottawa, Ottawa, ON, Canada, K1W6N5. Email:
9 gdoudak@uottawa.ca

10 Abstract

11 The present study develops a finite element model for the lateral torsional buckling
12 analysis of wooden twin beams braced by deck boards subjected to gravity or wind
13 uplift loading. The restraining action of the deck boards is modelled as continuous
14 partial lateral and twist restraints provided at the top of both beams that capture the
15 interaction between both beams. A parametric study is then conducted to examine the
16 effects of beam and deck spans, load type, load height, lateral restraint height and
17 stiffness and number of beam spans on the overall buckling capacity. The results
18 indicate that the restraining effects of deck boards significantly improve the lateral
19 torsional buckling capacity of twin-beam-deck assemblies.

20 Author Keywords

21 Lateral torsional buckling; wooden beam-deck system; finite element analysis; lateral
22 restraint; twist restraint

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26 **1. Introduction**

27 Timber floor and roof systems often consist of a set of relatively deep rectangular beams
28 braced at the top by tongue-and-groove wooden deck boards through nail connections.
29 While deep beams provide efficient means of material usage in bending, they are prone
30 to lateral torsional buckling (LTB) as a possible mode of failure. Deck boards can
31 potentially improve the buckling capacity of such systems by partially restraining the
32 twist and relative lateral displacements between the parallel beams. Despite the
33 common use of such systems in timber structures, contemporary design standards (NDS
34 [1], CAN/CSA O86 [2], Eurocode 5 [3]) recognize the beneficial role of deck boards in
35 suppressing LTB of beams when deck boards prevent the lateral displacement along the
36 beams' compression edge, but do not provide design guidelines for determining the
37 LTB capacity of wooden beams where (1) the tension edge is restrained as may be the
38 case for wind uplift on roof beams, or where (2) the deck boards provide partial lateral
39 and/or twist restraints for beams subjected to gravity loads. In this context, the present
40 study aims to develop a finite element solution for the LTB analysis of timber beam-
41 deck systems that incorporates (1) the characteristics of nailed joints in providing partial
42 lateral restraint, and (2) the beneficial effect of partial twist restraint provided by the
43 deck boards to the beams. Unlike previous LTB studies which treat each beam
44 separately while idealizing the effect of deck boards using lateral or rotational springs,
45 the present study treats both beams and the deck boards as an integrated system that
46 naturally captures the interaction between all three components throughout LTB, while
47 keeping the computational and modelling effort to a minimum. A pilot study on twin-
48 beam-deck systems (Du et al. [4]) based on ABAQUS showed that non-sway models
49 have a propensity buckle under net wind uplift but not under gravity loads, while sway
50 models can buckle under net wind uplift and gravity. Analytical and numeric models

51 for non-sway LTB were then formulated in (Du et al. [5]). The present study
52 complements past work by developing a finite element formulation for the LTB of twin-
53 beam-deck systems that captures the sway effects.

54 **2. Literature Review**

55 Consistent with the objective of the present study, a review is firstly presented for the
56 LTB literature for wood beams, followed by a review of stability research on
57 continuously restrained beams under wind and gravity loading. In general, continuously
58 restrained beams can be categorized into sway models, whereby the beams are free to
59 move laterally, and non-sway models where beams are fully restrained from lateral
60 movements at deck level. Since a recent review of non-sway models has been presented
61 in Du et al. [5], only research pertaining to sway models is presented herein. Finally, a
62 review of past research related to the strength and stiffness requirements for continuous
63 restraints is provided.

64 **2.1 Overview of stability research on wood members**

65 Past research on the LTB of wood beams has primarily focused on single beams and on
66 comparing test results against the classical LTB solution or relevant design standard
67 provisions. This includes the work of Hooley and Madsen [6] who developed design
68 equations for elastic and inelastic LTB of rectangular beams and verified those
69 equations experimentally. Hindman et al. [7] experimentally tested the LTB capacity of
70 unbraced rectangular cantilevers and compared test results with design standard
71 provisions. In a subsequent study, Hindman et al. [8] expanded their study to composite
72 I-joists. Burow et al. [9] and Burow et al. [10] evaluated the adequacy of LTB design
73 equations by testing simply-supported and cantilevered composite I-joists with a wide
74 range of slenderness ratios. Xiao [11] conducted a full-scale experimental investigation
75 on the elastic LTB of simply-supported beams with rectangular cross-section. A 3D

76 finite element model was also developed to study the sensitivity of material properties
77 and the effects of load height and boundary support height on the LTB capacity. St-
78 Amour [12] experimentally investigated the LTB of simply-supported composite I-
79 joists and developed an eigenvalue 3D finite element analysis (FEA) model for straight
80 beams, and a non-linear FEA model for beams with initial imperfections. Hu et al. [13]
81 and Hu et al. [14] developed lateral buckling solutions for wood beams with rigid and
82 flexible mid-span lateral braces offset from the shear center. However, none of the
83 above studies considered beams with continuous bracing. Such systems have been
84 studied by Zahn [15] who formulated the equilibrium equations governing the behavior
85 of rectangular beams partially restrained by the diaphragm action of deck boards. The
86 solution was experimentally verified by Jenkinson and Zahn [16] against the results of
87 a full-scale LTB test for a system consisting of two beams joined by deck boards and
88 subjected to uniform moments. An energy formulation was also developed by Zahn [17]
89 for the LTB of a single beam within a floor system by accounting for the internal strain
90 energy within the tributary decking strip of the beam. In summary, except for the study
91 of Jenkinson and Zahn [16], past experimental LTB studies on wood beams have
92 focused on individual beams as opposed to beam-deck systems.

93 **2.2 Sway models under gravity load**

94 Vlasov [18] formulated the general differential equations for a beam embedded in an
95 elastic medium and subjected to uniform moment loading. The critical moment was
96 determined for a beam continuously braced along its span through elastic lateral and
97 twist restraints. Pincus and Fisher [19] developed a LTB solution for twin beams braced
98 by a shear diaphragm and subjected to uniform moments. The solution captured the
99 shear and twist actions of the steel diaphragm. Errera et al. [20] extended the solution
100 of Pincus and Fisher [19] to other bracing scenarios. Floor assemblies consisting of two

101 beams laterally braced at their compression edge by the diaphragm shear action were
102 investigated by Apparao [21] who developed an eigen-value buckling solution for
103 straight beams and a nonlinear load-deformation solution for beams with initial
104 imperfections. Hancock and Trahair [22] formulated a finite element solution for
105 continuously braced beam-columns which characterized the deck action as elastic
106 restraints that partially restrain lateral displacement, twist along beam longitudinal axis,
107 weak-axis rotation and warping of the member. A closed-form solution was also
108 developed by Trahair [23] for beam-columns under constant loading. Assadi and
109 Roeder [24] investigated the stability of cantilevers with elastic lateral restraint at the
110 top edge. Albert and Dawe [25] developed a finite element solution for the buckling
111 analysis of a two-span I-section beam whose top flange is restrained by elastic lateral
112 and twist springs. For the purpose of capturing the interaction between the web and
113 flanges, the flanges were modelled by one-dimensional beam elements while the web
114 was modelled by plate elements. The inelastic behavior was captured by omitting
115 yielded zones in the beam. Khelil and Larue [26] developed an energy-based solution
116 for simply-supported beams with continuous partial lateral restraint. The solution was
117 able to predict the buckling capacity for beams with uniform and non-uniform moment
118 distributions.

119 **2.3 Sway models under wind uplift**

120 Studies on sway models with relatively low self-weight and subjected to net wind uplift
121 include the work of Pekoz and Soroushian [27] developed a buckling solution for steel
122 purlins restrained by sheeting through modelling the systems as beams on elastic
123 foundation. The problem was also investigated by Sokol [28] who idealized the purlin-
124 sheeting systems as columns on elastic rotational foundation. Lucas et al. [29]
125 formulated a non-linear elasto-plastic finite element model with geometric nonlinearity.

126 The model captured the interaction between purlins and sheeting, cross-sectional
127 distortion and local buckling effects. In a subsequent study, Lucas et al. [30] developed
128 a simplified model where the sheeting was idealized as elastic springs. Ye et al. [31]
129 adopted a finite strip model to investigate the local, distortional and lateral torsional
130 buckling of zed-purlins restrained by sheeting. The LTB of steel purlins laterally braced
131 by sheeting was investigated by Chu et al. [32], Li [33] and Chu et al. [34]. Basaglia et
132 al. [35] developed a solution for local, distortional and lateral torsional buckling of
133 channel and zed purlins restrained by steel sheeting. Apart from Sokol [28], Chu et al.
134 [32], Li [33], and Chu et al. [34], the above studies have focused on local or distortional
135 buckling of purlins. A common aspect among the above studies is that they were aimed
136 at steel purlin-sheeting systems. In contrast, the present study develops a LTB finite
137 element solution that tackles aspects specific to wooden beam-to-deck connections.

138 **2.4 Strength and stiffness requirements for continuous restraints**

139 Based on the studies of Errera et al. [20] and Apparao [21], Nethercot and Trahair [36]
140 as well as Errera and Apparao [37] proposed simplified design equations for the critical
141 loads of a single beam continuously braced by a shear diaphragm under uniform
142 moment loading. The stiffness requirement of the diaphragm to achieve desired moment
143 resistance was developed. The strength requirement was also proposed for initially
144 crooked beams. Starting from the models of Hancock and Trahair [22] and Trahair [23],
145 Trahair and Nethercot [38] added discrete restraint features to assess the effects of type
146 of continuous and discrete restraints on the buckling capacity of beam-columns.
147 Lawson and Nethercot [39] proposed a critical moment equation for a single beam
148 continuously braced by shear diaphragm. The solution incorporated moment gradient
149 and load height effects. Also developed were criteria to assess the adequacy of the
150 diaphragm shear stiffness. Helwig and Frank [40] developed a finite-element model for

151 the stability of twin-beam systems with sheeting acting as a shear diaphragm at the
152 beams' top edges. The authors confirmed the validity of the design equations proposed
153 by Nethercot and Trahair [36] for beams under uniform moments, and proposed
154 equations that account for loading heights. Also proposed was the diaphragm shear
155 stiffness required to achieve the design moment resistance. Helwig and Yura [41]
156 proposed a moment gradient equation for beams with stocky webs and suggested the
157 optimal diaphragm shear stiffness for beams with initial imperfections. In a later study,
158 Helwig and Yura [42] quantified the forces in shear diaphragms and outlined a
159 procedure for the design of shear diaphragms.

160 **2.5 Summary of literature**

161 Among the analytical solutions for the stability of continuously braced beams, the
162 studies of Vlasov [18], Albert and Dawe [25], Ye et al. [31], Li [33], and Basaglia et al.
163 [35] focused on the buckling of a single beam with continuous lateral and twist
164 restraints. In contrast, the present study tackles beam-deck systems. Analytical
165 solutions for the LTB of beam-deck systems consist of the work of Pincus and Fisher
166 [19], Errera et al. [20], Apparao [21], and Du et al. [5]. However, Pincus and Fisher
167 [19], Errera et al. [20] and Apparao [21] modelled the deck restraint through separate
168 rotational springs which do not fully capture the interaction between the two beams.
169 The study of Du et al. [5], while accounting for the interaction between both beams,
170 omitted the lateral sway effects. Another distinctive feature of the present study is that
171 it captures the partial relative lateral restraint between both beams provided by the wood
172 deck and the nailed connections, a feature absent in the work of Pincus and Fisher [19],
173 Errera et al. [20], Apparao [21] and Du et al. [5].

174 3. Constitutive Properties of Wood

175 Wood can be idealized as orthotropic with different mechanical properties in
176 longitudinal (L), radial (R) and tangential (T) directions. For a wood member, the
177 longitudinal axis is parallel to the grain, the radial axis is perpendicular to the grain and
178 normal to the growth rings, and the tangential axis is perpendicular to the grain but
179 tangent to the growth rings. Orthotropic constitutive behavior is characterized by twelve
180 properties: three moduli of elasticity in the longitudinal, radial, and tangential directions,
181 three shear moduli, and six Poisson's ratios. The moduli of elasticity and Poisson's
182 ratios are interrelated (e.g., FPL [43]), which reduces the number of independent
183 material properties to nine constants. Isopescu et al. [44] reported that the difference
184 between properties in radial and tangential directions is insignificant in most wood
185 species, thus reducing the number of independent properties to six. When no distinction
186 is made between properties in radial and tangential directions, the circular orthotropy
187 of wood coincides with rectangular orthotropy. Xiao et al. [45] conducted a LTB
188 sensitivity analysis for wood beams using a 3D finite element model within the
189 commercial software ABAQUS. Material was considered as orthotropic with six
190 independent constants. The study showed that only the modulus of elasticity E_L in the
191 longitudinal direction and the transverse shear modulus $G_T = G_{LR} = G_{LT}$ influence the
192 LTB capacity. Other material properties $E_T = E_R, G_{RT}, \mu_{LR} = \mu_{LT}, \mu_{RT}$ were shown to
193 have nearly no influence on the critical moments. Changes in these properties from the
194 default values by $\pm 50\%$ resulted in differences of less than 1% in the critical moment
195 predicted. The findings are also consistent with the observation from Hooley and
196 Madsen [6] and with the critical moment equation in Eurocode 5 [3] which depends
197 only on two constants E and G . The observation suggests some similarity with
198 isotropic beams where the elastic LTB capacity is characterized only by the modulus of

199 elasticity and the shear modulus. Nevertheless, isotropic beams differ from wood
200 orthotropic beams in that while the former has a G/E ratio ranging between 1/3 and
201 1/2, the later has a significantly lower G/E ratio in the neighborhood of 1/16.

202 **4. Assumptions**

203 The following assumptions have been adopted:

204
205 1. The formulation captures warping effects. While the study primarily targets beams
206 with rectangular sections for which warping has a small contribution, the model
207 remains equally valid for other doubly symmetric sections with more pronounced
208 warping effects.

209 2. Beam and deck materials are assumed linearly elastic. As discussed under Section
210 3, the orthotropic behavior of wood is characterized by six constants E_L , $E_T = E_R$,
211 $G_L = G_{LR} = G_{LT}$, G_{RT} , $\mu_L = \mu_{LR} = \mu_{LT}$ and μ_{RT} . Further, in the beam formulation
212 sought, since the normal stresses in the radial and tangential directions are
213 negligible compared to longitudinal stresses, the associated Young's moduli
214 $E_T = E_R$ will not contribute to the internal strain energy expression. Also, given
215 that the stress state in the beam is uniaxial, Poisson's ratios μ_{LR} , μ_{LT} , μ_{RT} will not
216 influence the internal strain energy. Finally, as the shear stresses acting on the radial
217 and tangential directions are negligible, the associated shear modulus G_{RT} does
218 not appear in the strain energy expression. In summary, of the six orthotropic
219 constitutive constants, only the modulus of elasticity in the longitudinal direction
220 E_L and shear modulus G_T will appear in the internal strain energy expression, in
221 a manner similar to isotropic materials.

222 3. Relative lateral displacements between beams and deck boards at the nailed joints
223 are allowed in order to account for the flexibility of connections. Deck boards are

224 assumed to provide partial relative lateral restraint due to the combined flexibility
225 of the nailed joints and deck boards (deck-joint assembly).

226 4. At each beam-deck joint, the deck board is assumed to rotate by the same angle of
227 twist for the beam (i.e., fully fixed connection is assumed). Consequently, deck
228 boards are assumed to provide continuous partial twist restraint to the beams.

229 5. The in-plane elastic shear restraint provided by the deck boards is negligible.

230 6. Throughout deformation, beam cross-sections remain rigid in their own plane, i.e.,
231 distortional effects are neglected.

232 7. Shear deformation effects within the beams are negligible, and

233 8. Pre-buckling deformation effects are omitted.

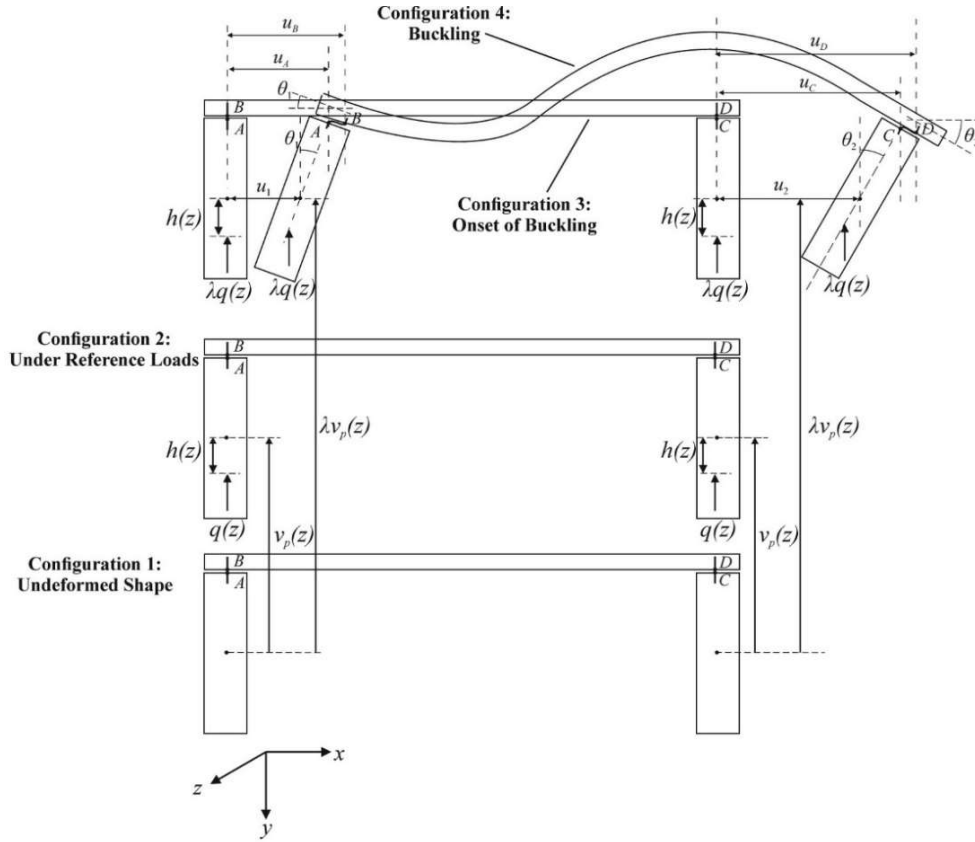
234 **5. Formulation**

235 **5.1 Problem description and notation**

236 The beam-deck system considered consists of two rectangular beams connected to deck
237 boards through nails. Fig. 1 shows a cross-sectional view of the twin-beam-deck
238 assembly. A left-hand coordinate system is used, with the z -axis oriented along the
239 longitudinal direction of the beam, and the x and y axes taken in the plane of the cross-
240 section. The positive sign convention adopted is consistent with that in Trahair [46]
241 where the angle of twist is clockwise. A transverse load $q(z)$ is assumed to be applied
242 to both beams at a distance $h(z)$ from the beam shear center. Under the external loads,
243 the system undergoes vertical displacement $v_p(z)$ in going from Configuration 1 to
244 2, in which subscript p denotes the pre-buckling displacements. The applied loads are
245 then assumed to be increased by a factor λ to attain the threshold value $\lambda q(z)$ at
246 the onset of buckling (Configuration 3). Under the load increase, pre-buckling

247 deformations are assumed to linearly increase to $\lambda v_p(z)$. The system then undergoes
248 LTB as manifested by lateral displacements $u_1(z), u_2(z)$ and angles of twist $\theta_1(z),$
249 $\theta_2(z)$ (Configuration 4) where subscripts 1,2 respectively denote the field variables
250 of the left and right beams. $u_1(z), u_2(z), \theta_1(z), \theta_2(z)$ are drawn in the positive
251 directions. In Fig. 1, points B and D are located at the deck underside and initially
252 coincide with points A and C located at the beams' top. At the onset of buckling, point
253 A at the top of beam 1 is assumed to undergo a lateral displacement u_A that differs
254 from that of point B at the deck underside. Also, point C located at the top of beam 2
255 undergoes a lateral displacement u_C which is different from that of point D , resulting
256 in a relative displacement between the top of both beams. The magnitude of relative
257 displacement depends on the lateral stiffness of the nailed joints, the axial and flexural
258 stiffness of the deck boards, including the effect of eccentricity between the deck board
259 centerline and beam-deck interface.

260



261

Fig. 1 Different stages of deformation

262 **5.2 Total potential energy**

263 The total potential energy Π of the twin-beam-deck system is

$$\Pi = U + V \quad (1)$$

264 in which U is the internal strain energy and V is the load potential energy. In Eq. (1),

265 the internal strain energy has four contributions $U = U_{b1} + U_{b2} + U_t + U_l$, where

266 U_{bi} ($i = 1, 2$) is the internal strain energy stored in beam i , U_t is the internal strain

267 energy for the partial twist restraint of the deck bending action, and U_l is the internal

268 strain energy for the relative partial lateral restraint of the deck-joint assembly. The load

269 potential energy V consists of two components, i.e., $V = V_{b1} + V_{b2}$, one for each beam.

270 The internal strain energy stored in beam i due to weak-axis bending, Saint-Venant
 271 torsion, and warping torsion takes the following form (Trahair [46])

$$U_{bi} = (1/2) \int_0^{L_b} \left(E_b I_y u_i''^2 + G_b J_b \theta_i'^2 + E_b C_w \theta_i''^2 \right) dz \quad (i=1,2) \quad (2)$$

272 where $E_b = E_{Lb}$ is the beam modulus of elasticity in the longitudinal direction, I_y is
 273 the moment of inertia about beam weak-axis, $G_b = G_{Lb}$ is the shear modulus for
 274 stresses acting on the plane normal to the longitudinal direction, J_b is the Saint-
 275 Venant torsional constant, C_w is the warping constant, L_b is the beam span, and all
 276 primes denote differentiation with respect to coordinate z taken along the beam
 277 longitudinal axis. For wood beams with rectangular cross-sections, material properties
 278 E_b and G_b are taken as the material properties of the wood type used. For doubly
 279 symmetric I-section members where the web and flanges have different materials, E_b
 280 and G_b are the transformed section properties (Du et al. [5]). The load potential energy
 281 for external loads applied at beam i (Trahair [46]) is

$$V_{bi} = \lambda \int_0^{L_b} \left(2M \theta_i u_i'' + qh \theta_i^2 \right) dz \quad (i=1,2) \quad (3)$$

282 where M is the bending moment induced by the reference load q (taken positive
 283 when acting downwards), and one recalls that λ is an unknown load multiplier to be
 284 obtained from the buckling analysis and h is the height of the loading point above the
 285 shear center (taken positive when the loading point is below beam shear center). The
 286 first term in Eq. (3) accounts for the destabilizing effect due to strong-axis flexure
 287 while the second term accounts for load potential energy gain due to the load offset
 288 from the beam shear center. For a given deck board located at a distance z_0 from the

289 beam end support, and angles of twist $\theta_1(z_0)$ and $\theta_2(z_0)$ the internal strain energy
 290 stored in this deck board is

$$U_i^* = \frac{E_d I_d}{L_d} \langle \theta_1(z_0) \quad \theta_2(z_0) \rangle \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{Bmatrix} \theta_1(z_0) \\ \theta_2(z_0) \end{Bmatrix} \quad (4)$$

291 where E_d is the deck modulus of elasticity, I_d is the moment of inertia about the
 292 deck board and L_d is the span of deck board. The internal strain energy U_i for the
 293 whole deck is simply the summation of the contribution from each deck board which
 294 can be approximately written in an integration form as

$$U_i = \sum U_i^* \approx \frac{E_d h_d^3}{12 L_d} \int_0^{L_b} \langle \theta_1(z) \quad \theta_2(z) \rangle \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{Bmatrix} \theta_1(z) \\ \theta_2(z) \end{Bmatrix} dz \quad (5)$$

295 where h_d is the deck board thickness. For a deck board at a distance z_0 from beam
 296 end support, the lateral displacements of points A and C at beams' top (Fig. 1) are
 297 expressed in terms of the lateral displacement u_i and angle of twist θ_i as

$$u_A(z_0) = u_1(z_0) + (h_b/2)\theta_1(z_0), \quad u_C(z_0) = u_2(z_0) + (h_b/2)\theta_2(z_0) \quad (6)$$

298 where h_b is the beam depth.

299 The internal strain energy U_i associated with the relative partial lateral restraint
 300 provided by the deck to the beams can be characterized by considering the lateral
 301 stiffness k of the deck-joint assembly as three springs in tandem (i.e., two springs,
 302 each representing the shear stiffness k_j of each joint at the beam-deck interface, and
 303 another spring representing the axial stiffness of the deck k_d). Stiffness k relates the
 304 relative displacement $u_c - u_A$ between the two beams proportionally to the lateral
 305 force $F = k(u_c - u_A)$. Adding the flexibility of the three spring components yields
 306 $1/k = 1/k_j + 1/k_d + 1/k_j$, which is solved for the overall stiffness k as

$$k = k_d k_j / (2k_d + k_j) \quad (7)$$

307 Assuming \bar{n} nails are installed at each beam-deck joint, each with the shear stiffness
 308 k_n , the shear stiffness k_j of the joint is $k_j = \bar{n}k_n$. The stiffness of a single-nail joint
 309 k_n can be estimated based on the load-displacement model in CAN/CSA O86 [2]

$$\Delta_n = 0.5d_f K_m (F/F_u)^{1.7} \quad (8)$$

310 where Δ_n is the shear deformation of a single-nail joint, d_f is the nail diameter in
 311 millimeter, K_m is the service creep factor accounting for load duration and moisture
 312 level, and F_u is the lateral strength resistance as governed by the smallest capacities
 313 amongst different failure modes. In the present study, the value $K_m = 1$ is used. The
 314 shear stiffness k_n is taken as the initial stiffness from Eq. (8), i.e., gradient that joins
 315 the origin and 10% of the lateral strength resistance F_u .

316 The model that quantifies the deck axial stiffness at the height of beam-deck interface
 317 is schematically depicted in Fig. 2 where the deck board is assumed to be subjected to
 318 a pair of eccentric loads F acting at the beam-deck interface. The lateral deformation
 319 Δ of the deck board is consisted of two components, an axial deformation Δ_1 , and
 320 another deformation Δ_2 associated with transverse bending, i.e.,

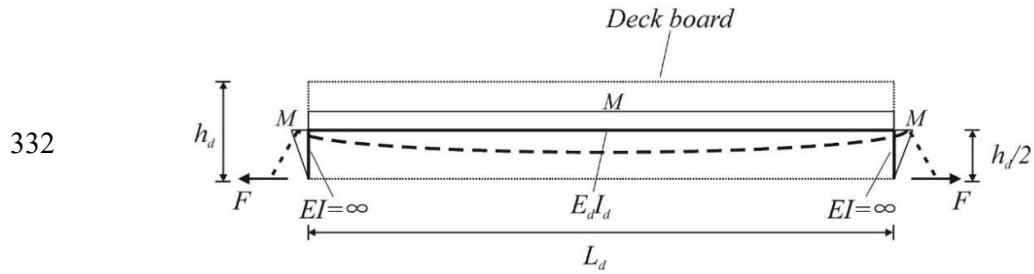
$$\Delta = \Delta_1 + \Delta_2 = FL_d/E_d b h_d + M_d M_0 L_d / E_d I_d \quad (9)$$

321 where M_d is the bending moment induced by eccentric lateral loads F and is
 322 expressed as $M_d = Fh_d/2$, b is deck board width, and M_0 is the moment for unit
 323 lateral loads. From the expressions for M_d and M_0 , by substituting into Eq. (9), one
 324 obtains $\Delta = FL_d/E_d b h_d + Fh_d^2 L_d / 4E_d I_d = 4FL_d/E_d b h_d$ and the axial stiffness of the
 325 deck board k_d at the height of beam-deck interface is $k_d = E_d b h_d / 4L_d$. From the

326 expression for k_d , by substituting into Eq. (7), one obtains the sought relative lateral
 327 stiffness as

$$k = \bar{n}E_d b h_d k_n / (2E_d b h_d + 4\bar{n}L_d k_n) \quad (10)$$

328 For example, for a 38 mm by 140 mm deck board spanning 2 m and having a modulus
 329 of elasticity of 10,000 MPa, the relative lateral stiffness of a single nail with a 3.66mm
 330 diameter at each joint is 1061 kN/m or 7579 kN/m/m (per unit length of deck), while
 331 that based on two nails at each joint is 1830 kN/m, or 13073 kN/m/m.



333 **Fig. 2** Deck board deformation due to eccentric lateral loads acting at the deck underside
 334 The corresponding internal strain energy stored in the deck board is

$$U_i^* = (k/2)(u_c - u_A)^2 = (k/2)\{u_2(z_0) - u_1(z_0) + (h_b/2)[\theta_2(z_0) - \theta_1(z_0)]\}^2 \quad (11)$$

335 For the whole deck, the internal strain energy is the summation of the energy
 336 contributions from all deck boards, which can be approximated in an integral form as

$$U_i = \sum U_i^* \approx (\bar{k}/2) \int_0^{L_b} [u_2 - u_1 + (h_b/2)(\theta_2 - \theta_1)]^2 dz \quad (12)$$

337 where the relative lateral stiffness per unit deck width $\bar{k} = k/b$ has been defined. In
 338 summary, from Eqs. (2),(3),(5), and (12), by substituting into Eq. (1), one obtains the
 339 total potential energy of the twin-beam-deck assembly as

$$\begin{aligned}
\Pi = & \frac{1}{2} \int_0^{L_b} E_b I_y (u_1''^2 + u_2''^2) + G_b J_b (\theta_1'^2 + \theta_2'^2) + E_b C_w (\theta_1''^2 + \theta_2''^2) \\
& + (E_d h_d^3 / 3L_d) (\theta_1^2 + \theta_1 \theta_2 + \theta_2^2) + \bar{k} [u_2 - u_1 + (h_b/2)(\theta_2 - \theta_1)]^2 \\
& + 2\lambda M (\theta_1 u_1'' + \theta_2 u_2'') + \lambda q h (\theta_1^2 + \theta_2^2) dz
\end{aligned} \quad (13)$$

340 6. Finite Element Formulation

341 The Hermitian polynomials are adopted to relate the lateral displacement fields u_i and
342 rotation fields θ_i ($i=1,2$) to the generalized nodal displacements, i.e.

$$u_i(z) = \langle L(z) \rangle_{1 \times 4}^T \{u_i\}_{4 \times 1}, \quad \theta_i(z) = \langle L(z) \rangle_{1 \times 4}^T \{\theta_i\}_{4 \times 1} \quad (14)$$

343 where

$$344 \langle L(z) \rangle_{1 \times 4}^T = \left\langle \left(1 - 3z^2/l^2 + 2z^3/l^3\right) \quad \left(z - 2z^2/l + z^3/l^2\right) \quad \left(3z^2/l^2 - 2z^3/l^3\right) \quad \left(z^3/l^2 - z^2/l\right) \right\rangle$$

345 is the vector of Hermitian interpolation functions and l is the element length,

346 $\langle u_i \rangle_{1 \times 4}^T = \langle u_0 \quad u_0' \quad u_l \quad u_l' \rangle_i$ is the generalized nodal lateral displacement vector,

347 $\langle \theta_i \rangle_{1 \times 4}^T = \langle \theta_0 \quad \theta_0' \quad \theta_l \quad \theta_l' \rangle_i$ is the generalized nodal angle of twist vector, and

348 subscripts 0 and l denote the nodal points of the element. From Eq. (14), by substituting

349 into Eq. (13), one obtains

$$\Pi = (1/2) \langle U \rangle_{1 \times 16}^T \left[[K_e] - \lambda [K_g] \right]_{16 \times 16} \{U\}_{16 \times 1} \quad (15)$$

350 where $\langle U \rangle^T = \langle \langle u_1 \rangle_{1 \times 4}^T \quad \langle u_2 \rangle_{1 \times 4}^T \quad \langle \theta_1 \rangle_{1 \times 4}^T \quad \langle \theta_2 \rangle_{1 \times 4}^T \rangle$ and the elastic stiffness matrix

351 $[K_e] = [K_b] + [K_t] + [K_l]$ has three contributions: $[K_b]$ which characterizes the beam

352 stiffness matrix, $[K_t]$ which characterizes partial twist restraint, and $[K_l]$ which

353 characterize the relative partial lateral restraint and are given by

$$[K_b] = \begin{bmatrix} E_b I_y [B_1] & [0] & [0] & [0] \\ [0] & E_b I_y [B_1] & [0] & [0] \\ [0] & [0] & E_b C_w [B_1] + G_b J_b [B_2] & [0] \\ [0] & [0] & [0] & E_b C_w [B_1] + G_b J_b [B_2] \end{bmatrix}_{16 \times 16} \quad (16)$$

$$[K_t] = \frac{E_d h_d^3}{6L_d} \begin{bmatrix} [0] & [0] & [0] & [0] \\ [0] & [0] & [0] & [0] \\ [0] & [0] & 2[B_3] & [B_3] \\ [0] & [0] & [B_3]^T & 2[B_3] \end{bmatrix}_{16 \times 16} \quad (17)$$

$$[K_l] = \frac{\bar{k}}{4} \begin{bmatrix} 4[B_3] & -4[B_3] & 2h_b [B_3]^T & -2h_b [B_3]^T \\ -4[B_3]^T & 4[B_3] & -2h_b [B_3]^T & 2h_b [B_3]^T \\ 2h_b [B_3] & -2h_b [B_3] & h_b^2 [B_3] & -h_b^2 [B_3] \\ -2h_b [B_3] & 2h_b [B_3] & -h_b^2 [B_3]^T & h_b^2 [B_3] \end{bmatrix}_{16 \times 16} \quad (18)$$

354 The geometric stiffness matrix $[K_g]$ takes the following form

$$[K_g] = \begin{bmatrix} [0] & [0] & -[B_4]^T & [0] \\ [0] & [0] & [0] & -[B_4]^T \\ -[B_4] & [0] & -h[B_5] & [0] \\ [0] & -[B_4] & [0] & -h[B_5] \end{bmatrix}_{16 \times 16} \quad (19)$$

355 and submatrices $[B_1], [B_2], [B_3], [B_4], [B_5]$ are defined as

$$356 \quad [B_1] = \int_0^l \{L''(z)\}_{4 \times 1} \langle L''(z) \rangle_{1 \times 4}^T dz, \quad [B_2] = \int_0^l \{L'(z)\}_{4 \times 1} \langle L'(z) \rangle_{1 \times 4}^T dz,$$

$$357 \quad [B_3] = \int_0^l \{L(z)\}_{4 \times 1} \langle L(z) \rangle_{1 \times 4}^T dz, \quad [B_4] = \int_0^l M(z) \{L(z)\}_{4 \times 1} \langle L''(z) \rangle_{1 \times 4}^T dz \quad \text{and}$$

$$358 \quad [B_5] = \int_0^l q(z) \{L(z)\}_{4 \times 1} \langle L(z) \rangle_{1 \times 4}^T dz. \quad \text{From Eq. (15), by evoking the stationarity of the}$$

359 total potential energy, i.e., $\delta \Pi = 0$, one obtains

$$\left([K_e] - \lambda [K_g] \right)_{16 \times 16} \{U\}_{16 \times 1} = \{0\}_{16 \times 1} \quad (20)$$

360 **7. Results**

361 A wooden twin-beam-deck assembly of common wood properties and practical
362 dimensions is chosen as a reference case to be used to assess the validity of the present
363 finite element solution through comparisons to the predictions of a model based on the
364 commercial software ABAQUS. The present finite element is then used to conduct a
365 parametric study by varying one parameter at a time from the reference case and
366 observing the critical moments and buckling mode shapes. For verification, the
367 transverse loads are applied at the beam shear center while for the parametric study,
368 unless specified otherwise, the transverse loads are applied at deck centerline.

369 The reference beam span is taken as 6m and the deck span is taken as 2 m. The beams
370 are assumed to be glue-laminated Spruce-Lodgepole Pine-Jack Pine 20f-EX with 570
371 mm in depth and 80 mm in width. The modulus of elasticity and shear modulus are
372 assumed to be 10,300 MPa and 474 MPa, respectively. The nominal bending moment
373 resistance based on material failure is 270 kNm. The deck boards are assumed to be
374 Spruce-Pine-Fir No. 2 grade and are 38 mm thick and 140 mm wide with a modulus of
375 elasticity of 10,000 MPa. All the material properties mentioned above are based on
376 CAN/CSA O86 [2] and FPL [43]. Two 3.66 mm nails are assumed to be installed at
377 each beam-deck intersection. Both uplift and downward buckling loads are investigated.

378 **7.1 Details of ABAQUS model**

379 The finite-element program ABAQUS was used to conduct an eigenvalue buckling
380 analysis of the twin-beam-deck assembly. The two-node B31OS element was used to
381 model the twin beams. Each node has seven degrees of freedom (i.e., three translations,
382 three rotations, and one warping deformation). The element accounts for shear
383 deformation only due to bending but ignores shear deformation induced by warping.
384 Two types of elements were chosen to capture the twist restraint of the deck and the

385 relative lateral restraint of the deck-joint assembly. The B31 element was used to model
 386 the flexural stiffness of the deck boards. The B31 element has two nodes with six
 387 degrees of freedom per node (i.e., three translations and three rotations). The two-node
 388 SPRING2 spring element was employed to simulate the partial relative lateral
 389 restraining action between the top of both beams for the deck-joint assembly. Fig. 3
 390 shows the twin-beam-deck model developed, with B31OS elements depicted at the
 391 centroidal axes of twin beams, B31 elements at the deck board centerlines, and
 392 SPRING2 elements spanning between the top of twin beams. The number of B31OS
 393 elements in each beam was chosen to be consistent with the number of deck boards.
 394 Each deck node (lying on line 2 or 5 in Fig. 3) was rigidly tied to the corresponding
 395 centroidal node (on line 1 or 4). Also, each spring node (on line 3 or 6) was tied to the
 396 corresponding centroidal node (on line 1 or 4). Both rigid links were defined using the
 397 BEAM type multi-point constraint (*MPC) feature in ABAQUS.

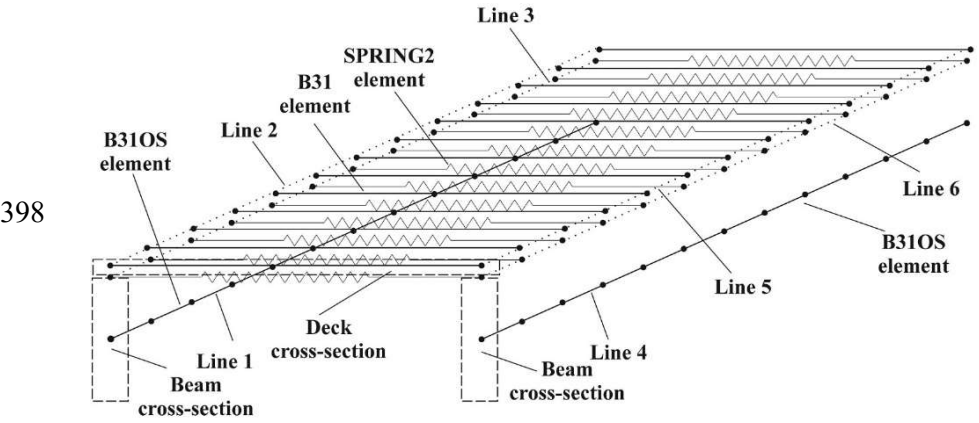


Fig. 3 ABAQUS twin-beam-deck model

400 **7.2 Mesh convergence study**

401 For the present FEA model, a mesh convergence study is performed for uniform
 402 moment loading (Fig. 4a), UDL (Fig. 4c), and mid-span concentrated loads (Fig. 4e).
 403 Systems with transverse loads applied at the beam shear center are observed to have
 404 equal critical positive moments (tension at the beams' bottom) and negative moments.

405 Thus, only the magnitudes of the moments are presented. Among all three load types
406 investigated, the present FEA results are observed to converge from above. No more
407 than 10 beam elements are found to be needed to attain convergence.

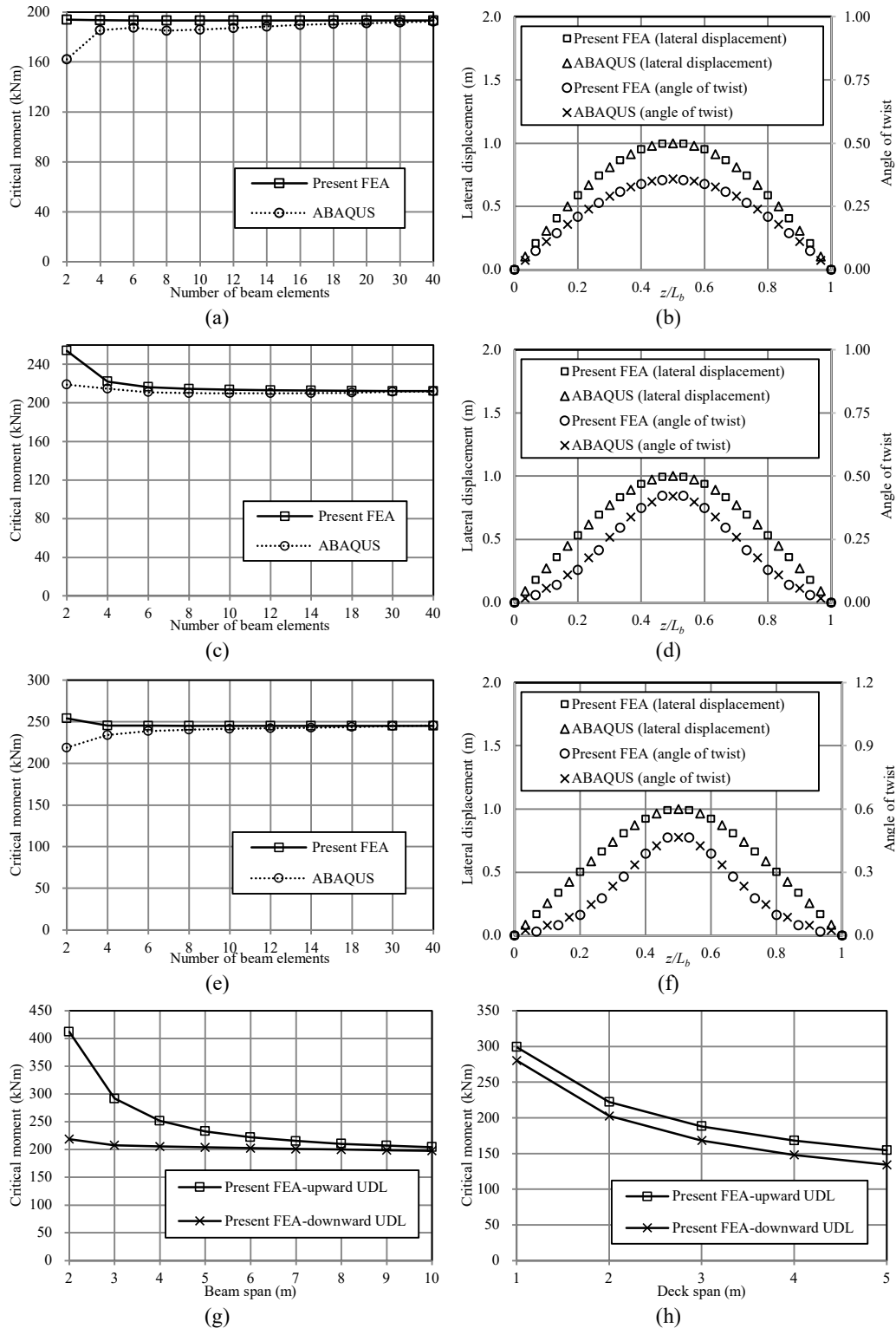
408 **7.3 Verification**

409 For uniform moment loading (Fig. 4a), the critical moments determined by the present
410 FEA solutions yield a value of 193 kNm, compared with a slightly lower value of 192
411 kNm predicted from ABAQUS. This is attributed to the fact that the ABAQUS B31OS
412 elements incorporate shear deformation effects due to flexure and thus provide a more
413 flexible representation of the system. The mode shapes as predicted by both solutions
414 are found in close agreement (Fig. 4b).

415 For UDL (Fig. 4c), the present FEA predicts a critical moment of 212.1 kNm which is
416 slightly higher than 211.8 kNm as predicted by the ABAQUS and the predicted mode
417 shapes nearly coincide (Fig. 4d). Similar observations are found for mid-span
418 concentrated loading (Fig. 4e, Fig. 4f). The above comparisons show good agreement
419 between the present FEA model and ABAQUS. The present FEA model, thus verified,
420 is subsequently used to investigate the effects of various parameters on critical moments.

421 **7.4 Effects of beam and deck spans**

422 The effects of beam and deck spans are illustrated in Fig. 4g and Fig. 4h. When the
423 beam span is increased from 2 m to 10 m (Fig. 4g), the buckling capacity for upward
424 UDL as predicted by the present FEA decreases from 412 kNm to 204 kNm. In contrast,
425 for downward UDL, the critical moment declines slightly from 219 kNm to 198 kNm.
426 For both upward and downward UDL, as the deck span is increased from 1 m to 5 m,
427 Fig. 4h shows a consistent decline in buckling capacity. The critical moment
428 corresponding to the case of a 5 m deck span is roughly half of the capacity of the case
429 of a 1 m deck span.



430 **Fig. 4** (a) Convergence study for uniform moments, (b) mode shapes for uniform moments, (c)
 431 convergence study for UDL, (d) mode shapes for UDL, (e) convergence study for mid-span
 432 concentrated loads, (f) mode shapes for mid-span concentrated loads, (g) critical moments for
 433 UDL with varying beam span, (h) critical moments for UDL with varying deck span

434 **7.5 Effect of relative lateral stiffness of the deck-joint assembly**

435 Fig. 5a shows the critical moments as a function of the relative lateral stiffness per unit
436 deck width \bar{k} for a beam-deck system subjected to UDL acting at the height of beam
437 shear center. Two beam spans were investigated, i.e., 2 m and 6 m. For the 2 m span,
438 the critical moments for downward loading increase with \bar{k} up to 100 kN/m/m, after
439 which the critical moment remains constant at 301 kNm. In contrast, the critical
440 moment for upward loading remains constant at 250 kNm irrespective of the magnitude
441 of \bar{k} . For the 6 m span, the critical moments for upward and downward UDL increase
442 as \bar{k} is increased. The increase in critical moments for downward loading is faster
443 than that of upward loading, resulting in a higher buckling capacity for a certain
444 stiffness range. The critical moments for upward and downward loading are found equal
445 with a magnitude of 212 kNm when \bar{k} exceeds the threshold value of 45.3 kN/m/m.
446 For the reference deck-joint assembly, \bar{k} corresponding to a single nail per joint is
447 7579 kN/m/m while that based on two nails per joint is 13073 kN/m/m. Both values are
448 depicted by the vertical lines on Fig. 5a. For the problem examined, both stiffness
449 values far exceed the threshold stiffness value, resulting in identical critical moment for
450 upward and downward loading. Further investigation on beam spans of 4 m and longer
451 suggests similar conclusions. Such results were omitted from Fig. 5a, to avoid clutter.
452 The present example suggests that a beam-deck connection with a single nail provides
453 enough lateral stiffness to develop their peak critical moments.

454 **7.6 Effect of lateral restraint height**

455 For the reference twin-beam-deck assembly subjected to UDL, the relative lateral
456 restraint of the deck-joint assembly is hypothetically moved from the beam shear center
457 to the deck centerline and the corresponding critical moments are presented in Fig. 5b.

458 Two categories of restraint stiffness were considered: below and above the threshold
459 value. Considered was the relative lateral stiffness per unit deck width $\bar{k} = 10$ kN/m/m,
460 a value below the threshold stiffness (discussed in the previous section). For upward
461 loading, the critical moments are observed to decrease from 205 kNm to 185 kNm as
462 the restraint moves from the shear center to the deck centerline. In contrast, for
463 downward loading, the critical moments are observed to increase from 168 kNm to 189
464 kNm. The results suggest that the restraint height has a moderate effect on the critical
465 moments for \bar{k} below the threshold value. For a relative lateral stiffness above the
466 threshold stiffness ($\bar{k} = 13,073$ kN/m/m for a two-nail connection), the critical moments
467 corresponding to upward and downward loading remain constant at 222 kNm and 203
468 kNm, respectively. Insight on the above observations can be provided by examining the
469 associated mode shapes. For a relative lateral stiffness below the threshold value, the
470 modes are observed to be symmetric (i.e., $u_1 = -u_2$ and $\theta_1 = -\theta_2$). Thus, the internal
471 strain energy for the partial lateral restraint (i.e., Eq.(12)) is non-zero and is influenced
472 by the restraint height. In contrast, when the relative lateral stiffness is beyond the
473 threshold value, the modes become anti-symmetric (i.e., $u_1 = u_2$ and $\theta_1 = \theta_2$), causing
474 Eq. (12) to vanish, thus eliminating the dependence of the solution on the restraint
475 height.

476 **7.7 Effects of partial relative lateral and twist restraints**

477 The critical moments for the reference twin-beam-deck assembly subjected to UDL
478 applied at the deck centerline are examined for various restraint scenarios: (1) partial
479 relative lateral and twist restraints included, (2) partial twist restraint with no relative
480 partial lateral restraint, (3) relative partial lateral restraint with no partial twist restraint,
481 and (4) no restraints. For downward loading (Fig. 5c), the contribution of the relative

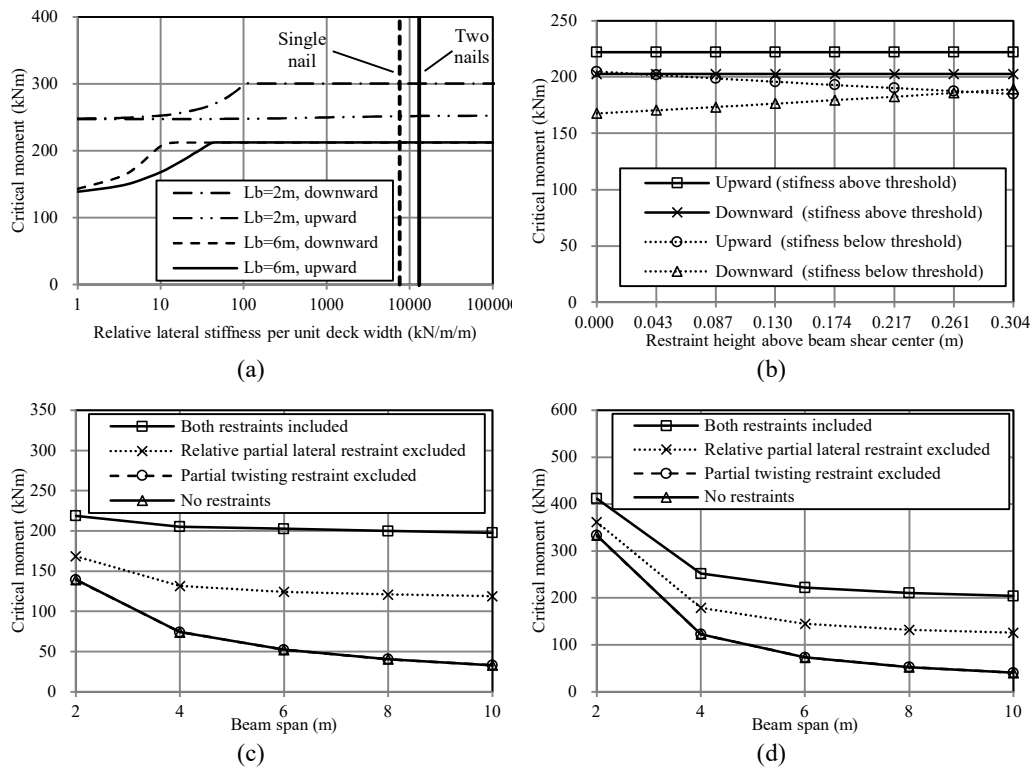
482 partial lateral restraint is illustrated by the gap between the upper solid line and the
483 dotted line. The contribution of the restraint grows as the span increases from 2 m to 6
484 m, after which the restraint contributes a steady 40% of the total capacity. Interestingly,
485 the results corresponding to twin-beams with relative partial lateral restraint alone (no
486 twist restraint) are observed to coincide with those of no restraints. This suggests that
487 the presence of relative lateral restraint alone between identical beams identically
488 loaded provides no LTB capacity increase. The results also suggest that the buckling
489 capacities with lateral and twist restraints are 1.6 times (beam with 2 m span) to 6 times
490 (beam with 10 m span) higher than those without restraints. Similar conclusions can be
491 drawn from the case of upward UDL loading (Fig. 5d).

492 **7.8 Effect of load position**

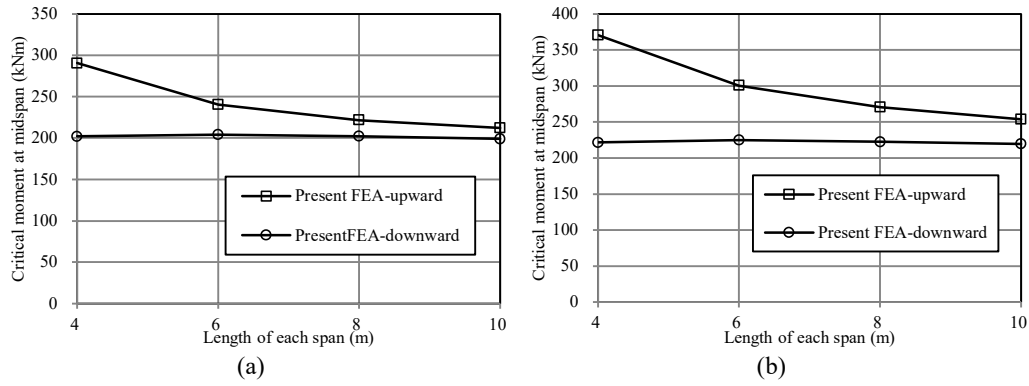
493 Three load positions are considered: (1) Deck centerline, (2) beam shear center and (3)
494 beam bottom. The critical moments are examined for beam spans of 4, 6, and 8 m and
495 deck spans of 1, 3, and 5 m and results are presented in Tables 1 and 2, respectively.
496 For downward loading, because of the destabilizing effect induced by deck centerline
497 loading, the buckling capacities are observed to be lower than those based on beam
498 shear center loading. The percentage decreases are 10.1%, 4.25% and 2.44% for 4, 6
499 and 8 m span, respectively. In contrast, loads applied at beam bottom is found to be
500 associated with a stabilizing effect which increases the buckling capacity. The
501 percentage increases are 9.65%, 4.72% and 2.44% for beam spans of 4, 6 and 8 m,
502 respectively. For upward loading, an opposite trend is observed where deck centerline
503 loading is found to correspond to the highest critical moments. The results in Table 1
504 suggest that the load position effect is more pronounced in short span beams. Table 2
505 suggests that the load position effect remains relatively steady (between 3%-7%) as the
506 deck span is varied from 1 m to 5 m.

507 **7.9 Effect of beam span on a two-span continuous twin-beam-deck system**

508 Fig. 6a and Fig. 6b show the buckling capacity of a twin-beam-deck system with two-
 509 span continuous beams under UDL and mid-span concentrated loads applied to both
 510 spans. For downward loading, the critical moments at mid-span remain nearly constant
 511 at 200 kNm and are found to be insensitive to the span. In contrast, for upward loading,
 512 critical moments are found to decrease as the span is increased from 4 to 10m. This
 513 trend is consistent with the observations from single-span beams. Table 3 provides a
 514 comparison of the critical moments for single-span and two-span twin-beam-deck
 515 systems with a UDL applied at the deck centerline. The results suggest that, in general,
 516 the critical moments for two-span twin-beam-deck assemblies are slightly higher than
 517 those based on single-span assemblies of the same span.



518 **Fig. 5** (a) Buckling capacity for (a) different lateral stiffness, (b) different lateral restraint
 519 height; Buckling capacity as influenced by inclusion/exclusion of partial lateral and twist
 520 restraints under (c) downward UDL with varying beam span, (d) upward UDL with varying
 521 beam span



522 **Fig. 6** Buckling capacity for two-span twin-beam-deck systems with varying spans under (a)
 523 UDL, (b) mid-span concentrated loads

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530 Table 1 Load position effect for different beam spans

Beam span (m)	Load position	Downward UDL		Upward UDL	
		Critical moment (kNm)	Difference	Critical moment (kNm)	Difference
4	Deck centerline	205	-10.1%	252	10.5%
	Beam shear center	228	/	228	/
	Beam bottom	250	9.65%	176	-22.8%
6	Deck centerline	203	-4.25%	222	4.72%
	Beam shear center	212	/	212	/
	Beam bottom	222	4.72%	203	-4.25%
8	Deck centerline	200	-2.44%	210	2.44%
	Beam shear center	205	/	205	/
	Beam bottom	210	2.44%	200	-2.44%

531

532 Table 2 Load position effect for different deck spans

Deck span (m)	Load position	Downward UDL		Upward UDL	
		Critical moment (kNm)	Difference	Critical moment (kNm)	Difference
1	Deck centerline	280	-3.21%	299	3.31%
	Beam shear center	290	/	290	/
	Beam bottom	299	3.10%	281	-3.01%
3	Deck centerline	168	-5.46%	188	5.73%
	Beam shear center	178	/	178	/

	Beam bottom	188	5.37%	169	-5.12%
5	Deck centerline	134	-6.84%	155	7.29%
	Beam shear center	144	/	144	/
	Beam bottom	154	6.82%	135	-6.42%

533

534 Table 3 Comparison between single-span and two-span twin-beam-deck systems

Beam span or length of each span (m)	Downward UDL			Upward UDL		
	Single-span beams (kNm) [1]	Two-span beams (kNm) [2]	[2]/[1]	Single-span beams (kNm) [3]	Two-span beams (kNm) [4]	[4]/[3]
4	205	202	0.985	252	291	1.12
6	203	204	1.01	222	240	1.08
8	200	202	1.01	210	222	1.06
10	198	199	1.01	204	212	1.04

535

536

537 8. Conclusions

538 A finite element formulation was developed for predicting LTB capacity of twin beams
539 continuously braced by deck boards. The solution captures the continuous relative
540 partial lateral restraint provided by the stiffness of deck-joint assembly and the partial
541 twist restraint provided by the deck bending stiffness. The results based on the present
542 solution are observed to be in close agreement with those based on ABAQUS while
543 involving a minimal computational effort. For wooden twin-beam-deck systems
544 investigated in the present study, the following observations were made:

- 545 1. The combined actions of continuous partial relative lateral and twist restraints
546 significantly increase LTB capacity. Compared with unrestrained beams, the
547 capacity increase can be as much as six folds. However, the provision of relative
548 lateral restraint alone was found ineffective in increasing the buckling capacity.
- 549 2. As the beam span is increased from 2 m to 10 m, the buckling capacity for upward
550 loading decreases by half while that based on downward loading remains

551 essentially constant. The buckling capacity for both upward and downward loads
552 declines significantly as the deck span is increased.

553 3. The load position greatly influences the buckling capacity for beam-deck systems
554 with short beam spans. Numerical results show that the buckling capacity can be
555 influenced by 10% to 23% when the load is moved away from beam shear center.
556 For beams with intermediate and long spans, the load position effect is below 5%.

557 4. The height of lateral restraint has no effect on the buckling capacity when the
558 relative lateral stiffness of the deck-joint assembly exceeds a threshold value.

559 5. A single nail connecting the beam to a 38 mm thick deck board provides lateral
560 stiffness that far exceeds the stiffness threshold to achieve maximum buckling
561 capacity. Additional measures to strengthen the beam-deck connections will not
562 lead to an increase in the buckling capacity.

563 6. The trend between beam span and buckling capacity for two-span twin-beam-deck
564 systems is similar to that of single span systems. The buckling capacity for two-
565 span systems is comparable to that of single-span systems with the same span.

566 **Notation**

567 A, B, C, D nodes located at beams upside or deck underside;
568 A_i, A_i' amplitudes of lateral displacement functions;
569 B_i, B_i' amplitudes of the angle of twist functions;
570 $[B_1], [B_2], [B_3], [B_4], [B_5]$ submatrices of the elastic and geometric stiffness matrices;
571 b deck board width;
572 C_w warping constant;
573 d_F nail diameter;
574 E_b, E_d modulus of elasticity of beam and deck, respectively;
575 E_L, E_T, E_R modulus of elasticity in longitudinal, tangential, radial directions
576 F lateral load at the beam-deck intersection;
577 F_u lateral strength resistance;
578 G_b shear modulus of the beam;
579 G_T shear modulus in transverse direction;

580	G_{LR}	shear modulus about longitudinal and radial direction;
581	G_{LT}	shear modulus about longitudinal and tangential direction;
582	G_{RT}	shear modulus about radial and tangential direction;
583	h	distance between loading point and beam shear center;
584	h_b, h_d	depth of beam and deck, respectively;
585	I_d	moment of inertia of a deck board in the strong-axis;
586	I_y	moment of inertia about beam weak-axis;
587	J_b	beam Saint-Venant torsional constant;
588	K_m	service creep factor;
589	$[K_b]$	beam stiffness matrix;
590	$[K_e]$	elastic stiffness matrix;
591	$[K_g]$	geometric stiffness matrix;
592	$[K_l]$	stiffness matrix for the relative lateral restraint;
593	$[K_t]$	stiffness matrix for the twist restraint;
594	k	relative lateral stiffness;
595	\bar{k}	relative lateral stiffness per unit deck width;
596	k_d	axial stiffness of deck board at the height of beam top;
597	k_j	shear stiffness at the beam-deck joint;
598	k_n	shear stiffness of a single-nail joint;
599	L_b, L_d	beam and deck span, respectively;
600	$\langle L(z) \rangle^T$	Hermitian polynomials;
601	M	reference strong-axis moment;
602	M_{cr}	critical moment
603	M_d	deck strong-axis moment under eccentric lateral loads;
604	M_{e1}, M_{e2}	deck board end moments;
605	M_0	deck strong-axis moment under unit eccentric lateral loads;
606	n	integer;
607	\bar{n}	number of nails at each joint;
608	$q(z)$	reference transverse load;
609	U	internal strain energy;
610	U_{bi}	internal strain energy in beam i ;
611	U_l	internal strain energy for the relative partial lateral restraint;
612	U_t	internal strain energy for the partial twist restraint;
613	U_l^*	internal strain energy for relative partial lateral restraint in one deck board;
614	U_t^*	internal strain energy for the partial twist restraint in one deck board;
615	u_A, u_C	lateral displacements for nodes A and C , respectively;
616	u_i	lateral displacement of beam i ;
617	$\langle u_i \rangle$	generalized nodal lateral displacement vector;

618	V	load potential energy;
619	V_{bi}	load potential energy for beam i ;
620	$v_p(z)$	prebuckling vertical displacement;
621	x, y, z	Cartesian coordinates;
622	z_0	distance from a given deck board to beam end-support;
623	Δ	deck lateral deformation;
624	Δ_1	deck axial deformation;
625	Δ_2	deck axial deformation due to transverse bending;
626	Δ_n	joint deformation of a single-nail joint;
627	θ_i	angle of twist of beam i ;
628	$\langle \theta_i \rangle$	generalized nodal angle of twist vector;
629	λ	load multiplication factor;
630	μ	Poisson's ratio;
631	μ_{LR}	Poisson's ratio in longitudinal and radial direction;
632	μ_{LT}	Poisson's ratio in longitudinal and tangential direction;
633	μ_{RT}	Poisson's ratio in radial and tangential direction;
634	Π	total potential energy.

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