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**Physical and Numerical Modeling of Large Marine Microplastics
Transport and Deposition**

by

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Abstract

The prevalence of micro-sized plastic pollution has raised significant concerns for aquatic ecosystems and their inhabitants, including humans, due to its distinct physical and toxicological characteristics. Large microplastics (MPs), primarily synthetic polymers measuring less than 5 mm, pose a considerable threat. According to the literature, rivers alone transport millions of tons of plastic into the ocean annually, regardless of other water bodies. During this transport, MPs undergo physical processes such as advection and diffusion, influenced by atmospheric and hydrological conditions. These processes are driven by larger-scale factors, including wind-induced drag, wave drift, and buoyancy, all stemming from currents, geostrophic circulation, and turbulence in extensive water bodies.

The inherent properties of MPs—such as size, shape, and density—along with fluid flow characteristics like velocity, depth, and pressure, play crucial roles in determining their transport and deposition patterns. Consequently, this thesis aimed to use physical and numerical modeling approaches to investigate MPs' transport and deposition under varying hydraulic parameters and bathymetric changes applied to the channel bed.

The experimental component of this research involved reviewing existing studies and conducting physical modeling of MPs transport and deposition. Various experimental test scenarios were carried out in a tilting flume at the National Research Council Canada Ocean, Coastal, and River Engineering Research Centre (NRC-OCRE) physical testing facility in Ottawa. Data on MPs transport and deposition were gathered using different physical setups in a straight open-channel flume to assess the effects of flow velocity, water levels, channel geometries, and particle shapes. This study contributed to validating numerical models for MPs transport and deposition.

The second part of the research, focusing on numerical modeling, followed three steps: (1) reviewing the literature, (2) using a machine learning (ML) approach to predict MPs transport and deposition dynamics for large spherical and cylindrical MPs, and (3) simulating MPs transport and deposition using a coupled computational fluid dynamics-discrete element method (CFD-DEM).

For the first step, a comprehensive review was conducted on the structure of four well-known Lagrangian particle tracking models (D-WAQ PART, Ichthyop, TrackMPD, and CaMPSim-3D) in simulating MPs' fate and transport. Each model was examined in terms of its ability to account

for key physical transport processes—such as advection, diffusion, windage, beaching, and washing off—as well as transformation processes like biofouling and degradation, which influence MPs' behavior in the marine environment. Additionally, the review evaluated the effects of MPs' physical properties (mainly size, diameter, and shape) on their fate and trajectories. This analysis identified gaps in particle tracking models, including insufficient consideration of homo- and hetero-aggregation, agglomeration, photodegradation, chemical and biological degradation, and additional advection caused by wave-induced drift.

In the second step of the numerical research, an ML-based study was conducted to predict the deposition patterns of spherical and cylindrical MPs using Extreme Learning Machines (ELM) for the first time, based on laboratory datasets. Various test scenarios were performed to collect data on multiple input parameters, including flow velocity, channel bed deepening, channel bed slope, water level, particle shape (spherical or cylindrical), near-bed flow velocity, and particle velocity at deposition. Eleven models, utilizing different dimensionless input combinations derived from Buckingham's theorem, were evaluated. The best-performing model incorporated all four dimensionless variables, achieving high accuracy for spherical ($R=0.94$) and cylindrical ($R=0.90$) MPs. Sensitivity analysis revealed that the ratio of water depth to particle diameter had the greatest impact on deposition patterns. This study highlighted the potential of ML in predicting MPs behavior, aiding environmental monitoring and management.

In the final step, numerical results were validated against laboratory experiments. A numerical modeling approach was employed to simulate the transport and deposition of spherical MPs under different flow and channel geometry scenarios. The unresolved CFD-DEM method was used for simulations, with results successfully validated through experimental measurements. Additionally, the model effectively captured the influence of hydraulic parameters on MPs deposition.

I would like to dedicate this thesis to my family.

To my dear mother and father, whose unwavering love and encouragement have supported me through every decision, even while I studied abroad. To my sisters, Mahsa and Maryam, for always being there for me. To my beloved wife, Mercedeh, whose belief in me and unwavering strength have been my greatest support. And to my son, Liam, who means the world to me.

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Publications from this research

The following publications have resulted from this research:

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- **Bigdeli, M.**, Mohammadian, A., Pilechi, A. (2022). “Numerical Modeling of Marine Microplastics Deposition using coupled CFD-DEM” accepted in 3rd IAHR Young Professionals Congress.
- **Bigdeli, M.**, Mohammadian, A., Pilechi, A. (2024). “A Laboratory Dataset on Transport and Deposition of Spherical and Cylindrical Large Microplastics for Validation of Numerical Models”, published by Journal of Marine Science and Engineering, MDPI.
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- **Bigdeli, M.**, Mohammadian, A., Pilechi, A., Bonakdari, H. (2025). “Dynamic Prediction of Large Spherical and Cylindrical Microplastic Deposition: A Machine Learning Approach for Transport and Deposition”, (submitted).
- **Bigdeli, M.**, Mohammadian, A., Pilechi, A. (2025). “Numerical Modeling of Large Spherical Microplastics Transport and Deposition using CFD-DEM”, (submitted).

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List of Abbreviations

ADV	Acoustic Doppler Velocimeter
AI	Artificial Intelligence
ANFIS	Adaptive Neuro-Fuzzy Inference Systems
ANN	Artificial Neural Networks
CaMPSim-3D	Canadian Microplastic Simulation
CFD	Computational Fluid Dynamics
Cyln.	Cylindrical
DEM	Discrete Element Method
D-WAQ PART	Delft3D - Water Quality Particle tracking module
ELM	Extreme Learning Machines
FFNN	Feedforward Neural Networks
FOV	Field Of View
GDP	Gross Domestic Product
Ichthyop	Ichthyoplankton
ML	Machine Learning
MPs	Microplastics
NRC-OCRE	National Research Council Canada Ocean, Coastal, and River Engineering Research Centre
PaTATO	Particle Tracking and Analysis Toolbox
PDSA	Partial Derivative Sensitivity Analysis
PEI	Prediction Interval Error
PET	Polyethylene Terephthalate
PTM	Particle Tracking Models
PVC	Polyvinyl Chloride
RA	Relative Accuracy
SLFFNN	Single-Layer Feedforward Neural Network
Sphr.	Spherical
SVM	Support Vector Machines

TE

TrackMPD

WUB

Trap Efficiency

Track Marine Plastic Debris

Width of Uncertainty Band

List of Symbols

d_p	Particle diameter
ν	Flow dynamic viscosity
w	Water depth in the channel
u_f	Flow velocity
h	Water depth in the channel at the deepening
s	Slope applied to the channel's bed
π_1	$\frac{w}{d_p}$
π_2	$\frac{h}{d_p}$
π_3	$\frac{u_f d_p}{\nu}$
π_4	s
R	Correlation Coefficient
MAE	Mean Absolute Error
$RMSE$	Root Mean Square Error
$BIAS$	Bias
e_i	Individual prediction error
\bar{e}	Mean prediction error
S_e	Standard deviation of prediction error
T_R	Rayleigh time
G	Shear modulus
Y	Young's modulus
R	Particle radius
ρ	Particle density
$t_{constact,H}$	Hertz contact time
ε_f	Void fraction
k	Turbulent kinetic energy
f_s	Sliding friction
f_r	Rolling friction

1 Introduction and Study Objectives

1.1 General Background

While most plastics are expected to float on the water's surface due to their lower density than water, research findings from (Kaiser et al., 2017) and (Zhang et al., 2019) have revealed a significant presence of plastic waste in the sediment beds of various water bodies, including coastal waters and rivers. Surprisingly, estimates suggest that around 99% of all plastics that have entered the world's oceans are unaccounted for, possibly having been deposited in sediments or deep ocean waters. A global analysis of plastic production and often unregulated disposal by (Geyer et al., 2017) projects that plastic debris accumulation in natural environments and landfills could reach a staggering 12,000 million metric tons (Mt) by 2050.

Rivers are recognized as one of the primary conduits for transporting plastic waste from land to oceans, with an estimated annual rate ranging from 0.41 to 4.0 Mt (Corcoran et al., 2019). Previous studies have emphasized that the mobility of micro-sized particles in aquatic environments is influenced by various factors, including particle-specific characteristics (such as size, density, shape, chemical composition), water chemistry, surface charge, surface coating, hydrodynamic conditions, and sediment properties like grain size distribution and organic matter content (Klaine et al., 2008; Masliyah, 1998; Wiesner & Bottero, 2017). It is challenging to accurately estimate the plastic load within sediments due to the absence of standardized data analysis procedures that account for the inherent variability in natural conditions and limited density of available samples. To make valid comparisons across different spatial and temporal scales, it is essential to consider hydrodynamic conditions at a local level. This consideration is crucial because the variations in microplastic abundance observed at small spatial scales can be more pronounced than those seen at larger spatial scales, as indicated by studies like (Enders et al., 2019; Martin et al., 2017). Further details have been provided in section 1.9.

Physical and numerical modeling serve as valuable tools for predicting the fate and transport of plastic debris, offering a potential solution to mitigate the plastic pollution problem. Therefore, it is crucial to thoroughly characterize and quantify the deposition of plastic debris in bed materials, especially when developing numerical models to simulate the fate and transport of plastic debris

30 in water bodies. This information is vital for tasks like mass balance calculations and estimating
31 the transport of plastic debris versus its deposition.

32 A comprehensive understanding of MPs' behavior in water resources is essential due to the
33 following challenges:

34 1. Widespread Distribution & Environmental Impact:

35 The high production rate and rapid industrial growth of plastics pose a substantial threat to marine
36 ecosystems and biodiversity, leading to severe pollution. The transport and deposition of MPs can
37 adversely affect water quality, making it unsuitable for drinking or other uses.

38 2. Bioaccumulation & Health Risks:

39 MPs can be absorbed by marine fauna, entering the food chain and persisting without degradation
40 (Alvarez-Zeferino et al., 2020). The distribution and accumulation of MPs—given their diverse
41 types, origins, and compositions—raise serious concerns. MPs can infiltrate the body tissues of
42 marine species and human lymph, liver, kidney, and circulatory systems (Vethaak & Legler, 2021).
43 Additionally, MPs can serve as platforms for the growth of harmful biofilms (Vethaak & Legler,
44 2021). It is, therefore, crucial to identify high-accumulation zones, assess associated risks, and
45 implement mitigation strategies.

46 3. Lack of Comprehensive Studies:

47 Limited reliable and comprehensive studies on MPs in water bodies hinder our ability to
48 understand their behavior and plan effective mitigation and remediation strategies. Existing studies
49 have primarily focused on collecting, pretreating, and analyzing MPs (Ballent et al., 2016;
50 Bujaczek et al., 2021; Castañeda et al., 2014; Corcoran et al., 2019; Vermaire et al., 2017).

51 4. Unclear Deposition & Accumulation Patterns:

52 Despite significant findings on the abundance, shape, color, and size of MPs (Ballent et al., 2016;
53 Bujaczek et al., 2021; Castañeda et al., 2014; Corcoran et al., 2019), their deposition and
54 accumulation patterns remain largely unknown, limiting our understanding of their environmental
55 impact on marine ecosystems.

56

57 5. Unresolved Deposition Characteristics:

58 The deposition behavior of MPs is not yet fully understood. Assessing water bodies using
59 laboratory experiments, modeling approaches, and numerical simulations can provide critical
60 insights into MPs' deposition mechanisms, making it a worthwhile research endeavor.

61 **1.2 Research Questions**

62 Several key questions arise when considering the transport and deposition of MPs in water bodies:

63 1. Which hydraulic factors have the greatest influence on the transport and deposition of MPs in
64 water bodies?

65 2. In the context of changes in water body bathymetry, how do changes in bathymetry impact
66 particle transport and deposition?

67 3. Is ML capable of sufficiently predicting the behavior of MPs under the influence of different
68 hydraulic flow conditions and different physical characteristics of the MPs?

69 4. How effective are numerical models in capturing the precise distribution and deposition of MPs
70 in water bodies?

71 **1.3 Research Objectives and Significance**

72 To address these questions comprehensively, it is imperative to establish research objectives in a
73 manner that facilitates solutions to the aforementioned concerns and contributes to the existing
74 body of knowledge. Accordingly, the research study's objectives are delineated as follows:

75 1. Develop and construct a physical model and conduct experimental trials to:

- 76
- 77 • Investigate the role of both transport and deposition in shaping the fate of marine MPs.
 - 78 • Identify the pivotal hydraulic parameters, considering water bodies' flow characteristics
79 that govern the transport and deposition of MPs particles.
 - 80 • Determine how bathymetric changes in water bodies contribute to the formation of dead
81 zones where MPs accumulate.

81 2. Verify and implement numerical simulations by integrating CFD-DEM model to:

- 82
- Explore the transport and deposition of MPs.

83 • Track particle behavior in turbulent marine conditions.

84 3. Evaluate the capability of the ELM model in predicting MPs deposition patterns based on a
85 comprehensive and reliable experimental dataset derived from controlled test scenarios.

86 **1.4 Research Outcomes**

87 The main results of this research can be summarized as follows:

- 88 1. Enhance our comprehension of the dynamics of marine MPs transport and deposition in
89 water bodies, considering different bathymetric conditions.
- 90 2. Present a robust and comprehensive physical model for a better understanding of the MPs’
91 behavior in real world conditions.
- 92 3. Provide the capacity of CFD-based and AI-based numerical models to simulate and predict
93 the transport and deposition of marine MPs.
- 94 4. Disseminate the research findings through the publication of scientific papers in peer-
95 reviewed journals.

96 **1.5 Novelty and Contribution of the Research**

97 In order to meet the research study objectives, there are some novel contributions that were verified
98 by lab tests. The main novelties are described below:

- 99 1. Elucidating the dynamic of marine MPs transport and deposition within the channel of water
100 bodies
- 101 2. Comprehending the influence of varied hydraulic conditions in water bodies on MPs transport
102 and deposition with different physical characteristics.
- 103 3. Identifying the most dominant parameter(s) governing the transport and deposition of MPs in
104 the channel of water bodies.
- 105 4. Determine the accuracy of CFD-based and AI-based models in simulating MPs transport and
106 deposition. This represents a state-of-the-art methodology in this field of study.

107 **1.6 Research Scope**

108 MPs present in water undergo various transformational processes, including biofouling (Kooi et
109 al., 2017), degradation, and fragmentation (Atugoda et al., 2022; Cole et al., 2011; Nikpay, 2022),
110 as well as hetero- and homo-aggregation (Atugoda et al., 2022; Besseling et al., 2017), commonly

111 referred to as particle-particle interactions (Van Sebille et al., 2020). These processes significantly
112 impact on the destiny of MPs within their receiving environment. Furthermore, MPs have become
113 widely distributed throughout marine ecosystems (Eriksen et al., 2014; Galgani et al., 2015a) due
114 to various physical transport mechanisms, including horizontal and vertical advection, diffusion,
115 and other factors. The transport of MPs in aquatic environments is influenced by a range of
116 processes, including advection and dispersion driven by ambient currents, Stokes drift, sinking,
117 resuspension, beaching, washing-off, deposition, aggregation, degradation, and fouling.

118 However, it's important to note that this study's specific focus is on addressing the existing
119 knowledge gap regarding the transport and deposition of marine MPs, particularly those
120 characterized by different bathymetric conditions. As a result, addressing transformational
121 processes or other factors influencing MPs transport and deposition in water bodies is beyond the
122 scope of this research. This limitation does not affect considerations related to the hydraulic
123 conditions of the flow and the physical properties of the particles.

124 **1.7 Thesis Outline**

125 The remainder of this thesis is structured in the following way. Section 2 presents a literature
126 review of the physical characteristics of MPs and the processes they undergo, as well as the
127 Lagrangian modeling of the fate and transport of marine MPs. Section 3 presents an experimental
128 study on the transport and deposition of large spherical and cylindrical MPs that can be employed
129 to validate numerical models. Section 4 discusses the dynamic prediction of large spherical and
130 cylindrical MPs deposition using the ELM. Section 5 deals with numerical modelling of the large
131 spherical MPs transport and deposition using CFD-DEM and Section 6 offers essential insights as
132 the concluding remarks.

133 **1.8 List of Publications (from the thesis)**

- 134 • **Bigdeli, M.**, Mohammadian, A., Pilechi, A., Taheri, M. (2022). “Lagrangian modeling of
135 marine microplastics fate and transport: the state of the science”, published by Journal of
136 Marine Science and Engineering, MDPI.
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143 Numerical Models”, published by Journal of Marine Science and Engineering, MDPI.
144 <https://doi.org/10.3390/jmse12060953>
- 145 • **Bigdeli, M.**, Mohammadian, A., Pilechi, A., Bonakdari, H. (2025). “Dynamic Prediction
146 of Large Spherical and Cylindrical Microplastic Deposition: A Machine Learning
147 Approach for Transport and Deposition”, (submitted).
- 148 • **Bigdeli, M.**, Mohammadian, A., Pilechi, A. (2025). “Numerical Modeling of Large
149 Spherical Microplastics Transport and Deposition using CFD-DEM”, (submitted).

150 **1.9 Appendix**

151 There are substantial disparities in the findings regarding whether MPs abundance is linked to
152 sediment composition, including grain size and organic matter content. Some studies, such as
153 (Strand et al., 2013) have suggested a potential relationship, while others, such as (Renzi et al.,
154 2018) have failed to establish such a connection. The existence of this controversy in literature
155 underscores the existing gaps in our knowledge regarding the possible association between MPs
156 and sediment.

157 MPs dynamics have previously been studied primarily in terms of rising and settling processes
158 (Yu et al., 2022). Considering the density of half of all plastics is higher than that of seawater
159 (Morét-Ferguson et al., 2010), the terminal settling velocity is a key parameter used to characterize
160 the migration of MPs into water bodies (Khatmullina and Isachenko, 2017). The density, shape,
161 and particle size of MPs have been found to have a significant impact on the terminal settling
162 velocity (Yu et al., 2022).

163 There is also a relationship between surface wettability and the settling velocity and drag
164 coefficient of micro-sized particles (Wang et al., 2020). In this respect, hyporheic zones have also
165 been investigated for the settling mechanisms of MPs. Thus, stagnant zones are vital reservoirs for
166 the long-term sink of MPs, which are remobilized and resuspended at higher flow rates (Kumar et
167 al., 2021). It has been observed that MPs with varying particle diameters settle at different speeds

168 (Kumar et al., 2021). In addition to natural sediments, MPs particles' settling and rising velocities
169 significantly affect their erosivity (Waldschläger and Schüttrumpf, 2019a). MPs settle and
170 sediment primarily as a result of gravitational forces (Frei et al., 2019). Also, MPs have been
171 observed to settle in rivers that have sloped beds, i.e., as flow velocity increases, MPs travel over
172 longer distances (Kumar et al., 2021). However, according to (Hoellein et al., 2019), there is a
173 positive correlation between MPs diameter and deposition velocity (V_{dep}), as opposed to a negative
174 correlation between deposition velocity and settling velocity (V_{fall}). Particles with diameters
175 between 50 and 500 meters are deposited under the influence of gravity when V_{dep}/V_{fall} ratio is
176 equal to 1. By contrast, when $V_{dep}/V_{fall} \gg 1$, particles with a diameter < 50 m have a high retention
177 rate while for $V_{dep}/V_{fall} \ll 1$, particles with a diameter > 500 μm are influenced by hydrologic
178 actions including momentum and resuspension which are both larger than gravitational settlings
179 (Hoellein et al., 2019).

180 As stated above, MPs and sediments are different in terms of the physical properties (e.g., density,
181 size, and shape) that affect their transport. The application of the basic and governing principles
182 of sediment transport to MPs transport has therefore remained in doubt for this reason and due to
183 the lack of sufficient knowledge of the prevailing transport mechanisms. Additionally, it is also
184 important to determine whether it is possible to calculate the terminal settling and rising velocities
185 of MPs, which have physical properties different from sediments, using theoretical methods that
186 are derived from sediment transport models. In order to investigate the aforesaid concerns, the
187 nutshell of the reviewed papers are summarized as below:

188 • It has been reported that there is often a correlation between MPs and sediment grain size.
189 Different conclusions have been reached regarding the impact of sediment composition, such as
190 grain size and organic matter content, on the abundance of MPs. Some of them reported indications
191 of such a relationship, but others were unable to prove it. There are currently knowledge gaps
192 regarding the potential relationship between MPs and sediment as a result of a controversy of this
193 nature in literature.

194 • The bed-load transport phenomenon occurs when successive particle contacts with the bed
195 are strictly limited by gravity while during suspended load transport, turbulent eddies impart
196 upward impulses to support all the excess weight of the particles. Depending on the particle size,

197 weight, shape, and ambient hydraulic conditions, an individual particle may be transported as a
198 suspended load or as a bedload.

199 • In partial transport, a type of bed load transport, finer material is mobilized, and coarser
200 material remains immobile when the bed shear stress is too low to mobilize it.

201 • Due to the fact that the start of a particle trajectory is determined by random fluid forces,
202 the length and height of the trajectory are also random variables. The saltation model can be used
203 to determine the heights, lengths, and velocities of hopping grains.

204 • Different MPs settle at different rates depending on their particle diameter. As well as
205 natural sediments, MPs particles' settling and rising velocities have a significant impact on their
206 erosivity.

207 • Given the aforementioned and due to the similar geometric characteristics between MPs
208 and non-cohesive sediment particles, principles associated with sediment incipient motion can be
209 applicable to MPs as well (Waldschläger and Schüttrumpf, 2019b, Yu et al., 2022). MPs particles
210 begin to move under the action of currents after they settle on the riverbed or seafloor. Incipient
211 sediment motion is extensively governed by the critical shear stress, which is the minimum flow
212 shear stress required for a single particle to move.

213 • Due to the hiding exposure effect, large grains may be exposed on the far side of the layer
214 as a result of MPs of smaller dimensions being layered over them. By varying the hiding-exposure
215 coefficient, the hiding-exposure effect can be fitted.

216 As reviewed in earlier sections, most previous studies have focused on applying sediment transport
217 equations to MPs by considering only the onset of sediment motion (incipient motion). However,
218 this leaves a significant knowledge gap regarding how MPs behave when they interact with the
219 bed surface. This study aimed to address that gap by providing deeper insights into how MPs are
220 trapped, rolled, slid, or saltated across a reliable range of conditions, using precise experimental
221 testing. Subsequently, numerical models were calibrated based on these experimental
222 observations.

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339 **2 Lagrangian Modeling of Marine Microplastics Fate and Transport: The** 340 **State of the Science**

341 **2.1 Abstract**

342 Microplastics pollution has led to irreversible environmental consequences and has triggered
343 global concerns. It has been shown that water resources and marine food consumers are adversely
344 affected by microplastics due to their physico-chemical characteristics. This study attempts to
345 comprehensively review the structure of four well-known Lagrangian particle-tracking models,
346 i.e., Delft3D - Water Quality Particle tracking module (D-WAQ PART), Ichthyoplankton
347 (Ichthyop), Track Marine Plastic Debris (TrackMPD), and Canadian Microplastic Simulation
348 (CaMPSim-3D) in simulating the fate and transport of microplastics. Accordingly, the structure of
349 each model is investigated with respect to addressing the involved physical transport processes
350 (including advection, diffusion, windage, beaching, and washing-off) and transformation
351 processes (particularly biofouling and degradation) that play key roles in microplastics' behavior
352 in the marine environment. In addition, the effects of the physical properties (mainly size,
353 diameter, and shape) of microplastics on their fate and trajectories are reviewed. The models'
354 capabilities and shortcomings in the simulation of microplastics are also discussed. The present
355 review sheds light on some aspects of microplastics' behavior in water that were not properly
356 addressed in particle-tracking models, such as homo- and hetero-aggregation, agglomeration,
357 photodegradation, and chemical and biological degradation as well as additional advection due to
358 wave-induced drift. This study can be regarded as a reliable steppingstone for the future
359 modification of the reviewed models.

360 **2.2 Introduction**

361 Aquatic environments, including oceans, coastal regions, estuaries, seas, and rivers, as major
362 planetary water resources, have been affected by marine plastic debris accumulation and
363 distribution, which has triggered global environmental concern and irreversible pollution
364 consequences (Bergmann et al., 2015; Panno et al., 2019; Wagner & Lambert, 2018). Macro (>
365 1 cm) (N. Hartmann et al., 2015), meso (5 mm – 2 cm) (N. Hartmann et al., 2015), micro (<
366 5 mm) (Rocha-Santos & Duarte, 2015), and nano-size (1 nm – 100 nm) (EU, 2019b) plastics
367 pollution has become a critical environmental problem due to their persistence, toxicological
368 properties, and destructive effects on not only Earth's hydrosphere (Barnes et al., 2009) but also

369 aquatic organisms (Koelmans et al., 2016; Ryan et al., 2009; Thevenon et al., 2014a; Wright et al.,
370 2013), wildlife, and human health (Lithner et al., 2011; Napper & Thompson, 2019; Rochman et
371 al., 2013; Wright & Kelly, 2017). The tendency of microplastics to adsorb and react with
372 wastewater agents is relatively high due to their physical properties, such as hydrophobicity
373 (Nikpay, 2021), as well as the large ratio of their surface area to volume, allowing them to adsorb
374 toxic pollutants (e.g., heavy metals) and invasive species (Bakir et al., 2014). Microplastics are
375 classified into two categories: primary and secondary. Primary microplastics are mainly detected
376 with an original size of less than 5 mm in textiles and some personal care products (Bergmann et
377 al., 2015; Cole et al., 2011). They can be transported through discharges of water treatment plants,
378 rivers, surface run-off, and wind into aquatic environments (Gall & Thompson, 2015). Large
379 plastic debris under fragmentation may produce secondary microplastics through some processes
380 including physical, chemical, and biological interactions and photo-degradation (Galgani et al.,
381 2013; Thompson et al., 2009). The secondary microplastics originate mainly from industrial resin
382 pellets, fishing nets, and other thrown-away plastic litters (Eerkes-Medrano et al., 2015), and they
383 can also be formed from the fragmentation of any larger plastic particle in the ocean, regardless of
384 its source. Microplastics in the environment majorly consist of secondary microplastics (Eriksen
385 et al., 2013). The possibility of microplastics breaking down into nanoplastics may increase during
386 their exposure to the environment, leading to higher environmental risks given the nano-size nature
387 of these materials (Li et al., 2018). In addition, plastic pollution impacts have been observed at
388 economic and social scales (Jeftic et al., 2009). The tourism industry accounts for a considerable
389 percentage of the gross domestic product (GDP) of many countries, which indicates the market
390 value of all services provided within a country's territory during a specific temporal period; thus,
391 governments schedule various plans to attract more tourists and introduce different sightseeing
392 places in order to boost tourism income. Nevertheless, beaches confront some challenges with
393 respect to their sustainability in terms of preservation and operation from the touristic and
394 economic perspectives. The uncontrolled release of microplastics causes the pollution of the
395 seawater and the shore.

396 Plastic debris enters the ocean directly or through rivers by uncontrolled dumping from land-based
397 or maritime sources, such as fishing or shipping sectors, and recreational activities (Galgani et al.,
398 2015b; Ivar do Sul & Costa, 2013; Li et al., 2020). Plastics will be degraded and broken down into
399 microplastics (i.e., particles less than 5 mm in size) after entering the ocean if not treated at early

400 stages. Microplastics are found throughout the water column in oceans and also on the seabed
401 (Desforges et al., 2014; Woodall et al., 2014). Numerous studies have been conducted to quantify
402 and characterize microplastics in aquatic environments (Eerkes-Medrano et al., 2015; Eriksen et
403 al., 2013; Hartmann et al., 2015; Thompson et al., 2009) including lakes (Eriksen et al., 2013; Free
404 et al., 2014), rivers (Klein et al., 2015; Kooi et al., 2016; Lechner et al., 2014; McCormick et al.,
405 2014; Yonkos et al., 2014), and shorelines (Browne et al., 2011) by employing different field
406 surveys and using a set of lab experiments. Recent advances in computer science and technology
407 have also promoted the application of Machine Learning (ML) approaches in the automated
408 quantification and characterization of microplastics (Girshick, 2015; Girshick et al., 2014; Liu et
409 al., 2016; Lorenzo-Navarro et al., 2021; Redmon et al., 2016; Ren et al., 2015).

410 Microplastics in water undergo some transformational processes such as biofouling (Kooi et al.,
411 2017), degradation and fragmentation (Atugoda et al., 2020; Cole et al., 2011; Nikpay, 2022), and
412 hetero- and homo-aggregation (Atugoda et al., 2020; Besseling et al., 2017), known also as
413 particle–particle interactions (Van Sebille et al., 2020), which influence their fate in their receiving
414 environment. Accordingly, their physical (mainly size, density, and shape), chemical, and
415 biological properties may change as a result of their ambient environment properties, including
416 sea surface temperature, sunlight, salinity, etc. The growth of ultrafine bacteria or microorganisms
417 and the attachment of organic material, invertebrates, and algae on the exterior surfaces of these
418 particles cause biofouling that results in not only density and size increases in the particles but
419 consequently decreases buoyancy (Rummel et al., 2017). Thus, it can be concluded that the sinking
420 process can be accelerated by biofouling (Ryan, 2015) and may result in the sedimentation of
421 fouled particles. In addition, microplastics have been widely pervasive in marine habitats (Eriksen
422 et al., 2014; Galgani et al., 2015b) due to physical transport processes including horizontal and
423 vertical advection and diffusion. Previous studies have shown that microplastics have reached very
424 remote regions in the Arctic (Obbard et al., 2014) and Antarctic (Lusher et al., 2015; Waller et al.,
425 2017). They can be found on the sea surface (Cózar et al., 2014), the seabed (Galgani et al., 2000),
426 and also in deep-sea sediments (Courtene-Jones et al., 2020; Peng et al., 2020). Predicting such
427 transport processes, given their dependencies on different hydrological and atmospheric variables
428 spanning over a wide temporal and spatial scale, is important for the accurate simulation of
429 microplastics' fate and transport in water bodies. The transport of microplastics in water bodies is
430 controlled by advection and dispersion driven by ambient currents, Stokes drift, sinking,

431 resuspension, beaching and washing-off, deposition, aggregation, degradation, and fouling, which
432 need to be simulated by numerical models.

433 Four sources for transport processes can be considered: (1) ambient currents (e.g., tidal and wind-
434 driven flows), leading to transport by advection; (2) wind-induced drift, i.e., windage, which is an
435 additional advection term because of wind drag and has an undeniable effect on plastic particles'
436 displacement (Van Utenhove, 2019)—in other words, the wind-driven drag force exerted on the
437 water's surface causes the drifting of floating particles exposed to wind; (3) wave action; and (4)
438 buoyancy, which influences the vertical transport. Indeed, the influence of buoyancy on vertical
439 transport and associated vertical mixing can be significant (Zhang, 2017).

440 Physical processes involved in the transport and fate of microplastics can occur at different spatio-
441 temporal scales, which makes the simulation challenging. In this respect, short-term processes
442 (e.g., particle–particle interactions at local and regional scales) may last for several hours or a few
443 days (Gouin et al., 2019); however, medium-term processes including biofouling, sedimentation,
444 and advection from freshwater bodies to marine environments may occur within several days or
445 weeks (Kooi et al., 2017). Long-term processes need months or even years to be completed (Jalón-
446 Rojas et al., 2019). For instance, some processes, such as degradation (or fragmentation), sinking,
447 and transport (due to advection and diffusion) from coastal water bodies to oceans are long-term
448 processes that increase computational costs. From the spatial scale perspective, the simulation of
449 physical processes in large study areas with global scales is of orders of thousands of square
450 kilometers, while the computational domains for regional and local areas are of orders of hundreds
451 and tens of square kilometers, respectively.

452 Numerical simulation has been widely employed in the past decades to predict marine
453 microplastics' behavior (Ballent et al., 2012; Critchell et al., 2015; Critchell & Lambrechts, 2016;
454 Declerck et al., 2019; Díez-Minguito et al., 2020; Pereiro et al., 2019; Raimundo et al., 2020).
455 There are also some studies that have predicted the main sources, transport, and fate of
456 microplastics under different environmental conditions based on numerical simulations (Isobe et
457 al., 2019; Lebreton et al., 2012; Liubartseva et al., 2016; Maximenko et al., 2012; Van Sebille et
458 al., 2015). Some of these studies have employed particle-tracking (Jalón-Rojas et al., 2019) and
459 hydrodynamic models (Bondelind et al., 2020a), individually or coupled with each other (Alosairi
460 et al., 2020; Collins & Hermes, 2019; Sousa et al., 2021). Numerical modeling is inexpensive by

461 far in comparison with experimental measurements and in situ observations and can predict
462 potential accumulation zones in remote areas where water quality monitoring is not feasible.
463 Furthermore, numerical models can simulate some key physical processes using fine grid
464 resolution, given the ever-increasing computational capacity and the ongoing development of
465 numerical techniques. Moreover, these models are capable of simulating different vertical and
466 horizontal flow patterns encompassing microplastics with good accuracy.

467 Numerical simulations of microplastics' fate and transport in water are generally carried out using
468 Eulerian and Lagrangian models. Unlike Lagrangian trajectory models that use a movable
469 coordinate system in space, Eulerian grid models typically employ a non-movable frame of
470 reference (Van Utenhove, 2019). In other words, a Lagrangian framework, known as the particle-
471 tracking method, tracks only individual particles, while an Eulerian approach addresses advection
472 and diffusion for given locations (computational nodes) (Van Sebille et al., 2018). Numerical
473 Eulerian hydrodynamic models can be coupled with Lagrangian models. Accordingly,
474 hydrodynamic modeling provides the ambient hydrological parameters as inputs for particle-
475 tracking models (PTM), which are known as promising, cost-effective tools for interpreting the
476 fate and transport of microplastics and providing better insights into microplastics' behavior in
477 dynamic flow systems (Hardesty et al., 2017). Although some studies have used an Eulerian
478 approach for the mixing and dispersion processes in dynamic water systems (Alosairi &
479 Alsulaiman, 2019), some other studies have used Lagrangian models for tracking particles to
480 investigate plastic particles' fate and transport (Van Utenhove, 2019). Lagrangian models solve
481 for a lower number of equations per node so that they do not need to solve the equations for each
482 location in the mesh. However, given the pros and cons of both Eulerian and Lagrangian
483 approaches, they can be employed to complement each other. The particle-tracking model coupled
484 with the hydrodynamic model can provide invaluable details about the potential trajectories of
485 plastic litter from local (Lammerts, 2016; Rodríguez-Díaz et al., 2020; Stuparu et al., 2015) to
486 global scales (Lebreton et al., 2012) that lead to the assessment of microplastics' transportation,
487 distribution, residence time, and fate under various conditions (Carlson et al., 2017; Critchell &
488 Lambrechts, 2016; Ebbesmeyer et al., 2012; Isobe et al., 2009; Lebreton et al., 2012; Liubartseva
489 et al., 2018; Van Utenhove, 2019).

490 Particle-tracking methods simulate the trajectory of particles and their transport in water bodies.
491 They are applicable to a wide range of plastic particles' transport in their surrounding environment
492 (Lynch et al., 2014; Van Sebille et al., 2018). PTM has been employed in many studies as a reliable
493 alternative for field surveys. This is due to the logistic challenges involved in field surveys
494 (Hoffman & Hittinger, 2017; Lebreton et al., 2012; Mason et al., 2020; Maximenko et al., 2012;
495 Van Sebille et al., 2012; Vennell et al., 2021). Most Lagrangian particle-tracking models do not
496 simulate the flow hydrodynamics. They can be considered as an important tool to simulate physical
497 processes and transport mechanisms that microplastics undergo.

498 Unlike the structured mesh, the unstructured mesh enables numerical models not only to mesh
499 complex geometries more easily but also to address arbitrary positions, even though it should be
500 noted that greater memory is needed for executing the simulation for unstructured mesh.
501 Hydrodynamic values in each element of the Eulerian mesh are interpolated in the same location
502 of the particle (as the host element) in PTM.

503 On the other hand, the spatio-temporal interpolation of the flow characteristics (mainly including
504 the velocity, diffusivities, salinity, sea surface height, temperature, etc.) in both horizontal and
505 vertical directions may be necessary, especially when the duration between successive
506 hydrodynamic outputs is different in comparison with the time step of particles' motion. Thus, the
507 hydrodynamic information provided on an Eulerian grid, as an input to the PTM model, should be
508 interpolated for the spatio-temporal location of the particles before being read by the PTM (Jalón-
509 Rojas et al., 2019). For instance, the particle's velocity at each location is found using the spatial
510 and temporal interpolation of the hydrodynamic model's velocity.

511 Most available numerical models have limitations in considering all effective factors relevant to
512 microplastics' fate and transport. For example, some models have ignored beaching, sinking, and
513 fragmentation processes as well as biofouling, deposition, resuspension, washing-off from the
514 beach, or wind drift (Critchell & Lambrechts, 2016; Ebbesmeyer et al., 2012; Murray et al., 2018).

515 Moreover, some physical properties of microplastics, including shape, size, and density, which are
516 dominant factors related to microplastics' buoyancy and mobility (Browne et al., 2010;
517 Chubarenko et al., 2016; Filella, 2015; Zhang, 2017), have not been parameterized in some
518 numerical models. Furthermore, most of the aforementioned models use a 2D approach, neglect
519 the dynamic behavior of microplastics, and only model marine debris neutral drifting within the

520 surface layers (Carlson et al., 2017; Ebbesmeyer et al., 2012; Isobe et al., 2009; Lebreton et al.,
 521 2012; Liubartseva et al., 2018; Murray et al., 2018).

522 Among the numerical models that are employed for the simulation of marine microplastics’
 523 behavior, we studied four PTMs, i.e., D-WAQ PART, Ichthyop, TrackMPD, and CaMPSim-3D.
 524 D-WAQ PART is the particle-tracking module of Delft-3D, which is an open-source simulating
 525 software equipped with different modules for predicting hydrodynamics in an aquatic
 526 environment. Ichthyop, as a free Java-based numerical model, has been employed for tracking
 527 plastic particles, although it has been originally used for studying the dynamics of ichthyoplankton.
 528 TrackMPD has been developed in recent years to simulate microplastic debris fates and transport.
 529 The latest developed numerical model that can be employed as a particle-tracking model is
 530 Canadian microplastic simulation (CaMPSim). It utilizes a coupling of an Eulerian and Lagrangian
 531 numerical model in the simulation of marine microplastics’ behavior.

532 The current paper systematically reviews recent progress in the numerical modeling of the
 533 transport and distribution of microplastics in marine environments with the aims of (1) reviewing
 534 the structure of Lagrangian particle-tracking models, (2) determining the substantial improvements
 535 that are still needed, and (3) discussing the consideration of physical processes that microplastics
 536 undergo and their physical properties in the simulation of microplastics’ fate and transport. In this
 537 respect, different 2D and 3D numerical models are categorized, as shown in Table 1.

538 **Table 1.** The most popular numerical 2D and 3D models used for simulating microplastics’ fate
 539 and transport.

Particle-Tracking Model	Hydrodynamic Model	Dimensionality	Method	Simulating Fate or Transport	Ref.
D-WAQ PART	Delft3D-FLOW	2D/3D	Lagrangian	Transport	(Alosairi et al., 2020; Sousa et al., 2021)
Ichthyop	ROMS	2D/3D	Lagrangian	Transport	(Collins & Hermes, 2019;

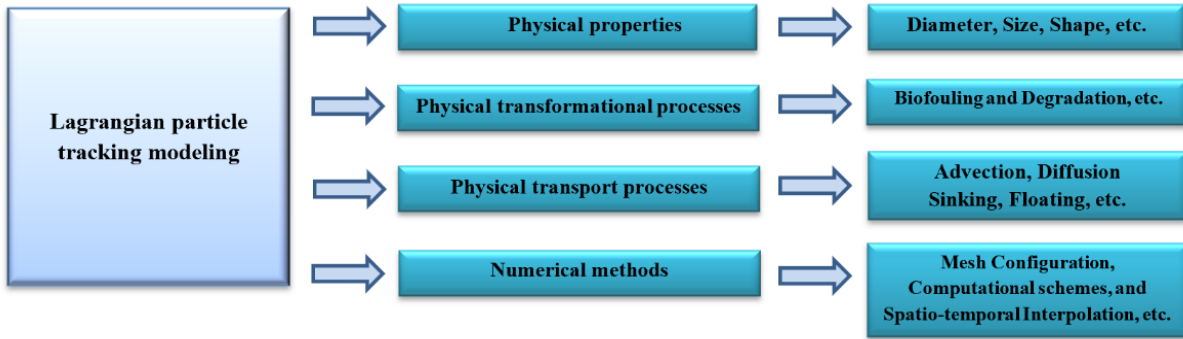
Particle-Tracking Model	Hydrodynamic Model	Dimensionality	Method	Simulating Fate or Transport	Ref.
					Lett et al., 2008)
TrackMPD	OGCM ¹	3D	Lagrangian	Both	(Jalón-Rojas et al., 2019)
CaMPSim-3D	TELEMAC	3D	Lagrangian	Both	(Pilechi et al., 2022)

540 1 Ocean General Circulation Models.

541 2.3 Materials and Methods

542 As stated above, the concept and framework of different numerical models for modeling
543 microplastics' fate and transport are reviewed in this study in detail, considering various ranges of
544 microplastics' behavior characteristics and key physical processes affecting their variations.
545 Lagrangian particle-tracking models are investigated in terms of considering physical properties
546 (and hydrodynamic aspects), physical transformational processes, physical transport processes,
547 and numerical schemes (Figure 1). In other words, the reviewed models have been compared based
548 on capability in considering the physical properties of microplastics (mainly size, shape, and
549 density), and in simulating the physical transformational and transport processes that microplastics
550 undergo. Moreover, the numerical methods that have been employed by each of them in terms of
551 considering the type of mesh they can read, their computational schemes, and the methods they
552 employ for implementing interpolation have been addressed. Finally, limitations and gaps in the
553 reviewed models and needed areas of improvement are determined and discussed in full.

554



555

556 **Figure 1.** The structure of Lagrangian particle-tracking models in the simulation of microplastics’
 557 physical properties, physical transformation processes, physical transport processes, and numerical
 558 methods.

559

560 **2.4 Physical Properties and Processes**

561 **2.1.1 Physical Properties**

562 The behavior and transport of microplastics mainly depend on major physical properties including
 563 the size, density, and shape (Ballent et al., 2012; Chubarenko et al., 2016; Khatmullina &
 564 Isachenko, 2017; Kooi et al., 2017) of the particles. Microplastics are identified in varied sizes and
 565 shapes, and denser particles normally settle faster even with similar shapes and sizes (Kooi et al.,
 566 2017). Moreover, larger particles in terms of size but with the same density experience an increase
 567 in the ratio of the gravitational force applied on the particle to the fluid viscous resistance, which
 568 results in higher settling velocities and faster settling (Jalón-Rojas et al., 2019). In addition, the
 569 shape of microplastics can affect their movements and settling velocity as well (Isachenko et al.,
 570 2016). It has been observed that elongated particles such as films and fibers, known as larger
 571 particles, settle faster than smaller, circular particles (Chubarenko et al., 2016). Microplastics come
 572 in various regular (e.g., spherules, cylinders, and beads) and irregular shapes with different
 573 geometries (pellets and fibers). Regular shapes usually belong to primary microplastics while
 574 irregular shapes often can be seen among the secondary particles as a result of degradation
 575 (Khatmullina & Isachenko, 2017). Moreover, the physical properties of microplastics undergo
 576 changes due to the impacts of different processes and ambient factors such as algae invasion (Nava
 577 & Leoni, 2021), salinity (Liu et al., 2020), UV index (Miranda et al., 2021), temperature (Miranda
 578 et al., 2021), and Stokes drift generated by wave action (De Leo et al., 2021). For instance, unlike

579 freshly emitted particles, the edges of microplastics become smooth during aging as a result of
 580 mechanical degradation since they are under polishing imposed by other fragments (Doyle et al.,
 581 2011). Different equations used in the reviewed models for modeling hydrodynamic aspects
 582 (including microplastics' physical properties) pertinent to the transport of microplastics have been
 583 summarized in this study in Table 2.

584 **Table 2.** Various equations used for the modeling of hydrodynamic aspects in the reviewed
 585 models. *

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
1	$v_s = c^n \times \left[A_0 + A_1 \sin \left(\frac{2\pi(t + \phi)}{T} \right) \right]$	Settling velocity	v_s (ms^{-1}) is the settling velocity of each particle at time t (hours), c indicates the local concentration of particles (kgm^{-3}), n is the exponent for adjusting concentration-dependent settling velocities, A_0 (Hall, 2017) is the non-cyclic component of the settling velocity (ms^{-1}), A_1 represents the amplitude of a periodic sinusoidal variation in time (ms^{-1}), T is the period of the sinusoidal variation (hours), and ϕ is the phase lag for the sinusoidal variation (hours).	D-WAQ PART	(Deltares, 2021)
2	$v_s = \frac{\vartheta}{2R} d_*^3 (38.1 + 0.93d_*^{12/7})^{-7/8}$	Settling velocity of	ϑ is the kinematic viscosity of water (m^2s^{-1}), d_* represents the dimensionless diameter of	TrackMPD	(Khatmullina & Isachenko, 2017)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
		spherical particles	the particle, and R is the radius of the particle (mm).		
3	$d_* = 2R(g(\rho_p - \rho_w)/\rho_w\vartheta^2)^{1/3}$	Dimensionless diameter of the particle	g is the gravitational acceleration (ms^{-2}), and ρ_p and ρ_w indicate the density of the particle and the suspending medium, respectively, with the same units (gm^{-3}).		(Khatmullina & Isachenko, 2017)
4	$v_s = \frac{\pi}{2} \frac{1}{\vartheta} g \frac{(\rho_p - \rho_w)}{\rho_w} \frac{2RL}{55.238L + 12.6}$	Settling velocity of cylindrical particles	L indicates the length of cylinders (mm).		(Khatmullina & Isachenko, 2017)
5	$v_s = \sqrt{\frac{4}{3} \frac{d}{C_D} \left \frac{\rho_p - \rho_w}{\rho_w} \right g}$	Settling velocity	d is particle diameter (mm), ρ_p (gm^{-3}) and ρ_w (gm^{-3}) indicate the density of the particle and the surrounding medium, respectively, g is the gravitational acceleration (ms^{-2}), and C_D represents drag coefficient.	CaMPSim-3D	(Pilechi et al., 2022)
6	$C_D = \frac{24}{Re_p} \left(\frac{1 - \Psi}{Re_p} + 1 \right)^{0.25} + \frac{24}{Re_p} (0.1806 Re_p^{0.6559}) \Psi^{-Re_p^{0.08}} + \frac{0.4251}{1 + \frac{6880.95}{Re_p} \Psi^{5.05}}$	Drag coefficient	Re_p is particle Reynolds number and Ψ is the particle shape factor (Dioguardi et al., 2018). Equation (6) has been reported as the best fit for modeling drag when interpreting microplastics'		(Pilechi et al., 2022)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
			sinking (Van Melkebeke et al., 2020).		

586 * The Ichthyop model uses the Stokes law in calculating the settling velocity.

587 **2.1.2 Biofouling and Degradation (Physical Transformation Processes)**

588 According to previous studies, microplastic particles are subjected to some natural transformation
589 processes including fouling, degradation, and mechanical breakdown during their residence time
590 in an aquatic ecosystem that may lead to changes in their physical properties (Van Utenhove,
591 2019). Biofouling is controlled by different ambient environmental factors such as temperature,
592 salinity, and the availability of sunlight, which is necessary for biofilm (including fine and ultrafine
593 organisms and plankton) growth (Collins & Hermes, 2019). On the other hand, particles become
594 more brittle and break apart into smaller pieces as a result of mechanical degradations by
595 weathering, wave actions, interaction with other particles, the seabed, and the shore (Andrady,
596 2011). Degradation is typically a slow process that takes over 50 years to be fully completed for
597 plastics since they are highly resistant to weathering (Andrady, 2011; Barnes et al., 2009).
598 However, in the swash zone (Efimova et al., 2018) and some energetic coastal areas (e.g., salt
599 marshes), natural abrasives including sediment and rock are present, and the degradation may be
600 completed sooner, even within 8 weeks based on (Weinstein et al., 2016), thereby affecting
601 seasonal microplastics' transport. The main equations employed in the reviewed models in this
602 study for considering biofouling and degradation processes are provided in Table 3. It should be
603 noted that none of the reviewed models are capable of simulating different types of aggregation.

604 **Table 3.** The main equations used in the reviewed models for the consideration of biofouling and
 605 degradation.

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
7	$\rho_p = \rho_0 \frac{R_0^3}{(R_0 + BT)^3} + \rho_D \left[1 - \frac{R_0^3}{(R_0 + BT)^3} \right]$	Biofouling of spherical particles	ρ_p (gm^{-3}) is the density of a fouled spherical particle, ρ_0 (gm^{-3}) is the density of a plastic particle, R (mm) indicates microplastic radius, and BT and ρ_D represent the thickness (μm) and density (gm^{-3}) of a biofilm layer, respectively.	TrackMPD–CaMPSim–3D	(Chubarenko et al., 2016)
8	$\rho_p = \rho_0 \frac{R_0^2}{(R_0 + BT)^2} + \rho_D \left[1 - \frac{R_0^2}{(R_0 + BT)^2} \right]$	Biofouling of cylindrical particles	ρ_p (gm^{-3}) is the density of a fouled cylindrical particle, ρ_0 (gm^{-3}) is the density of a plastic particle, R (mm) indicates microplastic radius, and BT (μm) and ρ_D (gm^{-3}) represent the thickness and density of a biofilm layer, respectively.	TrackMPD–CaMPSim–3D	(Chubarenko et al., 2016)
9	$BT = BT_0 + BR\Delta t$	Biofouling	BT_0 (μm) is the thickness of initial biofouling, BR is a constant biofilm thickness in stationary biofouling, and Δt indicates the time interval.	TrackMPD–CaMPSim–3D	(Jalón-Rojas et al., 2019)

10	$Size(D \text{ or } L) = Size_0(1 - DR \cdot T/100)$	Degradation	<p>D (mm) and L (mm) are the particle's diameter and length, respectively. $Size_0$ indicates the particle's initial diameter and length.</p> <p>DR as a constant rate shows the percentage of decrease in size of the particle per day (Weinstein et al., 2016), and T represents the duration of degradation from the beginning to the present time step.</p>	TrackMPD– CaMPSim– 3D	(Jalón-Rojas et al., 2019)
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606 **2.1.3 Physical Transport Processes**

607 In addition to the transformation processes (biofouling, degradation, aggregation), numerical fate
608 and transport models can simulate different physical processes driving microplastics' transport in
609 water such as advection, diffusion, windage, resuspension, beaching, and washing-off. The
610 transport of microplastics is generally governed by a generic advection–diffusion reaction, as
611 follows (Pilechi et al., 2022):

$$\frac{\partial C}{\partial t} = \nabla \cdot (CU) + \nabla \cdot (K\nabla CU) + \rho(C) \quad (11)$$

612 where C indicates the concentration at time t , U is the velocity of a surrounding medium, K
613 represents the dispersion coefficient, and $\rho(C)$ shows the reaction function.

614 Given the aforesaid concept, the main equations used in the reviewed models for addressing the
615 main physical processes controlling microplastics' transport in aquatic systems are provided in
616 Table 4.

617

618 **Table 4.** The mathematical or empirical equations employed in the reviewed models for
619 considering the physical transport processes that microplastics undergo in aquatic environments.

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
12	$\frac{d\vec{X}(t)}{dt} = \vec{u}(X, t) + \frac{d\vec{X}'(t)}{dt}$	Trajectory Calculation	\vec{X} is the 3D position vector with components in the horizontal plane, d is the total derivative, \vec{u} is the 3D velocity vector, and \vec{X}' is the fluctuation of the position vector.	D-WAQ PART	(Van Utenhove, 2019)
13	$V_{adv} = V_{flow} + V_{windage}$	Advection velocity	V_{adv} (ms^{-1}) represents the velocity field, which is responsible for advection. V_{flow} (ms^{-1}) is the velocity of flow, and $V_{windage} = C_{wd}(V_w - V_{flow})$. C_{wd} is an empirical wind drag coefficient that is related to particle characteristics and V_w (ms^{-1}) represents the wind velocity at 10 m above sea level.		
14	$K_{h(x,y)} = at^b$	Turbulent horizontal component of the diffusion coefficient	$K_{h(x,y)}$ is the horizontal component of the diffusion coefficient (m^2s^{-1}) at time t (which is defined from $t = 0$), a (which is equal to the	D-WAQ PART	(López et al., 2021)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
			<p>dispersion coefficient for small t), and b ($\in [0,1]$) are calibrating coefficients. It should be noted that the upper and lower limits of $K_{h(x,y)}$ are equal to at and a, respectively, and $K_{h(x,y)}$ increases with time. It has been used in Equation (11).</p>		
15	$D_z = \frac{c_\mu^{1/4} L \sqrt{k}}{\sigma_c}$	Turbulent vertical diffusion coefficient	<p>D_z is the vertical dispersion coefficient, C_μ indicates a calibrating constant for local equilibrium shear layers that is equal to 0.09, approximately. L is the mixing length (m), k represents turbulent kinetic energy, and σ_c is the Prandtl–Schmidt number. It has been used in Equation (11).</p>	D-WAQ PART	(Deltares, 2021)
16	$P = 0.5^{-t/T}$	The probability of being washed-off	<p>P is the probability of a particle being washed off, t is the time step, and T is</p>	D-WAQ PART—TrackMPD	(Liubartseva et al., 2018)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
			the half-life of plastic litter.		
17	$v_{part} = v_{water} + \frac{1}{24}gd^2 \frac{\Delta\rho}{\rho_w} \vartheta^{-1} \ln \left(\frac{2I}{d} + \frac{1}{2} \right)$	Vertical velocity of particle	v_{water} is vertical velocity of water, g is the gravitational force (ms^{-2}), d and I are prolate spheroid axes (indicating the area of a particle), $\Delta\rho = \rho_{part} - \rho_w$, ρ_{part} (gm^{-3}) is particle density and ρ_w (gm^{-3}) is the density of water, and ϑ indicates the kinematic viscosity.	Ichthyop	(Parada et al., 2003)
18	$K_h = \varepsilon^{1/3} l^{4/3}$	Horizontal turbulent diffusion coefficient	ε is the turbulent dissipation rate and l is the cell size. It has been used in Equation (11).	Ichthyop	(Peliz et al., 2007)
19	$\begin{aligned} x_{n+1} &= x_n + U\Delta t_i \\ y_{n+1} &= y_n + V\Delta t_i \\ z_{n+1} &= z_n + W\Delta t_i \end{aligned}$	First-order Euler method for advection	Δt_i is an internal time interval and $U = (U, V, W)$ is the velocity vector. This is a first-order discretization of Equation (11).	TrackMPD	(Jalón-Rojas et al., 2019)
20	$x_{n+1} = x_n + R(2r^{-1}K_h\Delta t_i)^{1/2}$	Calculation of horizontal displacement due to	R indicates a random number with an average of zero and r is its standard deviation that equals 1,		(Jalón-Rojas et al., 2019)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
		diffusion Equation (11)	and K_h is horizontal diffusivity (m^2s^{-1}).		
21	$z_{n+1} = z_n + R(2r^{-1}K_v\Delta t_i)^{1/2}$	Calculation of vertical diffusion term Equation (11)	K_v is vertical diffusivity (m^2s^{-1}).		(Jalón-Rojas et al., 2019)
22	$U = \frac{U_f + U_w \sqrt{\frac{\rho_{air}}{\rho_w} \frac{S_{above}}{S_{below}}}}{1 + \sqrt{\frac{\rho_{air}}{\rho_w} \frac{S_{above}}{S_{below}}}}$	Velocity field as a function of current and wave velocities	U_f is the current velocity, U_w is the wind velocity, ρ_{air} and ρ_w indicate the density of air and water, S_{above} and S_{below} represent the cross-sectional areas of spherical and cylindrical particles in dry and wet conditions, respectively. Note that Equation (17) and (21) use the same principles and their difference is only in some coefficients.	TrackMPD	(Jalón-Rojas et al., 2019)
23	$\frac{S_{above}}{S_{below}} = \frac{2\pi}{(\alpha - \sin\alpha)} - 1$	The ratio of the dry cross-sectional area of the particles to its wet cross-sectional area.	α is defined by Equation (24).		(Jalón-Rojas et al., 2019)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
24	$\alpha = 2\arccos\left(1 - \frac{h}{R}\right)$	Parameter for Equation (23)	h is the Archimedean force, R is the radius of the particle.		(Jalón-Rojas et al., 2019)
25	$\left(\frac{h}{R}\right)^2 \cdot \left(3 - \frac{h}{R}\right) = 4 \frac{\rho}{\rho_w}$	Archimedean force	ρ (gm^{-3}) and ρ_w (gm^{-3}) are the density of the particle and water, respectively.		(Jalón-Rojas et al., 2019)
26	$R_j = R_{jmax} \left(\frac{\tau}{\tau_{crit}} - 1\right)$	Resuspension flux	R_{jmax} is the maximum resuspension constant for microplastics. τ and τ_{crit} are shear stress and critical shear stress, respectively.		(Blom & Aalderink, 1998)
27	$\tau = \rho_w \left(\frac{g^{0.5} v_w}{Chezy}\right)^2$	Actual shear stress	ρ_w (gm^{-3}) is the density of seawater, g (ms^{-2}) is gravitational acceleration, v_w is the mean velocity of the seawater, and $Chezy$ is the Chezy coefficient ($\text{m}^{0.5}\text{s}^{-1}$).	CaMPSim-3D	(Blom & Aalderink, 1998)
28	$v_p = \frac{\sum_{i=1}^6 \frac{1}{\alpha_i} v_i}{\sum_{i=1}^6 \frac{1}{\alpha_i}} + v_s$	Vertical displacement velocity of particles	where v_p is the vertical displacement velocity of particles, v_i is particle velocity at the computational cells, α_i is vertex, and v_s is particles settling/rising velocity (obtained from Equation (5)).	CaMPSim-3D	(Pilechi et al., 2022)

No.	Mathematical or Empirical Relationships	Application	Defining Parameters	Employed in	Ref.
29	$Diff = \frac{\partial K}{\partial X} \Delta t + R \left\{ \sqrt{2K_p \left[X_p(t) + \frac{1}{2} \frac{\partial K}{\partial X} \Delta t \right] \Delta t} \right\}$	Calculation of diffusion term	where K is the diffusion coefficient (obtained from Equation (14)), Δt is the time step, R is a random number, and K_p is the diffusion coefficient.		(Pilechi et al., 2022)

620

621 2.5 Numerical Lagrangian Models

622 In PTM models, a velocity value is allocated to each particle. This is typically done by the
623 interpolation of values from an Eulerian hydrodynamic model onto the position of microplastic
624 particles. The advection velocity typically includes the settling velocity, the velocity of flow
625 provided by the hydrodynamic model, and a diffusion-based velocity that is randomly generated,
626 typically based on a Gaussian distribution (Collins & Hermes, 2019; Peliz et al., 2007). Many PTM
627 models use numerical methods such as the fourth-order Runge–Kutta method to solve equations
628 in time.

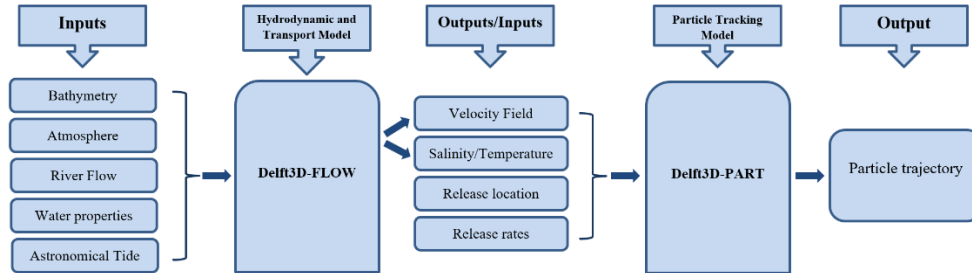
629 In most PTM models, the number of required particles will be determined based on several factors
630 such as the simulation period, the extent of particle dispersion, the cells' size in the computational
631 domain, computational resources, etc. (Deltares, 2021).

632 The majority of PTM models calculate the settling velocity considering the impacts of the
633 aforementioned physical properties such as density, size, and shape. The settling velocity of each
634 particle, v_s , is defined based on the relationships, as shown in Table 2.

635 2.5.1 D-WAQ PART

636 Delft3D, as open-source software, can be applied for simulating flows, the transport of sediments,
637 tidal currents, water quality, waves, etc. in various water settings. Delft3D has several modules as
638 shown in Figure 2. The modules are linked to each other in order to simulate the multi-dimensional
639 transport and hydrodynamics (in 2D or 3D) (Deltares, 2020). Particle transport and dynamical
640 spatial distribution are simulated by the PART module of Delft3D using particle-tracking

641 modeling using the inputs provided by the FLOW module (i.e., two- or three-dimensional flow
 642 data). In other words, velocity and other hydrodynamic parameters are computed by the FLOW
 643 module using an offline coupling and provided as inputs for the D-WAQ PART module.



644
 645 **Figure 2.** Schematic diagram of the Delft3D structure adapted from (Sousa et al., 2021), with
 646 permission from Elsevier, 2021.

647 D-WAQ PART is not capable of considering the impact of the biofouling process on the
 648 trajectories of micro-size particles via increasing the density and size of the particles since it cannot
 649 consider the thickness of the biofilm layer, and so, the density of a fouled particle. D-WAQ PART
 650 also lacks a specific tool to provide insight and appropriately predict complex processes including
 651 degradation as a long-term and aggregation as a short-term process.

652 Particle trajectories are computed over time in three dimensions by modeling the advection and
 653 dispersion processes (Rubinstein, 1981). As stated above, each individual particle can be advected
 654 by water currents and wind or wave drag as additional sources and also displaced randomly due to
 655 diffusion in order to explain unresolved processes.

656 The velocity of particles due to advection includes the local current velocity prepared by the
 657 hydrodynamic model in the Eulerian velocity field and the wind velocity at the location of a
 658 particle as shown in Equation (13) (Van Utenhove, 2019). The wind-drag coefficient is related to
 659 object characteristics as an empirical parameter. In addition, D-WAQ PART employs a random-
 660 walk particle-tracking model in order to describe the movement of particles due to diffusion
 661 processes. The diffusion is described using kinetic gas theory and the theory of Brownian motion
 662 (Deltares, 2021).

663 D-WAQ PART addresses beaching and washing-off using the probability of being washed off, as
 664 shown in Equation (16).

665 Many numerical models, including D-WAQ PART, are not capable of reading the unstructured
666 mesh information provided by an Eulerian model. However, some of the other Delft-3D modules
667 (e.g., D-Water Quality and GPP) can solve the advection–diffusion equation numerically on
668 arbitrary, unstructured grids and structured, shaped grids such as triangles, rectangles, and curved
669 elements.

670 On the other hand, D-WAQ PART utilizes the linear and block interpolation implemented for all
671 time functions between time intervals (Deltares, 2021). In the linear interpolation, D-WAQ PART
672 interpolates the hydrodynamic inputs linearly between the time records while it interpolates the
673 two different subsequent inputs with similar time as the block interpolation.

674 **2.5.2 Ichthyop**

675 Ichthyop is a free Java-based particle tracking numerical model which was originally developed
676 to study ichthyoplankton dynamics (Lett et al., 2008). Ichthyop has also been employed for
677 modeling the fate and transport of microplastics and toxic algae (Collins & Hermes, 2019; Frere
678 et al., 2017). The model can use three-dimensional hydrodynamic information including
679 temperature, velocity, and salinity, simulated by oceanic numerical models such as the Regional
680 Oceanic Modelling System, called ROMS, and the Model for Applications at Regional Scale,
681 named MARS, as input (Lett et al., 2008). Ichthyop is capable of providing two types of
682 microplastics simulations: (1) modeling plastic particle trajectories, and (2) ensemble simulation
683 (i.e., executing a series of numerous particle tracking simulations) (Lett et al., 2008). Ichthyop can
684 also simulate the trajectory of more than one million particles as passive tracers (Frere et al., 2017).

685 Similar to D-WAQ PART, Ichthyop does not consider the integration of complex particle
686 transformation processes including biofouling, different types of aggregation, degradation, and
687 fragmentation.

688 Ichthyop uses the Runge–Kutta method (Lett et al., 2008) for solving the advection equation
689 (Collins & Hermes, 2019). Horizontal and vertical advection and dispersion can be simulated in
690 Ichthyop (Collins & Hermes, 2019; Frere et al., 2017; Lett et al., 2008; López et al., 2021; Pilechi
691 et al., 2022; Previmer, 2007). The model follows D-WAQ PART and uses Equation (11) to
692 simulate the horizontal advection. Ichthyop utilizes a terminal velocity (v_{part}) in order to address
693 the vertical movement of particles, as defined in Equation (17) (Parada et al., 2003).

694 Ichthyop uses Equation (18) for the consideration of horizontal diffusion. Accordingly, horizontal
695 diffusion is defined as a function of the turbulent dissipation rate and the cell size of the
696 computational domain (Peliz et al., 2007). Furthermore, the vertical diffusion of particles is
697 predicted based on a proposed random-walk simulation (Visser, 1997).

698 Ichthyop is incapable of simulating some other effective physical processes on microplastics' fate
699 and transport, such as beaching and washing-off, windage, resuspension, and aggregation, which
700 can be regarded as an important limitation.

701 Hydrodynamic parameters are simulated on a rectangular structured mesh in some hydrodynamic
702 models such as ROMS and MARS. Then, they are interpolated linearly and extracted at time steps
703 in order to be prepared as inputs for Ichthyop.

704 **2.5.3 TrackMPD**

705 Amongst the three-dimensional numerical models, TrackMPD is an open-source model that has
706 been released in recent years (Jalón-Rojas et al., 2019) based on the framework of the Particle
707 Tracking and Analysis Toolbox (PaTATO) (Fredj et al., 2016). TrackMPD is compatible with
708 various ocean models in order to address the fate and transport of marine microplastic debris in
709 coastal areas. TrackMPD is capable of simulating microplastics' behavior in three dimensions
710 considering different physical processes, including beaching and washing-off, biofouling and
711 degradation, advection, windage, turbulent dispersion, and sinking, as well as different physical
712 properties such as density (Equations (7) and (8)), size (Equation (10)), and shape (Equations (2)–
713 (4), (7), and (8)). The physical properties of microplastics can be parametrized in TrackMPD,
714 particularly in the calculation of the settling velocity. TrackMPD utilizes the empirical equations
715 of (Khatmullina & Isachenko, 2017) and (Zhiyao et al., 2008) in estimating the settling velocity
716 of microplastics with different shapes as a function of particle density and radius. Accordingly, the
717 settling velocity (v_s) of spherical particles (ms^{-1}) is obtained based on a calibrated formulation
718 provided by (Zhiyao et al., 2008), as shown in Equations (2) and (3), while this is different, as
719 determined in Equation (4), for cylindrical fragments (mms^{-1}) (Khatmullina & Isachenko, 2017).

720 The effects of biofouling on the fate and transport of microplastics can be investigated by
721 TrackMPD. In other words, TrackMPD uses some empirical equations in order to approximate the
722 thickness of the biofilm layer. In this respect, the thickness of biofilm layers of spherical and

723 cylindrical particles can be estimated based on Equations (7) and (8), respectively (Chubarenko et
724 al., 2016). Accordingly, biofilm thickness can be given either a constant or variable value over
725 time. The case of constant biofilm thickness is referred to as stationary biofouling, while the case
726 of variable biofilm thickness is considered as non-stationary biofouling.

727 TrackMPD simply predicts mechanical degradation by using an empirical relation (Equation (10)).
728 In this respect, the size of plastic particles (their diameter or length) decreases based on a constant
729 rate and time. While TrackMPD can simulate biofouling and degradation, it suffers from
730 incapability in simulating some other transformational processes such as hetero- and homo-
731 aggregation.

732 To simulate the horizontal and vertical movements of plastic particles due to advection, TrackMPD
733 employs the Runge–Kutta method (Jalón-Rojas et al., 2019). In this respect, the velocity of the
734 flow is calculated iteratively in three directions at the location of a particle to consider velocity
735 values (ms^{-1}) at past and future times. Then, to calculate particle displacements (m) and obtain
736 their new location in each direction, the previous location of microplastics is added to the
737 multiplication of the internal time interval (Δt_i) and the flow velocity's components ($U =$
738 (u, v, w)), as shown in Equation (19). Moreover, the horizontal (in x or y directions) and vertical
739 diffusion of particles is predicted in TrackMPD based on a random-walk model, as demonstrated
740 in Equation (20) and (21), respectively (Jalón-Rojas et al., 2019).

741 As with D-WAQ PART, TrackMPD can consider the windage effect in the calculation of particle
742 velocity. Accordingly, the horizontal wind-induced drift of microplastics is considered using
743 Equations (22)–(25) (Jalón-Rojas et al., 2019).

744 TrackMPD estimates the probability of microplastics' washing-off based on a Monte Carlo
745 method, as shown in Equation (16) (Jalón-Rojas et al., 2019). The probability (P) of being washed
746 off in TrackMPD is an exponential function of the time step (t), and the half-life of plastic litter
747 (T) is equal to the time that microplastics remain on the beach (Liubartseva et al., 2018). If P is
748 higher than a random number, which varies from 0 to 1, beached microplastics can be transported
749 (i.e., washed off) seaward at extreme tides. D-WAQ PART uses a basically similar approach in
750 the calculation of the probability of wash-off.

751 Currently, TrackMPD can read hydrodynamic parameters off-line from some models such as
752 ROMS in only a rectangular structured mesh. Moreover, hydrological properties are interpolated
753 spatially at each time step, and then the hydrodynamic model assigns the values to particles located
754 in the corresponding grid cells of TrackMPD (Jalón-Rojas et al., 2019).

755 **2.5.4 CaMPSim-3D**

756 Canadian microplastic simulation (CaMPSim), the newest three-dimensional Eulerian–Lagrangian
757 framework (Pilechi et al., 2022), has been developed recently for simulating microplastics’ fate
758 and transport in different aquatic environments. CaMPSim-3D is based on the coupling of a three-
759 dimensional Lagrangian particle-tracking model and TELEMAC modeling system as the Eulerian
760 hydrodynamic numerical model. Similar to procedures followed by the aforesaid models,
761 hydrodynamic inputs such as atmospheric pressure, the direction and speed of the wind,
762 hydrodynamic and water quality parameters (including velocity, temperature, salinity, water level,
763 and algae concentration), wave height and direction as well as its time period, and characteristics
764 of suspended sediments that are all prepared by the Eulerian model are fed into the Lagrangian
765 particle-tracking model. On the other hand, CaMPSim-3D considers the physical properties of
766 microplastics (such as density, size, and shape) and physical processes (including the effect of
767 positive/negative buoyancy, advection, biofouling, spatio-temporal varying diffusion, degradation,
768 washing-off, and beaching, etc.).

769 In CaMPSim-3D, the physical properties of plastic particles such as density, diameter, gravity, and
770 shape are parameterized in the calculation of the rising or settling velocity. Unlike TrackMPD, the
771 formulation employed in CaMPSim-3D for calculating the settling velocity can be used for
772 particles with different shapes, such as spherules, films, fragments, and fibers (Pilechi et al., 2022).
773 CaMPSim-3D calculates the settling velocity (v_s) based on force balance formulation Equations
774 (5) and (6) provided for Newtonian fluids (Pilechi et al., 2022). CaMPSim-3D can predict the
775 formation of biofilm layers on plastic particles. In this respect, CaMPSim-3D utilizes some
776 proposed empirical equations for estimating the biofilm thickness, e.g., Equations (7)–(9). In
777 addition, CaMPSim-3D considers the degradation of microplastics based on empirical
778 relationships provided in the literature, e.g., Equation (10). However, it cannot simulate both types
779 of aggregation (i.e., hetero- and homo-aggregation) that microplastics undergo.

780 As with some of the aforementioned models, the Lagrangian particle-tracking model, CaMPSim-
781 3D, can simulate not only the advection and diffusion of particles but also beaching and washing-
782 off. It is also capable of considering the effect of buoyancy on particle movement (Pilechi et al.,
783 2022). In CaMPSim-3D, the displacement of particles due to horizontal and vertical advection is
784 calculated by solving the first- to the fourth-order of Equations (19)–(21). In addition, the vertical
785 displacement velocity of particles (v_p) is calculated based on the vertical component of flow
786 velocity, and particle settling/rising velocity (v_s , obtained from Equation (5)), as shown in
787 Equation (28) (Pilechi et al., 2022). Furthermore, CaMPSim-3D considers the impact of wind as a
788 driving factor on the transport of microplastics (see Equations (20)–(22). The turbulent diffusion
789 of plastic particles is simulated in CaMPSim-3D by using a modified relationship, as shown in
790 Equation (29) as a function of K (diffusion coefficient, see Equation (14)), Δt (the time step), R (a
791 random number obtained from a normal distribution, which varies between -1 and 1), and K_p
792 (diffusion coefficient at X_p as a function of t) (Pilechi et al., 2022).

793 K_p is calculated using the eddy viscosity (ϑ_t) calculated by the Eulerian hydrodynamic model
794 (Pilechi et al., 2022). Moreover, CaMPSim-3D predicts resuspension and deposition as well as
795 beaching and washing-off using the proposed relationships (see Equation (16)).

796 One of the advantages of CaMPSim-3D, in comparison with the aforesaid models, is its ability to
797 employ an unstructured grid (Pilechi et al., 2022). It should be mentioned that the CaMPSim-3D
798 model reads hydrodynamic parameters from an Eulerian hydrodynamic model (TELEMAC).
799 Although this approach has some benefits by reducing interpolation errors, it imposes more
800 computational time for point locations in an unstructured mesh. Furthermore, CaMPSim-3D
801 employs an inverse distance-weighted method for the interpolation (or map) of hydrodynamic
802 inputs in space, which is implemented for each particle (Pilechi et al., 2022).

803 **2.6 Results and Discussion**

804 Despite the fact that the reviewed models have generally acceptable accuracy in simulating the
805 physical transformational and transport processes, each has some noticeable limitations in this
806 regard. In this section, we have tried to discuss the reviewed models' limitations in detail regarding
807 the consideration of microplastics' physical properties and processes. It can be concluded that this
808 section can give better insights into the current gaps in the reviewed models.

809 D-WAQ PART and Ichthyop assume microplastics as passive particles without inertia when they
810 are floated at the sea surface, which is different in real ocean conditions (Frere et al., 2017; Van
811 Utenhove, 2019) since the shape and geometry of microplastics have impacts on their behavior,
812 particularly during the different physical processes that they undergo. TrackMPD considers plastic
813 particles as both active and passive particles using different scenarios. It means TrackMPD can
814 simulate microplastics with different shapes (spherical and cylindrical) and sizes (Jalón-Rojas et
815 al., 2019). CaMPSim-3D is more powerful in this regard in comparison with the aforesaid models.
816 This is because it utilizes the shape factor (Ψ) which capable CaMPSim-3D in considering
817 particles with different sphericity and circularity (Pilechi et al., 2022).

818 As stated in Section 2.5, D-WAQ PART and Ichthyop cannot simulate biofouling and degradation,
819 while both TrackMPD and CaMPSim-3D are capable of simulating these processes using the same
820 empirical relationships. It should be noted that TrackMPD and CaMPSim-3D cannot simulate
821 biofouling and degradation together within the same simulation. On the other hand, there are some
822 other transformational processes in the real ocean condition such as hetero- and homo-aggregation
823 (Domercq et al., 2022) and agglomeration (known as the formation of particle clusters in which
824 weak physical interactions are dominant between particles (Sturm et al., 2021)). In addition,
825 microplastics undergo weathering, photodegradation, chemical and biological degradation, and
826 defouling (known as the biofilm layer disintegration (Kooi et al., 2017)) in real conditions.
827 However, none of the reviewed models can simulate any of the aforementioned processes yet.

828 Unlike Ichthyop, D-WAQ PART, TrackMPD, and CaMPSim-3D can predict beaching and
829 washing-off using a probability empirical theory with the same accuracy. However, actual
830 beaching and washing-off in the real condition are highly complex processes since they depend on
831 many other factors, such as the characteristics of winds, tides, waves, and coastline structure (Van
832 Utenhove, 2019). Moreover, microplastics may undergo stranding at the shoreline and be washed
833 off from it repeatedly in subsequent episodes. Accordingly, the aforesaid models are currently
834 incapable of simulating beaching and washing-off more accurately.

835 All of the reviewed models can predict advection and diffusion; however, their accuracy in this
836 regard is different. All models can simulate advection and windage (except for Ichthyop).
837 Moreover, all of them are capable of modeling horizontal and vertical diffusion, although some of
838 them employ different methods and relationships in calculating the horizontal/vertical diffusion

839 coefficient. However, CaMPSim-3D utilizes a modified relationship for the calculation of
840 turbulent diffusion, which seems to be a more accurate approach in comparison with the methods
841 employed by other models. Among the reviewed models, only CaMPSim-3D can predict
842 resuspension. On the other hand, there are some additional important driving transport processes
843 such as Stokes drift (as a result of wave action (De Leo et al., 2021; Van Utenhove, 2019)), burial
844 into the deep sediment of the waterbodies (Domercq et al., 2022), and remobilization (Liro et al.,
845 2020), which none of the reviewed models can predict.

846 All of the reviewed models (except for CaMPSim-3D) are incapable of reading hydrodynamic
847 parameters from an unstructured mesh provided by an Eulerian model. It should be noted that
848 although reading hydrodynamic data from an unstructured mesh leads to a decrease in interpolation
849 errors, it increases the computational costs of particle-tracking simulations.

850 D-WAQ PART and TrackMPD can read and then simulate spatio-temporally interpolated
851 hydrodynamic parameters, while Ichthyop is only capable of utilizing temporally interpolated data.
852 CaMPSim-3D can consider only spatially interpolated hydrodynamic values in the particle-
853 tracking of plastic particles. It seems D-WAQ PART and TrackMPD utilize more accurate
854 numerical schemes in this regard. On the other hand, the efficiency and computational speed of
855 each of the reviewed models can be compared in systematic case studies. It should be noted that
856 there are other numerical models such as TRACMASS and OceanParcels which have not been
857 reviewed in this paper for the sake of brevity and to focus on models with more capabilities.

858 To make a general comparison among the reviewed models, the main characteristics and
859 capabilities of the reviewed models have been summarized in Table 5. It can be seen which of the
860 reviewed models are free for non-commercial use. Moreover, the capability of the reviewed
861 models in terms of the particle tracking of what type and shape of particles have been compared.
862 In addition, the main physical transformational and transport process each model can simulate are
863 shown.

864

865

866

Table 5. Summary of the main information of the reviewed models.

Particle-Tracking Model	Open Access (Free to Use)	Application (Particle Tracking of)	Shape and Type of the Simulated Particles	Main Mechanism Can be Simulated	Initial Year of Release *	Ref.
D-WAQ PART	No	Microplastics, Sediments (first-order decaying substances)	Different shapes and types	Advection (and windage), Diffusion, Beaching and washing-off, Sinking	2019	(Alosairi et al., 2020; Sousa et al., 2021)
Ichthyop	Yes	Microplastics	Prolate spheroids	Advection, Diffusion, Sinking	2008	(Collins & Hermes, 2019; Lett et al., 2008)
TrackMPD	Yes	Microplastics	Spherical and cylindrical particles	Advection (and windage), Diffusion, Beaching and washing-off, Sinking, Biofouling, Degradation	2019	(Jalón-Rojas et al., 2019)
CaMPSim-3D	Yes	Microplastics	Spherules, films, fragments, and fibers	Advection (and windage), Diffusion, Beaching and washing-off, Sinking,	2021	(Pilechi et al., 2022)

				Biofouling, Degradation		
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868 * Based on the publication date of the paper which introduces the model.

869 **2.7 Conclusions and Future Research Directions**

870 The current study aims, as a worthwhile endeavor, to systematically review some well-known
871 Lagrangian numerical particle-tracking models that have been used for predicting the fate and
872 transport of microplastics. In this regard, four Lagrangian particle-tracking models, i.e., D-WAQ
873 PART, Ichthyop, TrackMPD, and CaMPSim-3D, have been investigated in terms of addressing
874 physical transport processes (advection, diffusion, windage, beaching and washing-off, and
875 resuspension) and physical transformation processes (biofouling and degradation) with
876 consideration of the effects of the physical properties (mainly size, diameter, and shape) of
877 microplastics on their fate and transport. Moreover, a numerical scheme of each model is described
878 based on its spatio-temporal interpolation scheme and the structure of the adopted mesh.
879 Accordingly, it has been shown that some of these models are more powerful than others in
880 capturing the physical properties of particles and predicting their behavior under different physical
881 transport and transformation processes. Some concluding remarks for the current study are
882 summarized below:

- 883 • With respect to the consideration of the physical properties of microplastics, D-WAQ
884 PART and Ichthyop use simplified equations for calculating the settling velocity as a
885 function of the size and density of plastic particles, while TrackMPD and CaMPSim-3D
886 employ more accurate equations to meet the effect of particle shape as well. In other words,
887 D-WAQ PART utilizes an equation for calculating the settling velocity of plastic particles
888 with different shapes and types, while Ichthyop is suited for considering the settling
889 velocity of prolate spheroids. TrackMPD can differentiate spherical and cylindrical
890 particles in the simulation of the settling velocity. CaMPSim-3D can be applied for
891 simulating particles with different shapes, including spherules, films, fragments, and fibers.
- 892 • Among the transformation processes reviewed in this study (i.e., biofouling and
893 degradation), although biofouling has been regarded as one of the important processes in
894 the fate of microplastics, D-WAQ PART and Ichthyop are unable to predict the behavior
895 of particles under the influence of this process. However, TrackMPD and CaMPSim-3D
896 can simulate fouled particles using empirical equations. Unlike D-WAQ PART and

897 Ichthyop, TrackMPD, and CaMPSim-3D can predict the degradation of microplastics
898 based on the proposed relationship in the literature.

- 899 • Four particle-tracking models employ generally universal advection–diffusion equations,
900 while the equations for parameterizing various physical processes may be different.
901 However, only D-WAQ PART, TrackMPD, and CaMPSim-3D consider the effect of wind-
902 induced drift as an additional term of advection. In addition, all models are capable of
903 meeting horizontal and vertical diffusion (dispersion). All reviewed models use spatio-
904 temporally varying diffusion coefficients. D-WAQ PART, TrackMPD, and CaMPSim-3D
905 can predict beaching and washing-off using some probability-based relationships.
- 906 • D-WAQ PART, Ichthyop, and TrackMPD can read the extracted hydrodynamic data and
907 simulate particle behavior and particle trajectories only on a structured mesh, while
908 CaMPSim-3D reads hydrodynamic data from an unstructured mesh system. Moreover, the
909 hydrodynamic data are only interpolated linearly at each time step and prepared as inputs
910 for Ichthyop. D-WAQ PART and TrackMPD utilize the spatio-temporal interpolation of
911 the current characteristics, while CaMPSim-3D can use only spatially interpolated data.

912 Given the aforesaid descriptions, there are some physical transport and transformation processes
913 that some of the studied models are unable to simulate or exhibit less accuracy for in modeling. In
914 fact, in order to understand the impacts of microplastics on the sustainability of water resources
915 comprehensively, with higher levels of accuracy, we need to fully understand their behavior, which
916 can be elucidated by using a versatile and accurate modeling system. None of the reviewed models
917 can simulate homo- and hetero-aggregation or agglomeration. None of the four models can predict
918 photodegradation or chemical and biological degradation as important processes in the fate of
919 plastic particles. Wave-induced drift as an additional advection is observed only in CaMPSim-3D.
920 It is suggested that further studies be conducted in order to expand the current knowledge of
921 microplastics' fate and transport. There are some important processes that need to be further
922 studied with regard to the behavior of microplastics in the real condition (e.g., hetero- and homo-
923 aggregation, agglomeration, weathering, photodegradation, chemical and biological degradation,
924 defouling, resuspension, Stokes drift, and remobilization). More in-depth studies, as well as more
925 accurate numerical models in capturing the microplastics' behavior, along with significant
926 modifications in the numerical simulation structure of the reviewed models, are recommended for
927 the future.

928 **2.8 References**

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3 A Laboratory Dataset on Transport and Deposition of Spherical and Cylindrical Large Microplastics for Validation of Numerical Models

3.1 Abstract

The widespread presence of micro-sized plastic pollution has raised concerns due to their unique physical and toxic properties. Each year, water bodies carry millions of tons of plastic into the ocean. The inherent characteristics (such as size, shape, and density) of microplastics (MPs), along with flow factors like speed, depth, and pressure, significantly influence how MPs are transported and deposited. Therefore, this research aimed to gather experimental data on the transport and deposition of MPs to serve as a benchmark for numerical modeling. To achieve this goal, various test scenarios were set up in a straight channel flume to investigate different flow velocities, channel dimensions, and particle shapes. It was observed that cylindrical particles with the same density and similar size were more likely to become trapped compared to spherical particles. This study represents progress towards validating numerical models concerning the transport and deposition of microplastics.

3.2 Introduction

Oceans have been affected adversely because of the huge transport of plastic waste from lands, which occurs through different conduits, such as rivers, with an estimated annual rate ranging from 0.41 to 4.0 Mt (Corcoran et al., 2019). The presence of macro-(>1 cm) (N. I. B. Hartmann et al., 2015), meso-(5 mm–2 cm) (N. I. B. Hartmann et al., 2015), micro-(<5 mm) (Rocha-Santos & Duarte, 2015), and nano-sized (1–100 nm) (EU, 2019a) plastics has emerged as a significant environmental challenge due to their persistent nature, toxic properties, and detrimental impact on the Earth's hydrosphere (Barnes et al., 2009), aquatic organisms (Koelmans et al., 2016; Thevenon et al., 2014b), wildlife, and human health (Napper & Thompson, 2019; Wright & Kelly, 2017).

MPs exhibit diverse regular shapes (e.g., spherical, cylindrical, beads, and fibers) (Khatmullina & Isachenko, 2017) and irregular shapes (fragments) (Schwarzer et al., 2022). Furthermore, the physical properties of MPs undergo changes due to various processes and ambient environmental factors, such as algae invasion, salinity, UV index, temperature, and Stokes drift generated by wave action (De Leo et al., 2021; Liu et al., 2020; Miranda et al., 2021; Nava & Leoni, 2021). The interaction between microplastics and microalgae, including colonization and hetero-aggregation, alters the buoyancy of particles, affecting their settling and impacting their fate and transport in

1295 surface waters (Nava & Leoni, 2021). On the other hand, microplastics pose a threat to microalgae,
1296 impeding their growth, diminishing chlorophyll levels and photosynthesis, and inducing
1297 morphological changes (Nava & Leoni, 2021). Aging MPs, in contrast to freshly emitted particles,
1298 exhibit smoother edges as a result of mechanical degradation caused by interactions with other
1299 fragments (Doyle et al., 2011).

1300 The behavior and transport of MPs depend on their key physical attributes, including size, density,
1301 and shape (Bigdeli et al., 2022; Kooi et al., 2017). Denser particles settle faster, even when sharing
1302 similar shapes and sizes, and larger particles with the same density experience increased settling
1303 velocities due to the higher ratio of gravitational force to fluid viscous resistance (Jalón-Rojas et
1304 al., 2019). The shape of MPs also influences their movement and settling velocity, with elongated
1305 particles settling faster than smaller circular ones of comparable size (Decuzzi et al., 2005; Mason
1306 et al., 2020).

1307 The prediction of MP transport and deposition considering their different physical properties is of
1308 great importance. In this regard, there are several ways to investigate the behavior of MPs within
1309 different water bodies. Physical modeling, numerical simulations, and field surveys have been the
1310 most popular methods employed in previous studies. Physical and numerical modeling serve as
1311 valuable tools for predicting the fate and transport of plastic debris, offering a potential solution to
1312 mitigate the plastic pollution problem. Numerical fate and transport models can simulate different
1313 physical processes driving MP transport in water, such as advection, diffusion, windage,
1314 resuspension, beaching, and washing-off as well as the transformation (biofouling, degradation,
1315 and aggregation) (Bigdeli et al., 2022). Laboratory experiments have been acknowledged as one
1316 of the principal approaches to obtain reliable outcomes. However, there are limitations to physical
1317 modeling, numerical simulations, and field analyses. Conducting physical modeling is expensive
1318 and may raise questions regarding scaling waterbody geometry and particle transport and
1319 deposition patterns, although it can provide reliable results. Numerical modeling can predict
1320 potential accumulation zones in remote areas where water quality monitoring is not feasible.
1321 Furthermore, numerical models can simulate some key physical processes using fine grid
1322 resolution, given the ever-increasing computational capacity and ongoing development of
1323 numerical techniques. Moreover, these models can accurately simulate different vertical and
1324 horizontal flow patterns encompassing microplastics. In this respect, numerical models can be

1325 employed for predicting pathways, hot spots/accumulation zones, and potential sources of plastic
1326 pollution in different water settings promisingly (particularly in the calculation of mass balance
1327 and the estimation of moving plastic debris versus settled ones). Accordingly, numerical modeling
1328 can be considered a promising and relatively inexpensive tool in the simulation of MP behavior,
1329 but it must be validated using reliable experimental or field datasets. In other words, the lack of
1330 reliable datasets of different MPs in terms of physical properties limits the validation processes of
1331 numerical modeling approaches. It should be noted that such a validated numerical model,
1332 especially when coupled with field and lab data, can be substantially useful in providing accurate
1333 simulations. Field analysis is the most expensive choice when compared to the other two methods,
1334 and its feasibility may vary depending on the location and timing.

1335 From a physical modeling perspective, simplified and generic experiments can be useful for
1336 generating reliable data for the validation of numerical models. In this regard, several studies have
1337 been conducted to experimentally investigate the transport and fate of MPs, which can be helpful
1338 in developing numerical methods for predicting MP behavior. For instance, a research study (Choi
1339 et al., 2022) sought to uncover and describe how the secondary movements of microplastic fibers
1340 impact their movement. Secondary movements refer to the oscillating movements exhibited by a
1341 particle as it descends in calm or still water. Another study (Boos et al., 2021) introduced an
1342 experimental configuration designed for the real-time tracking and measurement of microplastic
1343 quantities at the pore scale (ranging from 1 μm to 10 μm) in three distinct environments: (a) surface
1344 flow, (b) at the interface of the streambed, and (c) within hyporheic sediments within an
1345 experimental flume setting. To investigate the transportation and burial of MPs by turbidity
1346 currents, a previous study (Pohl et al., 2020) introduced microplastic fragments and fibers into
1347 scaled-down turbidity currents during flume experiments. The primary objective of this study was
1348 to assess the presence of MPs in sediment samples collected from turbidity currents and their
1349 resulting deposits. Another study (Kerpen et al., 2020) investigated the driving forces and
1350 processes that control the movement of microplastic particles and their accumulation in the surf
1351 zone. To achieve this goal, systematic physical model tests were conducted using MPs of various
1352 sizes, shapes, and densities. Additional experimental investigations were conducted (Waldschläger
1353 & Schüttrumpf, 2019) to assess the settling and rising velocities of microplastic particles in relation
1354 to their density, diameter, and shape. The objective was to explore whether the equations
1355 commonly used in sediment transport studies could be applied to MPs and, if necessary, develop

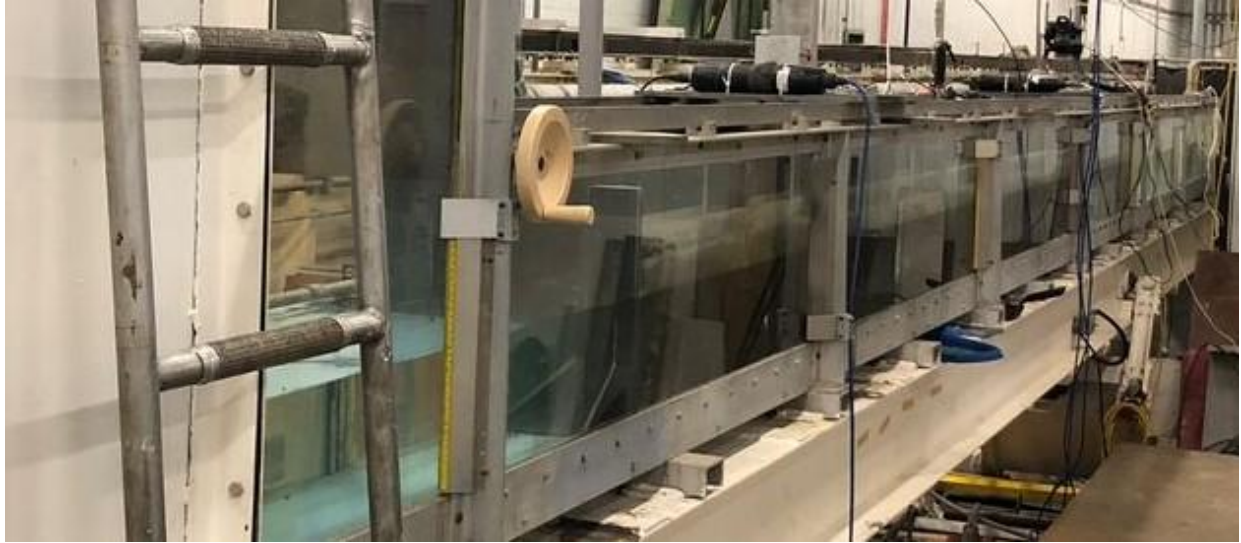
1356 more suitable formulas. To ensure robust data and accuracy, the settling or rising rate of each
1357 particle was measured three times in separate experiments.

1358 Given the aforementioned descriptions, it is crucial to thoroughly characterize and quantify the
1359 deposition of plastic debris in bed materials, especially when developing numerical models to
1360 simulate the fate and transport of plastic debris in water bodies. This information is vital for tasks
1361 such as mass balance calculations and the estimation of the transport of plastic debris versus their
1362 deposition. Accordingly, this study aimed to provide an experimental dataset that can be employed
1363 to validate numerical models utilized in the prediction of MP behavior in water bodies. In other
1364 words, we aimed to bridge the gap between our present knowledge regarding MP transport and
1365 deposition and what happens in real conditions when the velocity of the flow or the geometry of
1366 the channel changes by providing an experimental dataset.

1367 This paper is organized as follows. Section 3.3 describes the experimental setup including details
1368 of the instrumentation, the physical characteristics of particles, the flow conditions, and the
1369 employed geometrical characteristics of the channel. Section 3.4 presents the experimental
1370 outcomes derived from four distinct cases, varying in terms of the flow velocity and particle shape.
1371 Section 3.5 discusses the implications of these observations, and provides concluding remarks
1372 based on the conducted experiments.

1373 **3.3 Experimental Setup**

1374 Physical testing was conducted in a tilting flume at the National Research Council Canada Ocean,
1375 Coastal, and River Engineering Research Centre (NRC-OCRE) physical testing facility in Ottawa
1376 (Figure 3). The model was constructed in a flume measuring 10 m in length, 40 cm in height, and
1377 38 cm in width. It was built using plexiglass and PVC sheets with a thickness of 1 cm. The
1378 schematics of the model, illustrating its geometrical characteristics and dimensions of the different
1379 components, are shown in Figure 4.



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Figure 3. Tilting flume at NRC-OCRE.

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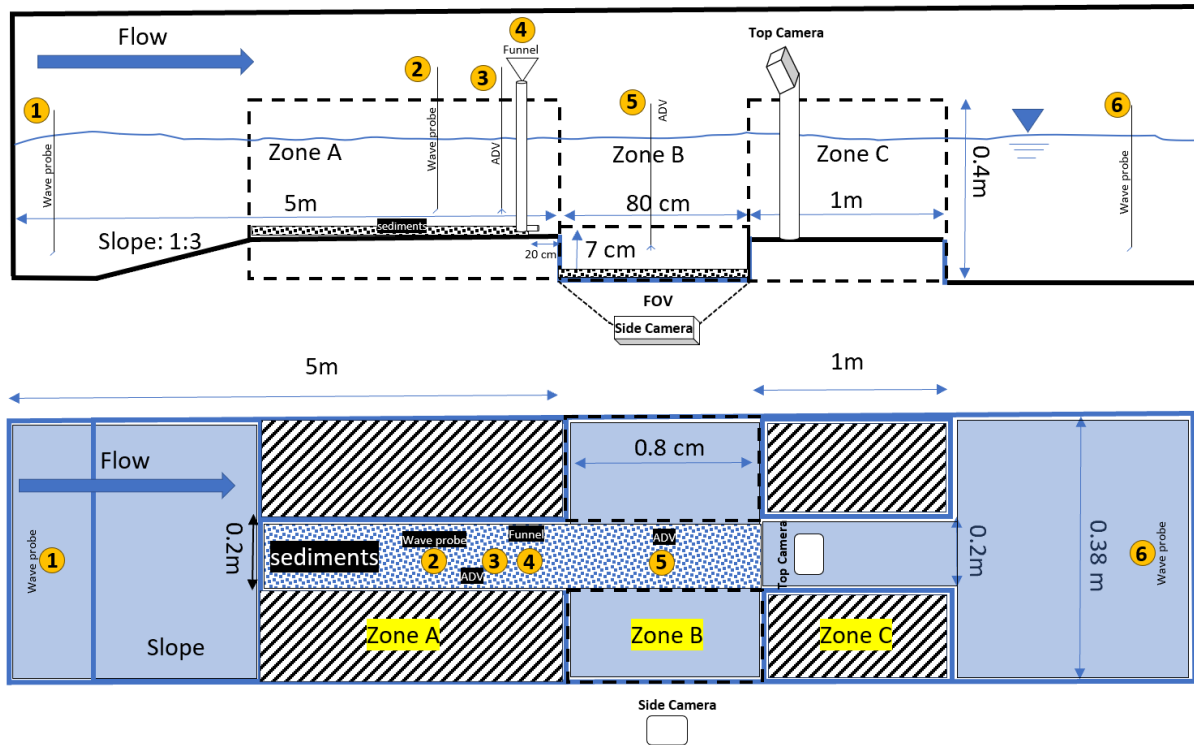
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The water surface elevations in the model were measured using three Nortek wave probes, which were mounted at different locations within the flume: first, at the very beginning of the flume (point 1 in Figure 4); second, in zone A, point 2, very close to the area where large expansion and deepening applied to the channel (i.e., zone B); and third (point 6), at the end of the flume just before the tilting gate (comprised of a panel supported by a hinge along its lower side, allowing water to pass over it when it is lowered) and a screen for capturing particles before entering the basin. Mesh screens were placed before the first wave probe to help maintain a uniform and steady flow within the model. Moreover, the flume inlet started five meters ahead of zone B. Thus, the flow was fully developed before it reached this area.

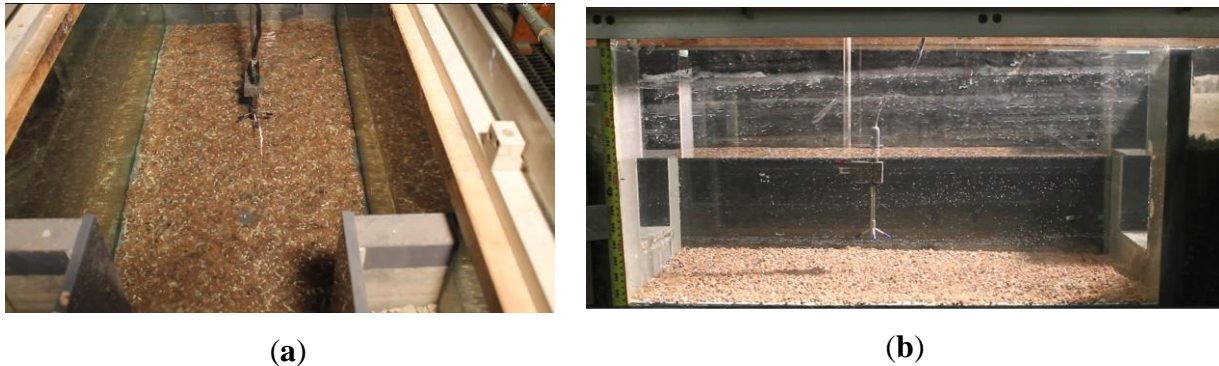


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1392 **Figure 4.** Schematic of the experimental setup: Top figure: Side view, Bottom figure: Top view.

1393 As depicted in Figure 4, a ramp with a slope of 1:3 has been added at the initial section of the
 1394 flume just before the channel that starts in zone A. This addition serves to elevate the bed of the
 1395 channel and enables the creation of a larger depth area in zone B. In order to replicate river
 1396 conditions, the model bed was covered with a layer of sediment (i.e., coarse sand (or pebble)
 1397 (Middleton, 1976) with a mean diameter range of $4 \text{ mm} \leq D_{50} \leq 6 \text{ mm}$) glued to the channel bed
 1398 in zones A and B. The movement of particles within the channel was recorded by two high-
 1399 resolution digital cameras positioned above (in zone C) and on the side (in zone B) of the flume at
 1400 the locations shown in Figure 4. MPs were injected through funnels and injection tubes that were
 1401 only five centimeters above the flume bed in zone A (Figure 4, point 4) and 20 cm away from the
 1402 area of the channel (i.e., zone B) where large and moderate expansion and large deepening
 1403 occurred. Thus, they were injected into the flow from a depth very close to the bed. This allowed
 1404 the particles to be transported and deposited freely on the bed or between the sediments. Water
 1405 velocity profiles were measured before and after the funnel and injection tubes to ensure that there
 1406 was no regime change in the flow's condition in zone A and just before zone B due to the presence

1407 of injection tubes. This setup covers the entire area of zone B. The field of view (FOV) for each
1408 of the mentioned cameras is illustrated in Figure 5.

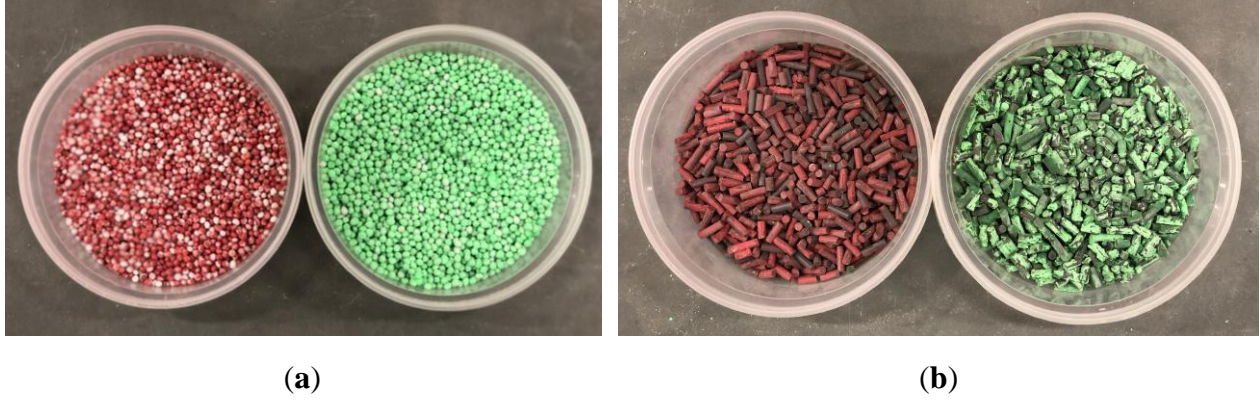


1409 **Figure 5.** FOV in the experimental setup from (a) top and (b) side view.

1410 To circulate the flow within the flume, we employed an 8 psi pressure pump. Water velocities
1411 were measured using two Nortek Acoustic Doppler Velocimeter (ADV), an adaptable and precise
1412 device utilized for measuring three-dimensional water velocity, probes—one deployed (in zone A,
1413 point 3) very close to the funnels and injection tubes (in zone A, point 4) and the other in the
1414 middle of zone B, in point 5 (see Figure 4). The probes were mounted at 0.6 of the total depth from
1415 the water surface. Velocity measurements for each test were recorded at a frequency of 25 Hz for
1416 a duration of 5 min. The selection of this duration and frequency was part of an iterative process
1417 aimed at ensuring that the resolution and duration of the recorded data were sufficient in the sense
1418 that the slope of the spectrum of the measured velocity satisfied Kolmogorov’s $-5/3$ law. The
1419 reported velocity magnitudes in this study are the averages of the recorded data over a 300 s
1420 measurement period.

1421 The time of each test started 100 s before releasing the particles and it lasted five minutes. The
1422 experiments were repeated three times for each test, and all were recorded using a digital camera
1423 to track the particle path, ensuring minimized errors and uncertainties.

1424 The experiments were conducted using fresh water with a density of 1000 kg/m^3 . Two types of
1425 microplastic particles, polyethylene terephthalate (PET) and polyvinyl chloride (PVC), each with
1426 an approximate density of 1380 kg/m^3 , were used for testing. The PET and PVC particles were
1427 shaped as spherical and cylindrical particles, respectively (Figure 6).



1428 **Figure 6.** Particles used in the experiments. (a) Spherical. (b) Cylindrical.

1429 The diameters of both the spherical and cylindrical particles were 3 mm, and the average length of
1430 the cylindrical particles was 4 mm.

1431 The spherical and cylindrical particles were coated both in green and red. They were coated with
1432 waterproof dyes in order to differentiate them along the channel's bed. The density of the coated
1433 particles was checked very carefully to see if the coating changed their density, and we found that
1434 the change in their density was negligible. The green particles were injected into the flow from the
1435 middle funnel while the red particles reached the channel through two side funnels, as shown in
1436 Figure 7. The experiments were performed for two flow velocities, 0.1 m/s and 0.5 m/s, at the
1437 location of the first ADV (i.e., zone A, point 3 in Figure 4). However, because of expansion and
1438 deepening, the velocities in zone B were much less than 0.1 m/s and 0.5 m/s; thus, the particles
1439 had more time to settle and dynamically interact with the near-bed flow.



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Figure 7. Injection tubes across the channel's width.

1442 **3.3.1 Experimental Cases**

1443 Multiple model setups were constructed, incorporating both expansions in the width of the channel
1444 along with increases in the depth of the channel. Four test series were defined for this study based
1445 on two particle shapes (spherical and cylindrical) and two flow conditions (0.1 m/s and 0.5 m/s),
1446 as outlined in Table 6. Each test series involved the examination of two width and depth values for
1447 the expansion region, resulting in a total of 16 tests.

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Table 6. Experimental cases implemented in this study.

Case	Flow Velocity (m/s)	Deepening in Depth (cm)	Expansion in Width (cm)	Material	Shape
Case 1—Spherical MP deposition in channel with a low velocity of flow	Low velocity U = 0.1 m/s	0 and 7	10 and 18	PET	Spherical
Case 2—Spherical MP deposition in channel with a high velocity of flow	High velocity U = 0.5 m/s	0 and 7	10 and 18	PET	Spherical
Case 3—Cylindrical MP deposition in channel with a low velocity of flow	Low velocity U = 0.1 m/s	0 and 7	10 and 18	PVC	Cylindrical
Case 4—Cylindrical MP deposition in channel with a high velocity of flow	High velocity U = 0.5 m/s	0 and 7	10 and 18	PVC	Cylindrical

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1456 On the other hand, the aforesaid cases were classified as shown in Table 7 in terms of large
 1457 deepening or moderate and large expansion in the width of the channel. Both expansion and
 1458 deepening were applied to the channel in zone B. The expansion applied to the width of the channel
 1459 and increased its value from 20 cm to 38 cm within zone B. In addition, deepening made the
 1460 channel bed 7 cm deeper in zone B.

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Table 7. Definition of large and moderate expansion in this study.

Case	Test	Flow Velocity (m/s)	Water Surface Elevation (cm)	Channel Depth at Zone B (cm)	Channel Width at Zone B(cm)	Large/Moderate Expansion/Deepening
Case 1	A	Low velocity U = 0.1 m/s	19.1	0	38	Large expansion
	B		19.8	7	38	Large expansion and deepening
	C		19.8	7	30	Moderate expansion and large deepening
Case 2	A	High velocity U = 0.5 m/s	35	0	38	Large expansion
	B		37	7	38	Large expansion and deepening
	C		37	7	30	Moderate expansion and large deepening
Case 3	A	Low velocity U = 0.1 m/s	19.1	0	38	Large expansion
	B		19.8	7	38	Large expansion and deepening
	C		19.8	7	30	Moderate expansion and large deepening
Case 4	A	High velocity U = 0.5 m/s	35	0	38	Large expansion
	B		37	7	38	Large expansion and deepening

1469 **3.4 Results**

1470 In this section, the observed trap efficiencies for the testing program are provided and discussed.

1471 In this study, 1000 particles were injected into the flow in each single test and also in the repeats.

1472 The number of settled MPs in zone B (as shown in Figure 4) was considered the trap efficiency

1473 (TE) in each test. It is important to highlight that when calculating the TEs, only the number of

1474 particles that both reached and potentially departed from this area was considered. In other words,

1475 particles deposited before this area were excluded from the study. As shown in Table 7, each case

1476 includes different tests (A to C) in terms of the depth and width of zone B, while the velocity of

1477 the flow and the shape of particles remained the same. In addition, each of these tests (i.e., Test A
1478 to C) was repeated three times. Thus, we named the first trial in implementing each test Trial 1.
1479 Trial 2 and Trial 3 indicated the second and third trials in each test, respectively. It should be noted
1480 that outliers (i.e., the unreasonable observations that may have occurred by laboratory errors) were
1481 excluded from the observations in each case when preparing the dataset. For instance, if the
1482 number of observed green MPs that were trapped in zone B was more than the red ones, the test
1483 was considered an outlier.

1484 In this study, the water velocity profiles for two flow conditions (i.e., the average velocity profile
1485 of 0.1 m/s and 0.5 m/s) were measured, as shown in Figure 12, in order to provide a better picture
1486 of the boundary conditions of the implemented tests. Moreover, the Reynolds stresses profiles are
1487 depicted in Figure 13 and Figure 14 (see Appendix).

1488 **3.4.1 Settling Velocity of Spherical and Cylindrical Microplastics**

1489 There are various equations for determining the settling velocity of spherical microplastics (Cheng,
1490 2009; Clift & Gauvin, 1971; Zhiyao et al., 2008), and there are some equations for cylindrical
1491 plastic particles (Khatmullina & Isachenko, 2017). Thus, the measured settling velocities of both
1492 spherical and cylindrical particles in this study were compared to the calculated velocities. The
1493 settling velocity was obtained using a graduated breaker with a height of 25 cm. MPs were injected
1494 into the flow from the water surface and the time that each of the spherical and cylindrical particles
1495 took to reach the breaker bottom was recorded. The settling velocity of the spherical and cylindrical
1496 particles was recorded as 0.14 and 0.18 m/s, respectively.

1497 The settling velocity, w_s (m/s), of the spherical particles was calculated using the Equation (1)
1498 (Zhiyao et al., 2008):

$$w_s = \frac{\nu}{2R} d_*^3 (38.1 + 0.93 d_*^{12/7})^{-7/8} \quad (1)$$

1499 where ν is the water kinematic viscosity (m^2/s), R is the radius of the spherical particle (m), d_*
1500 is the dimensionless diameter of the particle ($d_* = 2R(g(\rho_p - \rho_w)/\rho_w \nu^2)^{1/3}$), g is the gravity
1501 acceleration (m/s^2), and ρ_p and ρ_w are the particle and the water density (with the same units),
1502 respectively.

1503 In addition, some previous studies employed Equation (2) to determine the settling velocity of
 1504 spherical particles (Cheng, 2009; Clift & Gauvin, 1971; Shankar et al., 2021):

$$w_s = \sqrt{\frac{4(\rho_p - \rho_w)gd_n}{3\rho_w C_d}} \quad (2)$$

1505 where d_n is the dimensionless diameter of the volume equivalent sphere ($d_n = \sqrt[3]{6V/\pi}$) and V is
 1506 the volume of the particle. C_d is the drag coefficient and can be calculated using different methods
 1507 (e.g., (Cheng, 2009; Clift & Gauvin, 1971)). In this study, C_d was calculated using Equation (3)
 1508 (Clift & Gauvin, 1971), which covers the turbulent regime ($10^3 \leq Re_p < 2 \times 10^5$) as well:

$$C_d = \frac{24}{Re_p} + \frac{24}{Re_p} (0.15Re_p^{0.687}) + \frac{0.42}{1 + \frac{42,500}{Re_p^{1.16}}} \quad (3)$$

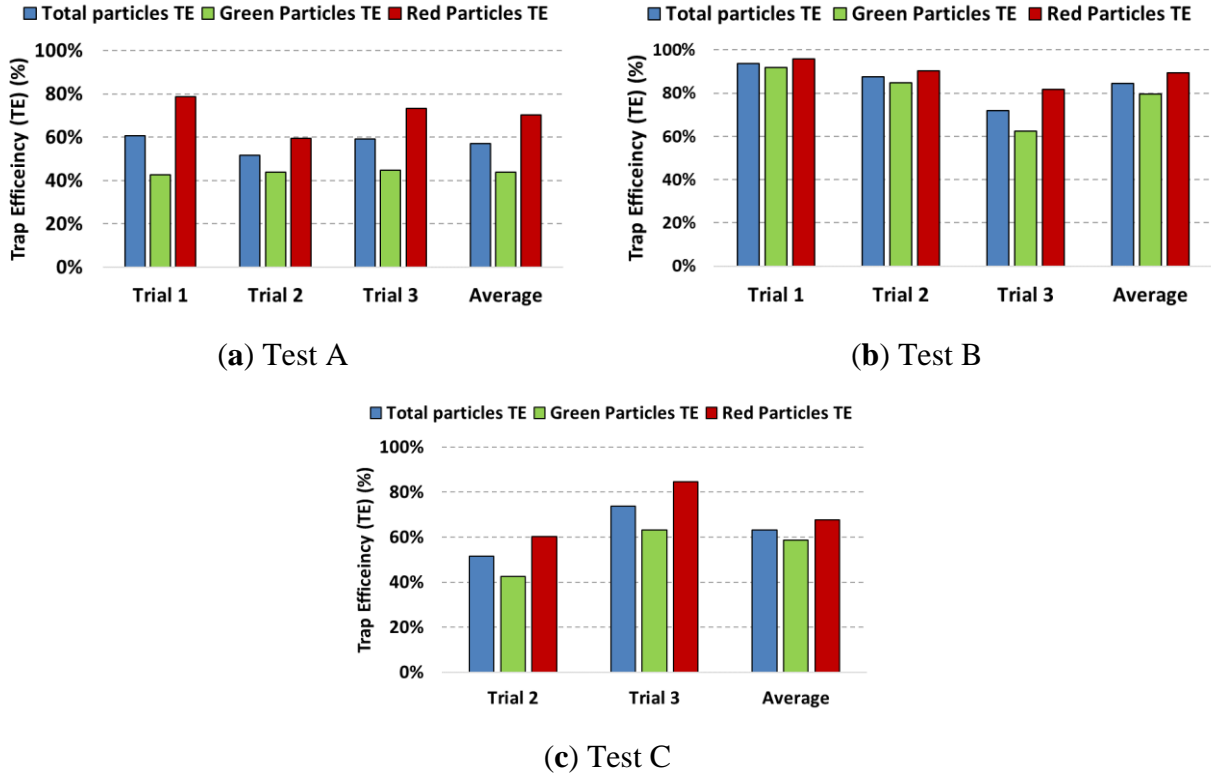
1509 where Re_p is the Reynolds number of the particle ($Re_p = \frac{\rho_w w_s d_n}{\mu}$). In this study, the settling
 1510 velocity was calculated as 0.11 and 0.13 m/s using Equations (1) and (2), respectively.

1511 On the other hand, in this study, the settling velocity of the cylindrical particles was calculated
 1512 based on the equation proposed in the previous study (Khatmullina & Isachenko, 2017), which
 1513 was not close to the measured settling velocity in the lab.

1514 3.4.2 Case 1—Spherical MP Deposition in Channel with a Low Velocity of Flow

1515 In this case, the observations demonstrate the deposition of spherical particles in a flow with a
 1516 mean streamwise velocity of 0.1 m/s, which is considered a low velocity in this study, by varying
 1517 the depth and width of the channel (Figure 8). Notably, across all the tests, the trap efficiency (TE)
 1518 of the red particles (i.e., particles came from the sides of the channel) consistently surpassed that
 1519 of the green particles that were injected into the flow from the middle of the channel. In Figure 8a,
 1520 the average values of TE, in Test A, for the green, red, and total particles are approximately 44,
 1521 70, and 57%, respectively. In Figure 8b, the observations in Test B (with large expansion and
 1522 deepening) indicated that the average TE was notably (over 20%) higher than that in Test A,
 1523 involving only a large expansion (Figure 8a). The average TEs for the green, red, and total particles
 1524 were approximately 80, 89, and 84%, respectively. In Test C, when moderate changes were applied
 1525 to the width of the channel while maintaining large deepening, the resulting TE values for the

1526 green, red, and total particles were considerably lower than that in Test B, at approximately 59,
 1527 68, and 63%, respectively, as indicated in Figure 8c. Trial 1 was excluded in Test C because it was
 1528 identified as an outlier.

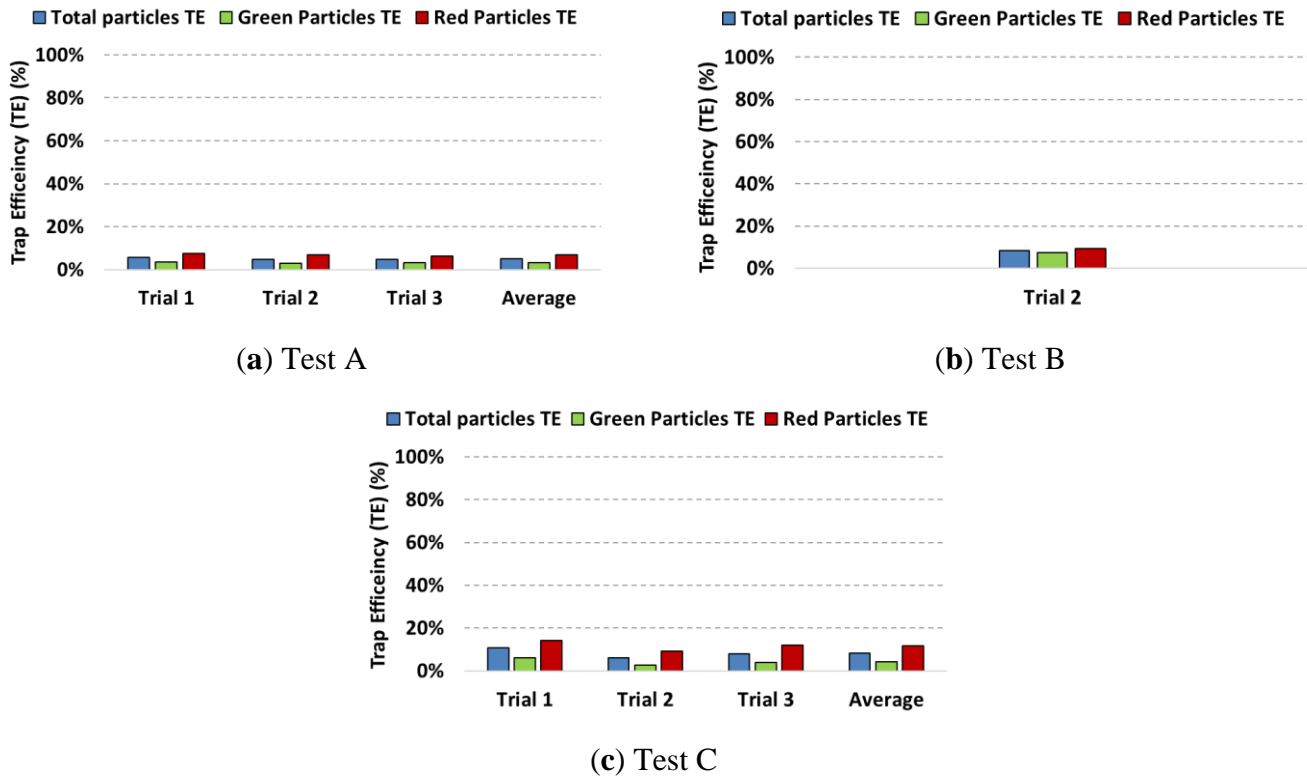


1529 **Figure 8.** Trap Efficiency of MPs in Case 1: Spherical MP deposition in the channel with a low
 1530 velocity of flow ($U = 0.1$ m/s). (a) Test A (large expansion). (b) Test B (large expansion and
 1531 deepening). (c) Test C (moderate expansion and large deepening). The graphs illustrate the results
 1532 of two repetitions of each test along with the average trap efficiency for both green and red
 1533 particles, with the total trap efficiency depicted in blue columns.

1534 3.4.3 Case 2—Spherical MP Deposition in Channel with a High Velocity of Flow

1535 The observations of spherical particle deposition in a flow with a mean streamwise velocity of 0.5
 1536 m/s, considered a high-velocity condition in this study, are shown in Figure 9, considering the
 1537 changes in the depth and width of the channel. Significantly, in all the tests, the TE for the red
 1538 particles injected from the sides of the channel consistently exceeded that of the green particles
 1539 (i.e., particles came into the flow from the middle of the channel) despite the minimal difference
 1540 between the observed values for the green and red particles. The average values of the TE for the

1541 green, red, and total particles were approximately 3, 7, and 5%, respectively, as shown in Figure
 1542 9a for Test A (i.e., where only a large expansion applies to the width of the channel). The
 1543 observations in Test B, where large deepening applies to the channel in addition to large expansion,
 1544 indicate that there is no significant change in the TE values (Figure 9b) compared to that in Test
 1545 A. The average TEs for the green, red, and total particles were approximately 7, 9, and 8%,
 1546 respectively. In Figure 9b, the observations in Trial 1 and Trial 3 were excluded because they
 1547 yielded outlier data. In the case of moderate changes applied only to the width of the channel while
 1548 the deepening in depth was large (i.e., Test C), the TE values for the green (4%), red (12%), and
 1549 total particles (8%) are depicted in Figure 9c.



1550 **Figure 9.** Trap Efficiency of MPs in Case 2: Spherical MP deposition in the channel with a high
 1551 velocity of flow ($U = 0.5$ m/s). (a) Test A (large expansion). (b) Test B (large expansion and
 1552 deepening). (c) Test C (moderate expansion and large deepening). The graphs illustrate the results
 1553 of two repetitions of each test along with the average trap efficiency for both green and red
 1554 particles, with the total trap efficiency depicted in blue columns.

1555 **3.4.4 Case 3—Cylindrical MP Deposition in Channel with a Low Velocity of Flow**

1556 As shown in Figure 10, the laboratory observations illustrate the deposition of cylindrical particles
1557 in the different conditions of the channel (in terms of different depths and widths) where the flow
1558 with a mean streamwise velocity of 0.1 m/s may translocate particles. Among all the tests in this
1559 case, the trap efficiency (TE) for the red particles injected into the flow from the sides of the
1560 channel was higher than that for the green particles injected into the flow from the middle of the
1561 channel. In Test A, the average values of the TE for the green, red, and total particles were
1562 approximately 77, 85, and 82%, respectively (Figure 10a). The application of a large expansion
1563 and deepening on the channel (i.e., Test B) led to a significant increase in the TEs of the green and
1564 red particles. Accordingly, the average TEs for the green, red, and total particles were
1565 approximately 96, 98, and 97%, respectively. While a moderate change was applied only to the
1566 width of the channel (i.e., Test C), compared to the previous test (i.e., Test B), the resulting TE
1567 values for the green, red, and total particles were considerably smaller (i.e., 78%, 84%, and 81%,
1568 respectively), as indicated in Figure 10c. Two replicates of this test (i.e., Trials 2 and 3) were
1569 excluded from the dataset because they were identified as outliers.

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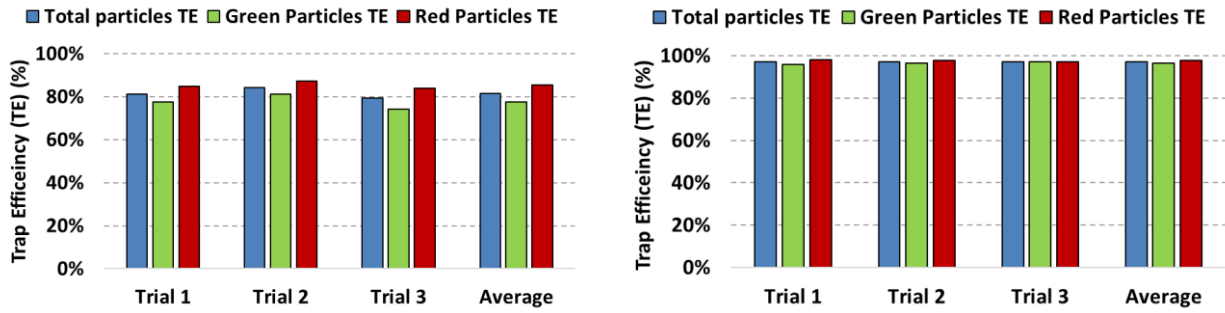
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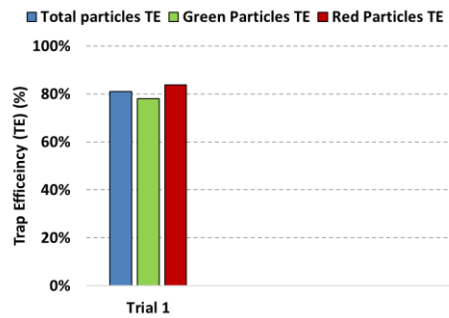
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(a) Test A

(b) Test B



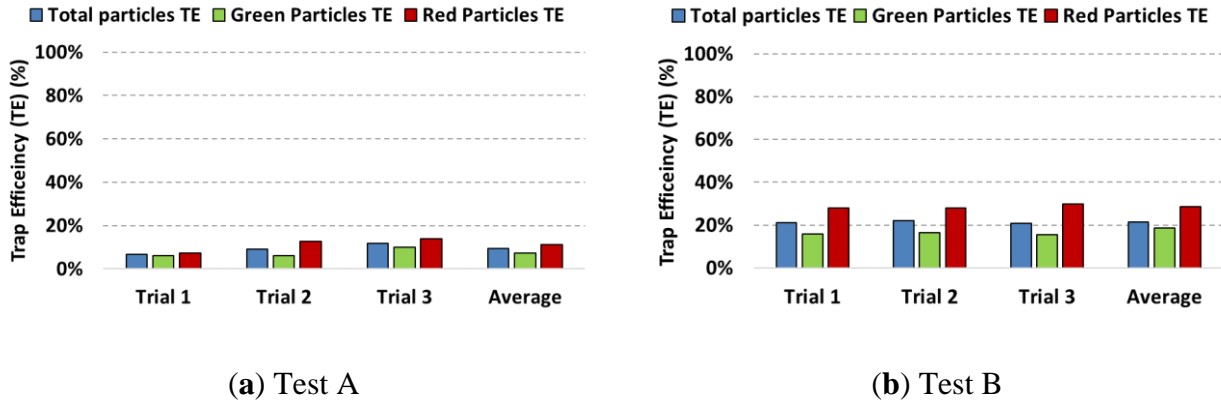
(c) Test C

1577 **Figure 10.** Trap Efficiency of MPs in Case 3: Cylindrical MP deposition in the channel with a low
 1578 velocity of flow ($U = 0.1$ m/s). (a) Test A (large expansion). (b) Test B (large expansion and
 1579 deepening). (c) Test C (moderate expansion and large deepening). The graphs illustrate the results
 1580 of two repetitions of tests along with the average trap efficiency for both green and red particles,
 1581 with the total trap efficiency depicted in blue columns.

1582 3.4.5 Case 4—Cylindrical MP Deposition in Channel with a High Velocity of Flow

1583 In this case, it was observed that some cylindrical particles were deposited on the bed in a flow
 1584 even with a mean streamwise velocity of 0.5 m/s, as shown in Figure 11. As stated above, the red
 1585 and green particles were injected into the flow from the sides and the middle of the channel,
 1586 respectively. In both tests (i.e., Test A and B) of this case, the values of the TE for the red particles
 1587 were consistently higher than those of the green particles, as in the previous cases. When only a
 1588 large expansion was applied to the width of the channel (i.e., Test A), the average values of the TE
 1589 for the green, red, and total particles were approximately 7, 11, and 9%, respectively, as shown in
 1590 Figure 11a. By applying both a large deepening and expansion to the depth and width of the
 1591 channel (i.e., Test B), different TE values were observed for both the green and red particles. In

1592 this regard, the average TEs for the green, red, and total particles were approximately 19, 29, and
 1593 21%, respectively, as shown in (Figure 11b). One of the limitations of this study is the lack of
 1594 observations of cylindrical particles when they undergo moderate expansion in width and a large
 1595 increase in the depth of the channel.



1596 **Figure 11.** Trap Efficiency of MPs in Case 4: Cylindrical MP deposition in the channel with a
 1597 high velocity of flow ($U = 0.5$ m/s). (a) Test A (large expansion). (b) Test B (large expansion and
 1598 deepening). The graphs illustrate the results of two repetitions of tests along with the average trap
 1599 efficiency for both green and red particles, with the total trap efficiency depicted in blue columns.

1600 3.5 Discussion and Conclusions

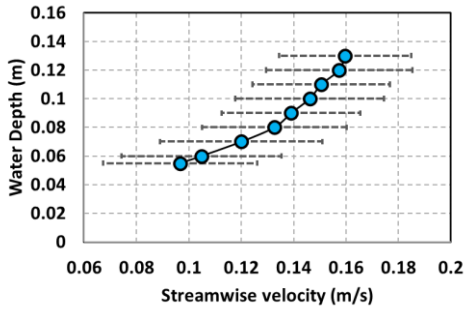
1601 It is crucial to know how different shapes of MPs are transported and deposited in channels. Here,
 1602 a simple and typical geometry was used without any complexity. Numerical models need to be fed
 1603 by experimental or field data in order to yield reliable results. Experimental tests can validate
 1604 numerical simulations so that they can be employed when there is no access to experimental results
 1605 observed in real conditions. The results of this study shed light on the behavior of MPs when they
 1606 undergo different flow velocities. On the other hand, MP transport and deposition were observed
 1607 when bathymetrical changes were applied to the channel's bed. Accordingly, numerical
 1608 simulations can be calibrated and validated based on experimental observations. The applications
 1609 of the generated dataset are the calibration and validation of numerical models. In this regard, the
 1610 initial conditions and boundary conditions in numerical models, particularly CFD models, need to
 1611 be defined properly in order to simulate the real conditions. This is an essential step to determine
 1612 some parameters in the numerical model, including the initial flow velocity and pressure, the flow
 1613 and particle density, the size of the particles, their material, and the static and rolling friction

1614 between the particles and the walls (or boundaries) of the channel. According to the dataset, it is
1615 possible to define the parameters for the numerical models that affect particle transport and
1616 deposition appropriately so that numerical models can be validated and employed for other cases
1617 where there is no access to experimental data.

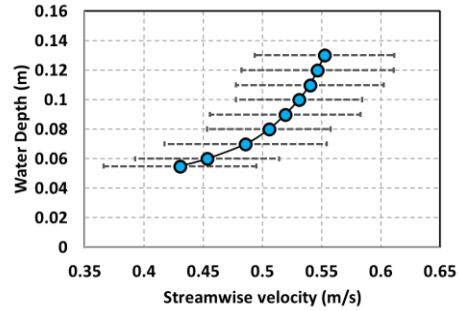
1618 According to the aforementioned observations, it can be seen that by increasing the velocity of the
1619 flow from 0.1 m/s to 0.5 m/s, the values of the TE for both the green and red particles become
1620 smaller among all the cases, as expected. It was observed that fewer green particles injected into
1621 the middle of the channel were trapped compared to the red particles injected into the sides of the
1622 channel because the velocity at the centerline of the channel was higher than that at the sides.
1623 Among the different tests for each case, the values of the TE for both the green and red particles
1624 were the highest when large expansion and deepening were applied simultaneously to the width
1625 and depth of the channel, respectively. By comparing the observed results between the cases
1626 implemented using spherical and cylindrical particles, it was observed that more cylindrical
1627 particles were trapped, which is compatible with previous studies (e.g., (Mason et al., 2020)).
1628 However, it should be noted that since PVC and PET have the same density of 1380 kg/m³, there
1629 is no significant difference in the settling velocity of these materials.

1630 There are some limitations in this study that need to be investigated further in the future. For
1631 instance, this study focused on the investigation of the deposition of microplastics of the same
1632 density and size, overlooking the diverse range of microplastics observed in water bodies, which
1633 vary in terms of both density and size. In addition, there are various shapes of microplastics,
1634 including regular and irregular, found in aquatic environments, whereas spherical and cylindrical
1635 particles were solely investigated in this study. On the other hand, the effect of a slope on the
1636 bottom of the channel was not investigated. In this study, only coarse sand was considered as the
1637 sediment on the bed of the channel. Moreover, a very limited range of Froude numbers (using two
1638 flow velocities and water levels) were included in implementing the tests.

1639 It has been suggested that more in-depth studies are needed in future to obtain a better insight into
1640 MP transport and deposition.

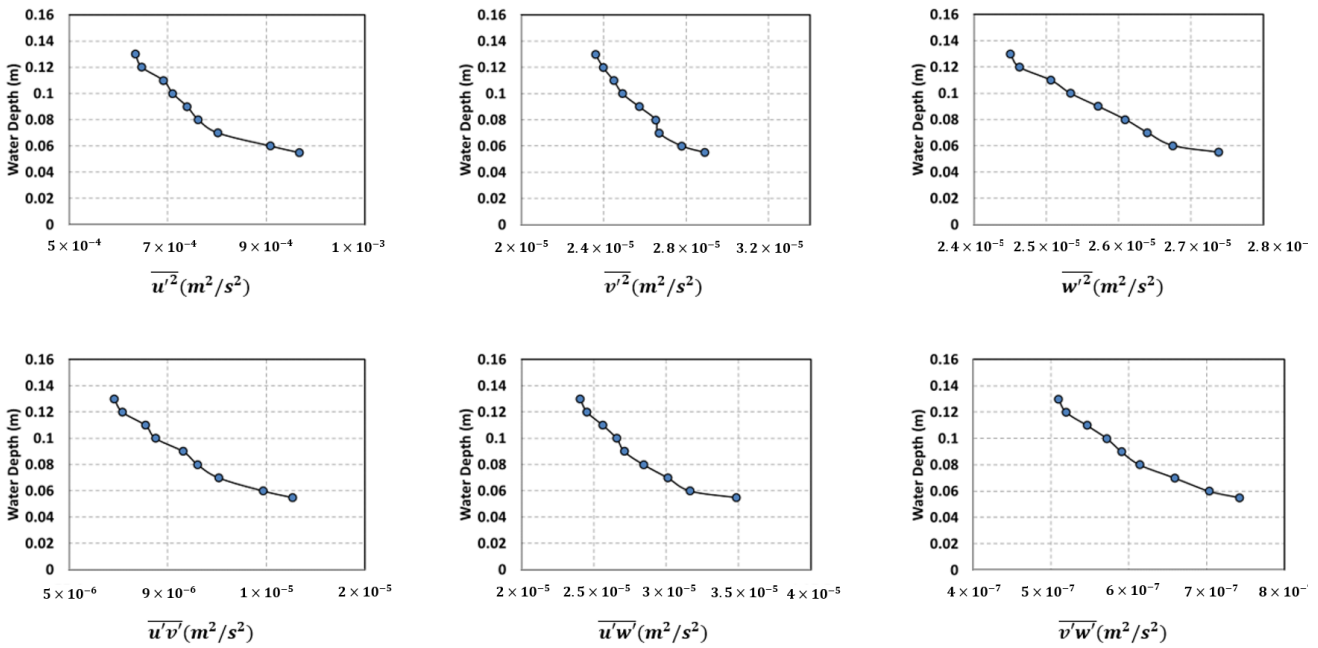


(a)

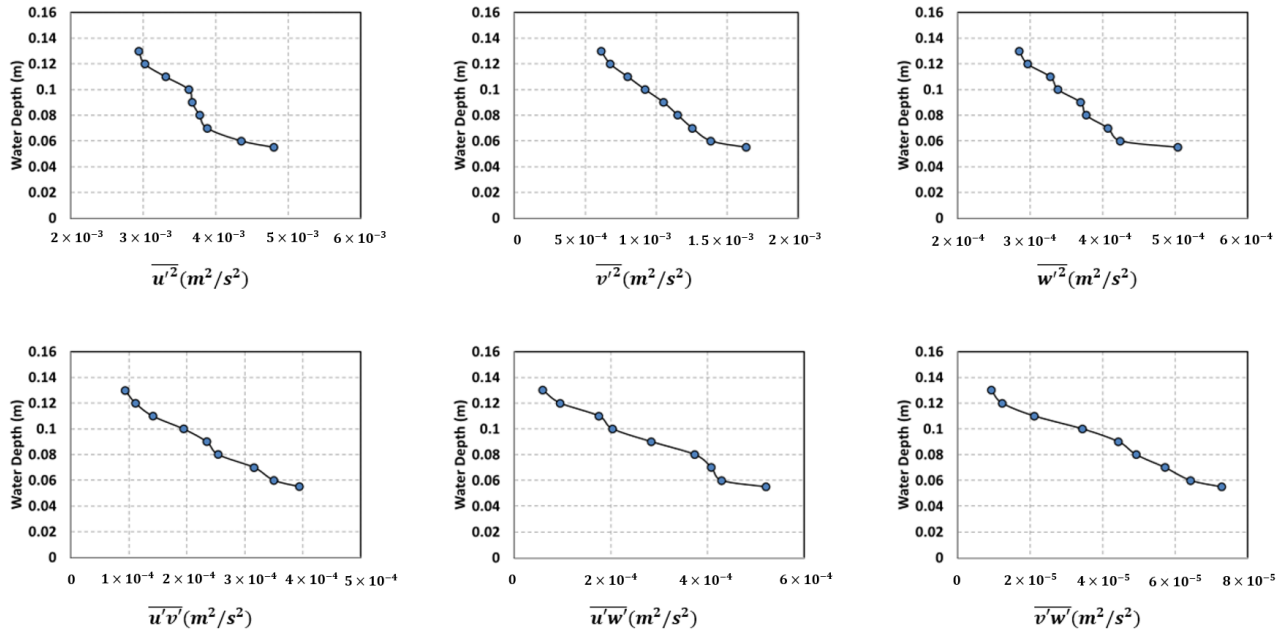


(b)

1642 **Figure 12.** The streamwise water velocity profile in the channel with (a) a low velocity of flow (U
 1643 $= 0.1$ m/s) and (b) a high velocity of flow ($U = 0.5$ m/s). The error bars indicate the standard
 1644 deviation.



1645 **Figure 13.** The profiles of Reynolds stresses in the channel with a low velocity of flow ($U = 0.1$
 1646 m/s).



1647 **Figure 14.** The profiles of Reynolds stresses in the channel with a high velocity of flow ($U = 0.5$
 1648 m/s).

1649 **3.7 References**

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4 Dynamic Prediction of Large Spherical and Cylindrical Microplastic Deposition: A Machine Learning Approach for Transport and Deposition

4.1 Abstract

The prevalence of microplastics (MPs) pollution has cast a shadow over aquatic ecosystems and their inhabitants, including humans. Each year, water bodies transport millions of tons of plastic into the ocean, with a considerable portion of this plastic settling in aquatic environments, leading to complex deposition patterns that impact aquatic ecosystems. This study aimed to develop a novel machine learning model to predict the deposition patterns of spherical and cylindrical microplastics (MPs) with identical particle diameters (d_p) and flow dynamic viscosities (ν), utilizing laboratory-generated datasets. To achieve this, different generic test scenarios were conducted to collect data for various cases of the water depth in the channel (w), the flow velocity (u_f), the water depth in the channel where it undergoes deepening (h), the slope applied to the channel's bed (s), and the particle shape (including spherical and cylindrical). Eleven models including different dimensionless combinations (obtained using Buckingham theorem) of input variables (i.e., $\pi_1 = \frac{w}{d_p}$, $\pi_2 = \frac{h}{d_p}$, $\pi_3 = \frac{u_f d_p}{\nu}$, $\pi_4 = s$) were taken into account. Statistical evaluation metrics were used to determine the best model for predicting either spherical or cylindrical MPs. The model including all four dimensionless inputs was found to be the best model based on either spherical ($R = 0.94$, $MAE = 0.06$, $RMSE = 0.09$, and $BIAS = 0.01$) or cylindrical data ($R = 0.90$, $MAE = 0.09$, $RMSE = 0.11$, and $BIAS = 0.02$). The sensitivity analysis and confidence intervals revealed that the ratio of the water depth in the channel's bed deepening to the particle diameter had the most significant influence on the deposition patterns of both spherical and cylindrical MPs. These findings underscore the potential of machine learning approaches in advancing our understanding and prediction of microplastic behavior in aquatic environments, contributing to improved environmental monitoring and management strategies.

Keywords: Spherical, Cylindrical, Microplastics, Transport, Deposition, Machine Learning

4.2 Introduction

Large MPs (≤ 5 mm) can easily enter rivers, lakes, and coastal areas through direct deposition, wastewater effluent, surface runoff, infiltration, wind, etc., due to their small size and density as well as the huge uncontrolled release (Bigdeli et al., 2024; Bigdeli et al., 2022; Fazil et al., 2024;

1777 Fiore et al., 2022; Gan et al., 2023; Xi et al., 2023). Subsequently, they are deposited or transported
1778 between various water bodies. MPs deposition and transport cause different serious challenges
1779 such as environmental pollution and outbreak diseases. For instance, urban water management
1780 systems, such as water distribution systems and wastewater treatment plants, serve as transition
1781 points for MPs, connecting natural and human water cycles (Ou et al., 2024). On the other hand,
1782 MPs can be absorbed by aquatic fauna, entering the food chain and persisting without degradation
1783 (Alvarez-Zeferino et al., 2020). Consequently, they can infiltrate the body tissues of human organs
1784 such as lymph, liver, kidney, and circulatory systems (Vethaak & Legler, 2021). Furthermore, MPs
1785 can serve as platforms for the growth of harmful biofilms (Vethaak & Legler, 2021).

1786 Several methods have been employed to study MPs behavior in different water bodies, including
1787 laboratory modeling, machine learning (ML) modeling, numerical simulations, and field surveys
1788 (Atugoda et al., 2022; Bigdeli et al., 2024; Bigdeli et al., 2022; Boos et al., 2021; Cao et al., 2024;
1789 Eriksen et al., 2013; Fazil et al., 2024; Issac & Kandasubramanian, 2021; Jalón-Rojas et al., 2019;
1790 Mai et al., 2018; Neder et al., 2022; Shamskhany & Karimpour, 2022). Physical and predictive
1791 modeling (including ML methods), are particularly valuable for predicting the transport and fate
1792 of plastic debris, offering potential solutions to mitigate plastic pollution. Laboratory experiments
1793 are recognized as a primary method to obtain reliable results. In physical modeling, conducting
1794 simplified and generic experiments is useful for (1) observe the behavior of MPs under real-world
1795 conditions and (2) provide data for predictive models using ML approaches (Bigdeli et al., 2024).
1796 However, the main question here is whether ML is capable enough in predicting the behavior of
1797 MPs under the influence of different hydraulic flow conditions and various physical properties of
1798 MPs?

1799 Artificial intelligence (AI) and its subset (i.e., ML) have emerged as promising tools in various
1800 scientific contexts in recent years, particularly in addressing the transport and deposition of MPs.
1801 Previous studies (Fazil et al., 2024; Qian et al., 2024) reported that ML has significant potential in
1802 addressing sediment and MPs deposition and transport, and providing proactive risk assessment.
1803 This leads to significant improvements in real-time monitoring of MPs, as well as in understanding
1804 their deposition patterns and transport trajectories, which are crucial in adopting timely actions to
1805 preserve human well-being and aquatic ecosystems.

1806 Size, density, and shape are important physical properties that influence the fate and transport of
1807 MPs (Bigdeli et al., 2024; Bigdeli et al., 2022). In addition, the flow velocity, the water depth, and
1808 bathymetry changes exerted on the bottom of a water body including deepening and slope, are the
1809 other parameters that influence MPs deposition and transport. The complexity of their independent
1810 and combined effects makes it difficult to develop predictive models using traditional theoretical
1811 analysis and data fitting methods. However, machine learning can address these challenges by
1812 effectively handling regression tasks, inferring patterns from large datasets, and transforming them
1813 into promising models (Fazil et al., 2024; Qian et al., 2024; Xu et al., 2022). Moreover, machine
1814 learning models exploit the feature of interpretability, which can shed light on the significance and
1815 impact of parameters that play key roles in the deposition and transport of MPs.

1816 One of the fast and promising ML methods is Extreme Learning Machine (ELM) which was
1817 proposed by Huang et al. (2006). ELM is a rapid technique that typically offers very high training
1818 speeds, significantly faster than other popular AI-based methods including adaptive neuro-fuzzy
1819 inference systems (ANFIS), support vector machines (SVM), and artificial neural networks (ANN)
1820 (Ebtehaj et al., 2023; Ebtehaj et al., 2018; Milukow et al., 2019). ELM works within the framework
1821 of a single hidden layer multi-layer perceptron, and this is why ELM is efficient and robust (Aredah
1822 et al., 2022). Furthermore, when modeling a nonlinear and complex system with ELM, the user
1823 only needs to adjust one parameter as well as the number of hidden neurons compared to other AI
1824 techniques that have multiple adjustable parameters the affect model performance (Aredah et al.,
1825 2022; Ebtehaj et al., 2023; Ebtehaj et al., 2018; Milukow et al., 2019).

1826 The ELM consists of three layers of an input, a hidden, and an output (Ebtehaj et al., 2018). The
1827 input layer stores the input variables and makes the data available to the ELM network. The main
1828 computations of the ELM occur in the hidden layer that connects the input and output layers
1829 (Aredah et al., 2022). The information processed in the hidden layer is then passed to the output
1830 layer, which contains target parameters and where the results of ELM are compiled. In other words,
1831 the ELM algorithm starts by randomly initializing its weights between the input and hidden layers
1832 and then analytically approximates the output weights (Huang et al., 2006). It then transforms the
1833 optimization of the weight vector that is needed to be assigned between the hidden and output
1834 layers into a straightforward mathematical problem. Unlike traditional ANNs that use gradient-
1835 based optimization, the ELM method employs a least-squares training approach to develop a

1836 single-layer feedforward network. ELM eliminates numerous complexities associated with
1837 gradient-based algorithms, such as local minima and overfitting, as well as the need for iterative
1838 adjustment of learning parameters including learning rate and epochs (Ebtehaj et al., 2018).
1839 Therefore, ELM is remarkably efficient and tends to achieve a global optimum using commonly
1840 employed activation functions. Generalization efficiency and performance are the advantages of
1841 ELM compared to traditional feedforward neural networks (FFNN) algorithms which have been
1842 proven effective on a variety of topics in different fields (Huang et al., 2006).

1843 This study aimed to evaluate the capability of the ELM model in predicting MPs deposition
1844 patterns based on a comprehensive and reliable experimental dataset derived from controlled test
1845 scenarios. The dataset was generated through the development and construction of a scaled
1846 physical model designed to simulate realistic aquatic environments. A series of experimental test
1847 scenarios were conducted, varying key parameters such as flow conditions, particle characteristics,
1848 and channel geometries, to capture a wide range of deposition behaviors. These tests provided a
1849 robust dataset to rigorously assess the accuracy and efficiency of the ELM model in predicting
1850 MPs deposition patterns, highlighting its potential for applications in environmental monitoring
1851 and water resource management.

1852 **4.3 Materials and methods**

1853 **4.3.1 Experimental setup**

1854 In this study, 57 unique test scenarios (including changes in the water depth, the flow velocity, the
1855 elevation and slope of the channel's bed) took place in a tilting flume at the National Research
1856 Council Canada Ocean, Coastal, and River Engineering Research Centre (NRC-OCRE) facility
1857 (Figure 15). The flume is 10 meters long, 40 centimeters high, and 38 centimeters wide including
1858 1-centimeter-thick plexiglass and PVC sheets.



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Figure 15. NRC-OCRE Tilting flume adapted from Bigdeli et al. (2024)

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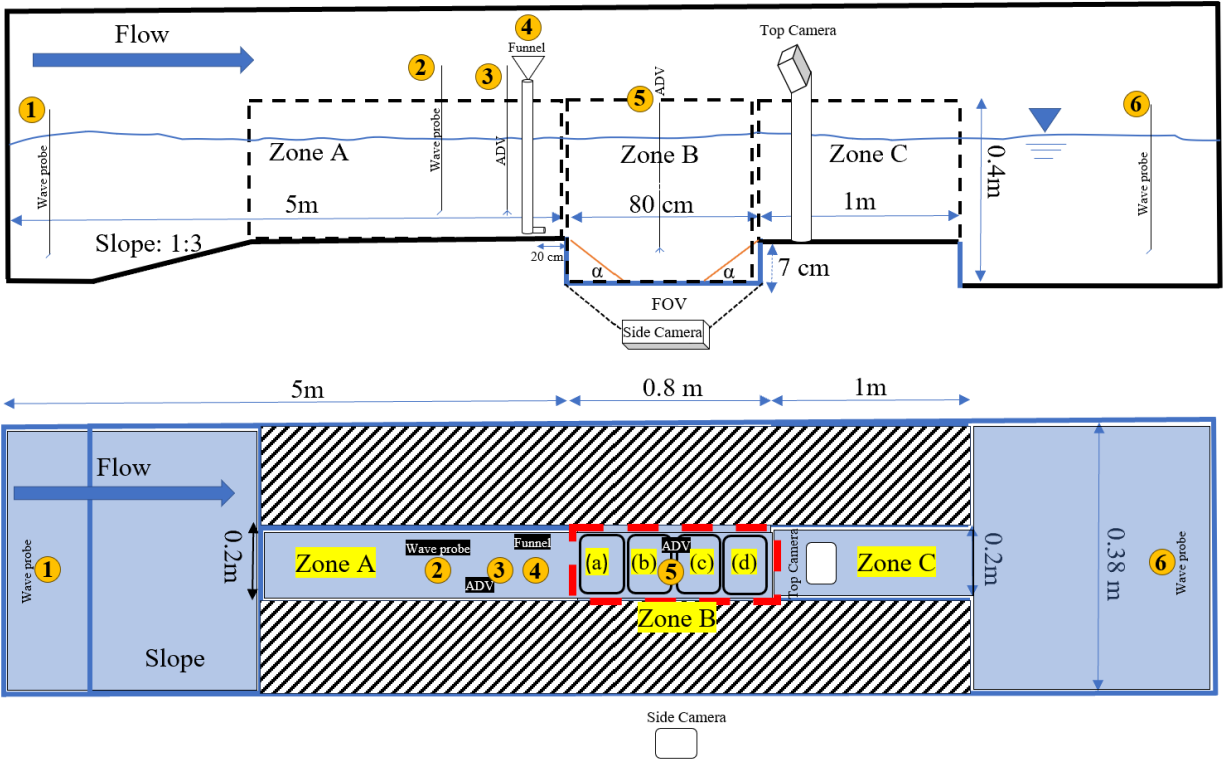
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Figure 16 shows schematically the physical model, detailing its geometric properties and various components as well as the dimensions in each zone. Three Nortek wave probes were employed to measure the water surface elevations in the model. Wave probes were installed at various locations in the flume: the first at the very beginning of the flume (point 1 in Figure 16); the second in zone A at point 2, near the area (i.e., zone B) where the channel deepens sharply (7 cm) and moderately (3.5 cm); and the third (point 6) at the end of the flume, just before the tailing gate. A tailing gate consisting of a panel hinged at the bottom allows water to pass over it when lowered. Additionally, two mesh screens were positioned within the flume: one before the first wave probe to maintain a uniform and steady flow, and another after the tailing gate to capture particles before they enter the basin. It should be noted that each test scenario was started only after ensuring that the flow was fully developed.

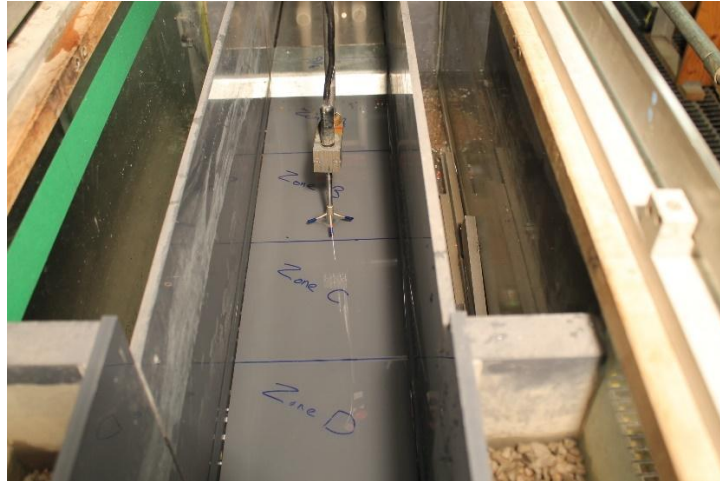


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Figure 16. Schematic of the physical model setup: Top figure: Side view, Bottom figure: Top view.

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As shown in Figure 16, a ramp with a slope of 1:3 was added to the initial section of the flume, just before the start of zone A. This ramp raises the channel bed, allowing a deeper area in zone B. The motion of particles within the field of view (FOV) was monitored using two high-resolution digital cameras placed on the top (in zone C) and side (in zone B) of the flume at the positions indicated in Figure 16. FOV was divided into four zones (a), (b), (c), and (d), each 20 cm in length as shown in Figure 17. Three slopes ($\alpha = 10^\circ, 16.25^\circ, \text{ and } 22.5^\circ$) were applied to the bed of the channel within FOV. MPs were injected into the flow through funnels and injection tubes positioned just five centimeters above the flume bed in zone A (Figure 16, point 4) and 20 cm away from zone B. Therefore, they were injected into the flow from a position very near the bed, enabling the particles to be transported and deposited freely on the bed.



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Figure 17. FOV in the physical model setup from top view

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An 8-psi pressure pump was used to circulate the water flow within the flume. Water velocity profiles were evaluated both before and after the funnels and injection tubes to confirm that the flow conditions in zone A and just before zone B were unaffected by the presence of the injection tubes. In each test scenario, flow velocities were measured using two Nortek Acoustic Doppler Velocimeters (ADV). The first ADV was positioned near the funnels and injection tubes in zone A (point 3), while the other was placed at the center of FOV at point 5 (see Figure 16). The probes were installed at a depth of 0.6 times the total water depth from the surface. Velocity measurements for each test were recorded following the same procedure described by Bigdeli et al. (2024).

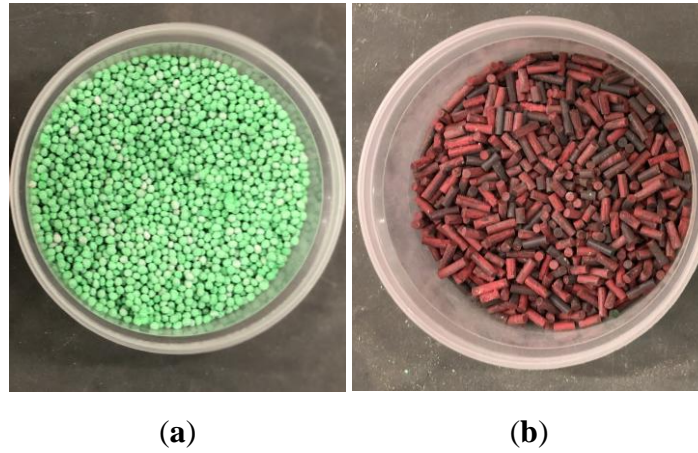
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Each test began 100 seconds before particles release and lasted for five minutes. The test scenarios were repeated three times, and all observations were recorded to investigate MPs deposition and ensure that errors and uncertainties are minimized. It is important to note that the result for each test scenario was confirmed after the investigation of measurements obtained from three repeats and the repeat(s) which included outliers was excluded from the dataset. To measure MPs deposition within FOV, trap efficiency (TE), i.e., the ratio of deposited number MPs to total number of MPs entered FOV, was measured in each zone. The fresh water with a density of 1000 kg/m³ was used for conducting test scenarios. Two widely observed types of MPs in water bodies, polyethylene terephthalate (PET) and polyvinyl chloride (PVC), were used each having a density of 1380 kg/m³, approximately. The shape of PET and PVC particles were considered as spherical and cylindrical, respectively (Figure 18). A total of 1000 MPs particles were injected into the flow

1906

1907 in each test. Both the spherical and cylindrical MPs were 3 mm in diameter. Furthermore, the
1908 average length of the cylindrical particles was 4 mm.



1909 **Figure 18.** MPs used in test scenarios. (a) Spherical (b) Cylindrical adapted from Bigdeli
1910 et al. (2024)

1911 4.3.2 Data Acquisition

1912 Two separate datasets were generated, one for spherical and the other for cylindrical MPs based
1913 on observations in each test scenario. Accordingly, the values of the statistical indices of the
1914 parameters used in this study are shown in Table 8. The effective hydraulic parameters on
1915 deposition of MPs have been given in this table including w (water depth in the channel), u_f (the
1916 velocity of the water flow), h (water depth in the channel where it deepens in FOV), s (the slope
1917 applied to the channel's bed in FOV), d_p (the particle diameter corresponding to 3 mm), ν
1918 (dynamic viscosity of water), TE_a (trap efficiency of MPs in zone a), TE_b (trap efficiency of MPs
1919 in zone b), TE_c (trap efficiency of MPs in zone c), TE_d (trap efficiency of MPs in zone d), and TE_t
1920 (total trap efficiency of MPs in FOV).

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Table 8. Statistical indices of all measured parameters in experimental test scenarios and dimensionless inputs

Parameter	Min		Max		Mean		Standard Deviation (SD)		Variance	
	Sphr.*	Cyln.**	Sphr.	Cyln.	Sphr.	Cyln.	Sphr.	Cyln.	Sphr.	Cyln.
w	0.09	0.10	0.24	0.25	0.16	0.17	0.04	0.04	0.001	0.001
u_f	0.07	0.07	0.36	0.38	0.20	0.20	0.09	0.08	0.01	0.01
h	0.09	0.10	0.31	0.31	0.21	0.21	0.04	0.04	0.002	0.002
s	0	0	22.5	22.5	13.39	12.80	7.77	7.66	60.39	58.69
w/d_p	0	0	1	1	0.49	0.46	0.25	0.24	0.06	0.06
h/d_p	0	0	1	1	0.54	0.53	0.21	0.21	0.04	0.04
$u_f \cdot d_p / \nu$	0	0	1	1	0.49	0.44	0.32	0.27	0.10	0.07
TE_a	0	0	1	1	0.45	0.47	0.46	0.39	0.21	0.15
TE_b	0	0	1	1	0.14	0.15	0.30	0.21	0.09	0.05
TE_c	0	0	0.25	0.42	0.02	0.08	0.04	0.11	0.002	0.011
TE_d	0	0	0.25	0.61	0.03	0.09	0.05	0.15	0.002	0.023
TE_t	0.01	0.02	1	1	0.67	0.77	0.43	0.32	0.18	0.10

1923 *Spherical

1924 **Cylindrical

1925 **4.3.3 ELM Architecture**

1926 ELM, known as a novel alternative algorithm for a single-layer feedforward neural network
1927 (SLFFNN), utilizes the activation function of $f(x)$ as below using N arbitrary training samples,

1928 included as a pair of $(x_i, y_i) \in R^n \times R^m (i = 1, 2, \dots, n)$, with L hidden nodes to predict N different
 1929 samples with an error equal to zero approximately (Milukow et al., 2019):

$$\sum_{i=1}^L \beta_i f_i(x_j) = \sum_{i=1}^L \beta_i f_i(a_i \cdot b_i \cdot x) \quad j = 1, 2, \dots, N \quad (1)$$

1930 where $\beta_i = [\beta_{i1}, \beta_{i2}, \dots, \beta_{im}]^T$ and $a_i = [a_{i1}, a_{i2}, \dots, a_{in}]^T$ denote the matrices of output and input
 1931 weights that are connected both to the hidden layer nodes. b_i represents the bias of the hidden layer
 1932 nodes.

1933 The equation 1 can be rewritten as follows (Milukow et al., 2019):

$$\sum_{i=1}^L \beta_i f_i(x_j) = H\beta \quad (2)$$

1934 where H and β indicate the output matrix of the hidden layer and the output weight matrix,
 1935 respectively, as shown in below (Ebtehaj et al., 2018; Milukow et al., 2019):

$$H = \begin{bmatrix} f(a_1 \cdot b_1 \cdot x_1) & \cdots & f(a_L \cdot b_L \cdot x_1) \\ \vdots & \ddots & \vdots \\ f(a_1 \cdot b_1 \cdot x_N) & \cdots & f(a_L \cdot b_L \cdot x_N) \end{bmatrix}_{N \times L} \quad (3)$$

$$\beta = \begin{bmatrix} \beta_1^T \\ \vdots \\ \beta_L^T \end{bmatrix}_{L \times m} \quad (4)$$

1936 It should be noted that ELM can calculate the output weight with a very small error ($\varepsilon > 0$) using
 1937 the pseudo-inverse of the H output.

1938 In this model, the first step during the training phase is to determine a_i and b_i as the matrices of
 1939 the input weights and the bias of the hidden layer nodes, respectively. Then hidden layer matrix is
 1940 obtained by employing input variables. With the input weights and hidden layer biases established,
 1941 the SLFFNN training process is converted to a least squares solution. The target function is defined
 1942 subsequently as below by utilizing regularization theory in ELM models (Milukow et al., 2019):

$$\min L_{ELM} = \frac{1}{2} \|\beta\|^2 + \frac{c}{2} \|T - H\beta\|^2 \quad (5)$$

1943 where c is a regularization parameter which is used for increasing the generalization capability of
 1944 the ELM (Salimi et al., 2024). T denotes the label matrix as below (Ebtehaj et al., 2018; Milukow
 1945 et al., 2019):

$$T = \begin{bmatrix} y_1^T \\ \vdots \\ y_L^T \end{bmatrix}_{N \times m} \quad (6)$$

1946 ELM can assign random parameters to hidden nodes if the number of discrete samples is more
 1947 than the number of hidden neurons (Ebtehaj et al., 2018).

1948 4.3.4 Methodology

1949 4.3.4.1 Dimensional Analysis

1950 One of the basic theorems in dimensional analysis is called the pi-theorem (Buckingham's π
 1951 theorem), because it applies to groups of products of quantities. This technique allows for the
 1952 development of an equation containing only products of variables, and occasionally, a single non-
 1953 dimensional product, by utilizing a functional transformation. According to dimensional analysis,
 1954 a natural phenomenon can be described by the following very general equation (Gibbins, 2011):

$$f(Q_1, Q_2, \dots, Q_i, \dots, Q_N) = 0 \quad (7)$$

1955 where each of the variables, Q_j , $j = 1, 2, \dots, N$ represents either concepts, properties or units-
 1956 conversion factors.

1957 In dimensional analysis, the π -theorem provides the following transformation of equation 7
 1958 (Gibbins, 2011):

$$f(\Pi_1, \Pi_2, \dots, \Pi_i, \dots, \Pi_N) = 0 \quad (8)$$

1959 Each of Π_i variables form a non-dimensional product of one or more variables in equation 7 where
 1960 $p < N$ and this means that any of Π_1, Π_2, \dots can be related to the others (e.g., $\Pi_1 = \phi(\Pi_2, \Pi_3, \dots)$).
 1961 This transformation theorem is known as the π -theorem. In other words, we can show a
 1962 phenomenon as a function of some dimensionless parameters, taking into account the fundamental
 1963 dimensions. Accordingly, there are n independent variable with m fundamental dimensions (for
 1964 instance time and length). Thus, we can conclude that the phenomenon can be a function of $n - m$
 1965 dimensionless parameters which help us to investigate its trends, variations, and sensitivity in
 1966 appropriate manner. In this study, the dimensional analysis (using Buckingham theorem) was

1967 conducted to reduce scale effects of experimental datasets on the performance of the ELM model.
 1968 In this regard, seven independent variables were considered, as shown below, in order to find
 1969 dimensionless parameters as well as the relationship between the hydraulic characteristics of the
 1970 flow and other effective parameters on MPs transport and deposition:

$$f(v, w, u_f, h, d_p, s) = 0 \quad (9)$$

1971 where v is the kinematic viscosity, w is water depth in the channel, u_f represents the velocity of
 1972 the water flow, h is the water depth in the channel's bed deepening in FOV, d_p is the particles
 1973 diameter, and s is the slope applied to the channel's bed in FOV. By using Buckingham π theorem,
 1974 the dimensionless groups were defined as:

$$TE = f(\pi_1, \pi_2, \pi_3, \pi_4) \quad (10)$$

$$\pi_1 = \frac{w}{d_p}, \pi_2 = \frac{h}{d_p}, \pi_3 = \frac{u_f d_p}{v}, \pi_4 = s$$

1975 Therefore, after generating the dataset for both spherical and cylindrical MPs, the effective
 1976 variables on TE were made dimensionless and connected to it through the Eq. 10. The spherical
 1977 and cylindrical datasets were split separately into two groups of training and testing data in which
 1978 70% of the data were used for training. Then the model optimization parameters were configured,
 1979 and orthogonality was enabled. The number of hidden layer neurons was determined by trial and
 1980 error. It should be noted that for practical applications, the model should be compact, with the
 1981 number of hidden neurons (L) much lower than the number of input samples (N). If N matches L ,
 1982 zero error is achieved in the training dataset and often results in high error in the testing dataset
 1983 which is known as overfitting (Salimi et al., 2024). The methodology of this study includes nine
 1984 steps as depicted in Figure 19.

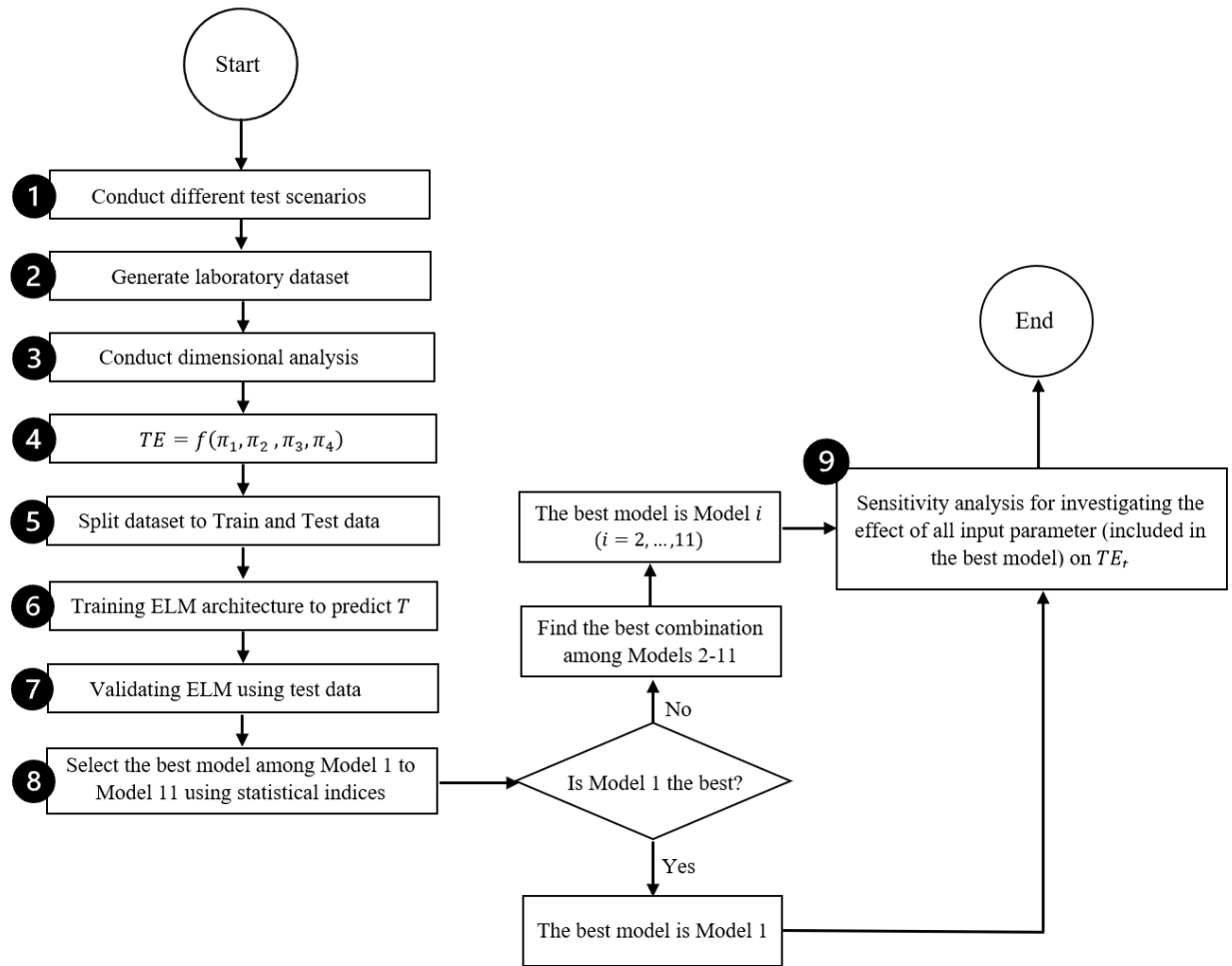
1985 The training data were used to build the model based on 11 combinations, and the performance
 1986 was evaluated using the testing dataset. In order to investigate the effect of each input variable
 1987 included in Eq.10 on the performance and accuracy of the ELM in predicting TE_t , different sets
 1988 of combinations were proposed. First, the ELM was modeled using the combination included in
 1989 Model 1 (i.e., all four dimensionless input variables included in Eq.10). By removing one to three
 1990 input variables, ten more models (Model 2 to 11) were proposed in addition to Model 1, as
 1991 described in Table 9.

1992

Table 9. The input variables included in Models 1 to 11

Model	Included input variables	Excluded input variable(s)
Model 1	$\pi_1, \pi_2, \pi_3, \pi_4$	-
Model 2	π_2, π_3, π_4	π_1
Model 3	π_1, π_3, π_4	π_2
Model 4	π_1, π_2, π_4	π_3
Model 5	π_1, π_2, π_3	π_4
Model 6	π_3, π_4	π_1 and π_2
Model 7	π_2, π_4	π_1 and π_3
Model 8	π_2, π_3	π_1 and π_4
Model 9	π_1, π_4	π_2 and π_3
Model 10	π_1, π_2	π_3 and π_4
Model 11	π_1, π_3	π_2 and π_4

1994 If any of Models 2 through 11, which have one to three fewer parameters than Model 1,
 1995 outperformed Model 1, the better-performing model in predicting TE_t was selected. By applying
 1996 the reduction of input parameters, the increase in modeling accuracy was investigated and the best
 1997 model was selected. This investigation was continued to determine which combination of input
 1998 variables yielded the highest accuracy for ELM, as these had the greatest impact on modeling
 1999 performance. It is important to note that the performance of ELM in prediction of TE_a , TE_b , TE_c ,
 2000 and TE_d was investigated in addition to TE_t as shown in Table 10 and Table 11.



2001

2002

Figure 19. The methodological structure of the study.

Table 10. Statistical characteristics for different models using various combinations of spherical input variables.

Model No.	<i>R</i>					<i>MAE</i>					<i>RMSE</i>					<i>BIAS</i>				
	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>
Model 1	0.96	0.91	0.94	0.93	0.95	0.12	0.09	0.01	0.01	0.09	0.14	0.13	0.02	0.02	0.12	0.03	0.00	0.01	0.00	0.01
Model 2	0.91	0.86	0.75	0.68	0.82	0.15	0.10	0.02	0.03	0.18	0.19	0.15	0.03	0.03	0.24	0.03	0.01	0.01	0.01	0.02
Model 3	0.79	0.90	0.83	0.69	0.81	0.20	0.10	0.02	0.03	0.21	0.28	0.14	0.03	0.03	0.27	0.00	0.01	0.00	0.00	0.04
Model 4	0.76	0.69	0.82	0.85	0.70	0.23	0.15	0.01	0.02	0.24	0.33	0.21	0.02	0.03	0.29	0.14	0.00	0.00	0.00	0.00
Model 5	0.56	0.46	0.89	0.69	0.77	0.31	0.17	0.02	0.02	0.20	0.38	0.25	0.02	0.03	0.27	0.02	0.00	0.00	0.01	0.06
Model 6	0.69	0.76	0.72	0.57	0.65	0.24	0.13	0.02	0.02	0.22	0.34	0.22	0.03	0.03	0.31	0.09	0.00	0.00	0.01	0.08
Model 7	0.77	0.80	0.73	0.55	0.41	0.23	0.17	0.01	0.03	0.33	0.31	0.27	0.02	0.04	0.40	0.07	0.10	0.00	0.01	0.10
Model 8	0.43	0.43	0.44	0.48	0.67	0.38	0.15	0.03	0.03	0.23	0.42	0.20	0.04	0.05	0.31	0.03	0.05	0.01	0.02	0.03
Model 9	0.73	0.77	0.60	0.55	0.49	0.25	0.17	0.02	0.03	0.32	0.32	0.27	0.03	0.04	0.35	0.05	0.11	0.01	0.02	0.09
Model 10	0.33	0.32	0.49	0.62	0.36	0.39	0.19	0.02	0.02	0.33	0.43	0.25	0.02	0.03	0.41	0.07	0.08	0.00	0.01	0.03
Model 11	0.40	0.46	0.46	0.47	0.77	0.35	0.18	0.03	0.03	0.21	0.41	0.25	0.04	0.03	0.30	0.04	0.04	0.01	0.01	0.08

2004

2005

2006

2007

2008

Table 11. Statistical characteristics for different models using various combinations of cylindrical input variables.

Model No.	<i>R</i>					<i>MAE</i>					<i>RMSE</i>					<i>BIAS</i>				
	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>	<i>TE_a</i>	<i>TE_b</i>	<i>TE_c</i>	<i>TE_d</i>	<i>TE_t</i>
Model 1	0.94	0.92	0.91	0.83	0.91	0.11	0.10	0.04	0.07	0.11	0.12	0.13	0.05	0.10	0.14	0.01	0.04	0.01	0.02	0.03
Model 2	0.75	0.90	0.84	0.65	0.82	0.23	0.07	0.05	0.08	0.12	0.28	0.11	0.06	0.12	0.17	0.10	0.00	0.00	0.00	0.03
Model 3	0.79	0.91	0.84	0.71	0.86	0.18	0.08	0.05	0.10	0.14	0.23	0.10	0.07	0.14	0.18	0.05	0.02	0.02	0.03	0.00
Model 4	0.79	0.77	0.82	0.71	0.75	0.18	0.12	0.05	0.06	0.16	0.23	0.16	0.07	0.10	0.21	0.02	0.01	0.01	0.00	0.00
Model 5	0.47	0.49	0.47	0.62	0.82	0.31	0.10	0.08	0.09	0.12	0.37	0.16	0.10	0.13	0.19	0.01	0.01	0.02	0.00	0.00
Model 6	0.63	0.86	0.72	0.69	0.85	0.21	0.10	0.05	0.10	0.14	0.29	0.15	0.07	0.17	0.20	0.04	0.06	0.02	0.05	0.04
Model 7	0.64	0.66	0.77	0.51	0.54	0.24	0.11	0.05	0.09	0.20	0.32	0.15	0.07	0.11	0.24	0.00	0.01	0.01	0.03	0.06
Model 8	0.40	0.53	0.48	0.52	0.67	0.30	0.12	0.08	0.11	0.19	0.35	0.17	0.10	0.16	0.26	0.15	0.04	0.00	0.01	0.03
Model 9	0.73	0.68	0.83	0.58	0.56	0.23	0.11	0.04	0.08	0.24	0.28	0.16	0.06	0.10	0.29	0.04	0.04	0.02	0.04	0.06
Model 10	0.39	0.37	0.52	0.45	0.53	0.31	0.14	0.07	0.09	0.20	0.35	0.21	0.09	0.10	0.27	0.07	0.07	0.02	0.07	0.03
Model 11	0.42	0.57	0.38	0.57	0.73	0.32	0.14	0.07	0.11	0.15	0.36	0.18	0.09	0.15	0.21	0.03	0.01	0.01	0.01	0.02

2012 **4.3.5 Sensitivity analysis**

2013 As shown in step 9 in Figure 19, sensitivity analysis was employed to investigate how changes in
2014 the parameters affect the models' performance. In this study, the partial derivative sensitivity
2015 analysis (PDSA) method was used which is known as one of the commonly used techniques for
2016 analyzing the sensitivity (Azimi et al., 2017; Ebtehaj et al., 2018). In this technique, the partial
2017 derivative of the target parameter (i.e., TE_t) with respect to each input parameter (i.e.,
2018 $\pi_1, \pi_2, \pi_3, \pi_4$) was calculated to assess the sensitivity of ELM to each input variable. The PDSA
2019 value determines the impact of input variables on the model's predictions. In this regard, smaller
2020 or larger PDSA values obtained with respect to each input variables imply lower or higher effect
2021 of that variable on the model's predictions for TE_t , respectively. In addition, positive or negative
2022 PDSA values for each input variable indicate that an increase or decrease in the input variable
2023 causes a corresponding increase or decrease in the target variable (TE_t), respectively.

2024 In this study, the derivative of the obtained equation (Eq. 11) with ELM based the activation
2025 function of Sigmoid, which was in the form of a matrix function, with respect to each effective
2026 parameter was calculated using PDSA technique. This allowed for the investigation of how the
2027 considered function varies with each parameter. The matrix of partial derivatives was employed
2028 in applying PDSA.

$$f(x) = \exp \left(\left[1 / \left(1 + \exp \left(-((InW \times InV) + BHN) \right) \right) \right]^T \times OutW \right) \quad (11)$$

2029 where InW denotes the $p \times q$ matrix of the input weights, InV represents the $q \times 1$ matrix of input
2030 variables (x), BHN indicates the $p \times 1$ matrix of the bias of the hidden neurons, and $OutW$ shows
2031 the $p \times 1$ matrix of the output weights. It is important to note that InV consists of x_1, x_2, \dots, x_q
2032 input variables so that each of which included k sample values so that they are formed as $k \times 1$
2033 matrices. Moreover, p includes the number of q input variables in addition to the number of the
2034 target variable.

2035 **4.3.6 Statistical evaluation metrics**

2036 Different statistical indices were employed to assess the performance of the ELM using the
2037 predefined optimization parameters in the prediction of TE_a, TE_b, TE_c, TE_d , and TE_t by including
2038 various combinations of input variables. The statistical indices including the correlation coefficient

2039 (R), Mean Absolute Error (MAE), Root Mean Square Error ($RMSE$), and $BIAS$, were calculated
 2040 as below:

$$R = \frac{\sum_{i=1}^n (x_i - \bar{x})(y_i - \bar{y})}{\sqrt{\sum_{i=1}^n (x_i - \bar{x})^2 \sum_{i=1}^n (y_i - \bar{y})^2}} \quad (12)$$

$$MAE = \frac{\sum_{i=1}^n |y_i - x_i|}{n} = \frac{\sum_{i=1}^n |e_i|}{n} \quad (13)$$

$$RMSE = \sqrt{\left(\frac{1}{n}\right) \sum_{i=1}^n (x_i - y_i)^2} \quad (14)$$

$$BIAS = \frac{\sum_{i=1}^n (y_i - x_i)}{n} \quad (15)$$

2041 where x_i and y_i are observed and predicted TE values, respectively, and \bar{x} and \bar{y} are the mean
 2042 values of observed and predicted values of TE , respectively. $e_i = y_i - x_i$ denotes the error. n is
 2043 the number of samples.

2044 4.3.7 Uncertainty analysis

2045 Previous research in the field of machine learning has indicated that recognizing the inherent
 2046 uncertainty within machine learning algorithms is essential for improving their reliability
 2047 (Habibian, 1994; Huang et al., 2022). Accordingly, in this study, the uncertainty analysis
 2048 (Shaghghi et al., 2017) was implemented using parameters such as individual prediction error
 2049 (e_i), mean prediction error (\bar{e}), and standard deviation of prediction error (S_e) as shown below
 2050 (Ebtehaj et al., 2018; LaMorte, 2021):

$$e_i = P_i - T_i \quad (16)$$

$$\bar{e} = \sum_{i=1}^n e_i / n \quad (17)$$

$$S_e = \sqrt{\sum_{i=1}^n (e_i - \bar{e})^2 / n} \quad (18)$$

2051 where P_i and T_i indicate predicted and observed values, respectively.

2052 **4.4 Results and discussion**

2053 In this section, the laboratory observations as well as the ELM performance in prediction of TE
2054 values for both spherical and cylindrical particles using different combinations of input variables
2055 were investigated. Then, the best model (i.e., combination) based on each of spherical and
2056 cylindrical datasets was introduced. Furthermore, sensitivity analysis was performed for both
2057 spherical and cylindrical data to examine the variation trend of the target variable (TE_t), based on
2058 the derivative of Eq. 11, with respect to each input variable. Finally, uncertainty analysis was
2059 performed to evaluate the performance of the proposed ELM-based model for both spherical and
2060 cylindrical particles.

2061 **4.4.1 Laboratory observations**

2062 In this study, three water depths ($w = 0.1, 0.15, \text{ and } 0.2 \text{ m}$), three velocities of flow ($u_f =$
2063 $0.1, 0.2, \text{ and } 0.3 \text{ m/s}$), six channel's bed deepening in FOV ($h =$
2064 $0.135, 0.17, 0.185, 0.220, 0.235, \text{ and } 0.270 \text{ m}$), and four slopes applied to the channel's bed in
2065 FOV ($s = 0^\circ, 10^\circ, 16.25^\circ, \text{ and } 22.5^\circ$) were included in conducting the experiments. As stated
2066 above, 57 out of 72 possible test scenarios excluding outlier repeats were considered in this study.
2067 As stated in Table 8, the mean values of TE values for cylindrical particles were more than
2068 spherical MPs in all zones which was consistent with the results obtained in previous studies (e.g.
2069 Bigdeli et al. (2024)). This was confirmed by the observed average \pm standard deviation values of
2070 TE_t for spherical (0.67 ± 0.43) and cylindrical (0.77 ± 0.32) particles. The calculated standard
2071 deviation of the spherical data was slightly larger than that of cylindrical data indicating the wider
2072 range of TE_t values because of the shape on the results of test scenarios. The other important point
2073 that should be noted was the similarity of the measured mean values for two hydraulic parameters,
2074 i.e., w and u_f . The average values for w and u_f in spherical tests were observed to be very close
2075 to those in cylindrical experiments, indicating consistency in the test scenario implementations. In
2076 other words, the predefined values for w and u_f in spherical and cylindrical tests were set to be
2077 the same while the measured values of these two parameters were observed with a very small
2078 difference. The other two effective parameters, h and s , were predefined in each test scenario due
2079 to the geometric conditions of the channel.

2080 Among the observed experimental test scenarios, there are several key findings regarding the
 2081 behavior of both spherical and cylindrical MPs: (1) In cases with low flow velocity ($u_f =$
 2082 0.1 m/s): Most MPs were deposited in the FOV due to the slower flow velocity near the bed,
 2083 which has greater influence on the TE_t values than other hydraulic factors. (2) In high flow velocity
 2084 cases ($u_f = 0.3 \text{ m/s}$): The deepening of the channel along the FOV has significant impact on
 2085 increasing the deposition of MPs, especially at slopes of 16.25° and 22.5° . (3) In steeper slope
 2086 cases ($s = 16.25^\circ$ and 22.5°): TE_t values were increased, as the particles velocity was reduced due
 2087 to the continuity equation, independent of factors such as u_f and h .

2088 4.4.2 Hyperparameters Tuning

2089 Hyperparameter tuning is crucial for optimizing the performance of machine learning models. The
 2090 optimized parameters for ELM training are presented in Table 12. It is important to highlight that
 2091 identical parameter values were employed for both spherical and cylindrical datasets to ensure
 2092 consistency in the modeling process.

2093 For the ELM, separate training sessions were conducted with the spherical and cylindrical datasets.
 2094 For both datasets, a total of 64 samples for training and 27 samples for testing were allocated. This
 2095 distribution was achieved using a random selection method, designating 70% of the data for
 2096 training and the remaining 30% for testing.

2097 During the hyperparameter tuning, it was observed that setting the number of hidden neurons to
 2098 25 and the number of iterations to 25,000 provided the optimal performance for the ELM. These
 2099 settings were determined after extensive experimentation and evaluation of the model's
 2100 performance metrics. Table 12 lists the key hyperparameters that were fine-tuned to achieve the
 2101 best results. The sigmoid function with 4 input neurons and 5 output neurons was used as the
 2102 activation function, which followed the ELM learning rule.

2103 **Table 12.** User-specified parameters for the ELM model

Parameter	Value
Hidden Neuron	25
Activation Function(s)	Sigmoid Function

Parameter	Value
Neurons - Input	4
Neurons - Output	5
Learning rule	ELM
Number of iterations	25,000

2104

2105 4.4.3 Selecting the best Model

2106 As mentioned earlier, 30% of the data (27 samples) was used to test the performance of the model.
2107 The prediction results obtained based on those samples of spherical and cylindrical MPs deposition
2108 in different zones of FOV using Model 1 are depicted in Figure 20 and Figure 21, respectively.
2109 Model 1 contained all four effective and dimensionless parameters included in Eq. 10 and
2110 estimated TEs for both spherical and cylindrical MPs. As shown in Figure 20, TE_a , TE_b , TE_c ,
2111 TE_d , and TE_t were estimated almost close to the observed values, $R = 0.96, 0.91, 0.94, 0.93$ and
2112 0.95 , respectively, in spherical tests. The Model 1 also performed well in prediction of TE_a , TE_b ,
2113 TE_c , TE_d , and TE_t by $R = 0.94, 0.92, 0.91, 0.83$ and 0.91 , respectively, in cylindrical tests (Figure
2114 21). This similarity provided the first evidence of the accuracy and validity of predictive model.
2115 Additionally, MPs deposition patterns were very similar between the physical model and the
2116 corresponding ELM model. These observations demonstrated that the ELM model was promising
2117 for predicting MPs deposition. The relatively small errors were obtained for predicting TE_a , TE_b ,
2118 TE_c , TE_d , and TE_t in spherical tests with $MAE = 0.12, 0.09, 0.01, 0.01$, and 0.09 , and $RMSE$ of
2119 $0.14, 0.13, 0.02, 0.02$, and 0.12 , respectively. This was followed by the obtained results for TE_a ,
2120 TE_b , TE_c , TE_d , and TE_t in cylindrical tests with $MAE = 0.11, 0.10, 0.04, 0.07$, and 0.11 , and
2121 $RMSE$ of $0.12, 0.23, 0.05, 0.010$, and 0.14 , respectively. $BIAS$ was obtained to be very small and
2122 positive in the results for both spherical and cylindrical tests. This indicated that ELM had a slight
2123 overestimation in predicting all TEs , resulting in insignificant impact to the prediction results. In
2124 spherical tests, the $BIAS$ values were found to be $0.03, 0.004, 0.01, 0.001$, and 0.01 for $TE_a, TE_b,$
2125 TE_c, TE_d , and TE_t , respectively. $BIAS$ values of predictions for cylindrical particles for the

2126 quantities TE_a , TE_b , TE_c , TE_d , and TE_t were found to be 0.01, 0.04, 0.01, 0.02, and 0.03,
 2127 respectively.

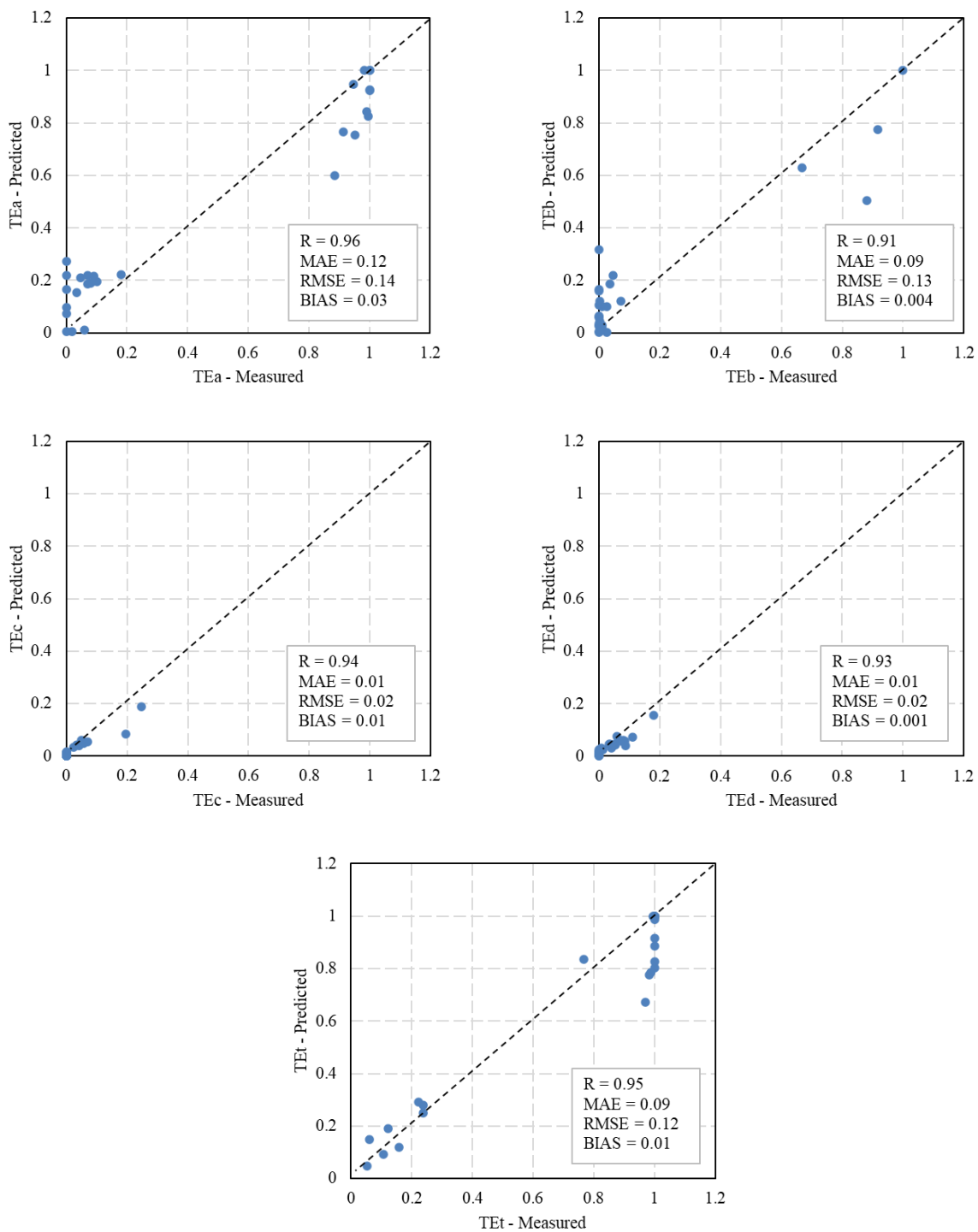


Figure 20. The prediction of spherical MPs deposition in different zones of FOV using ELM, Model 1.

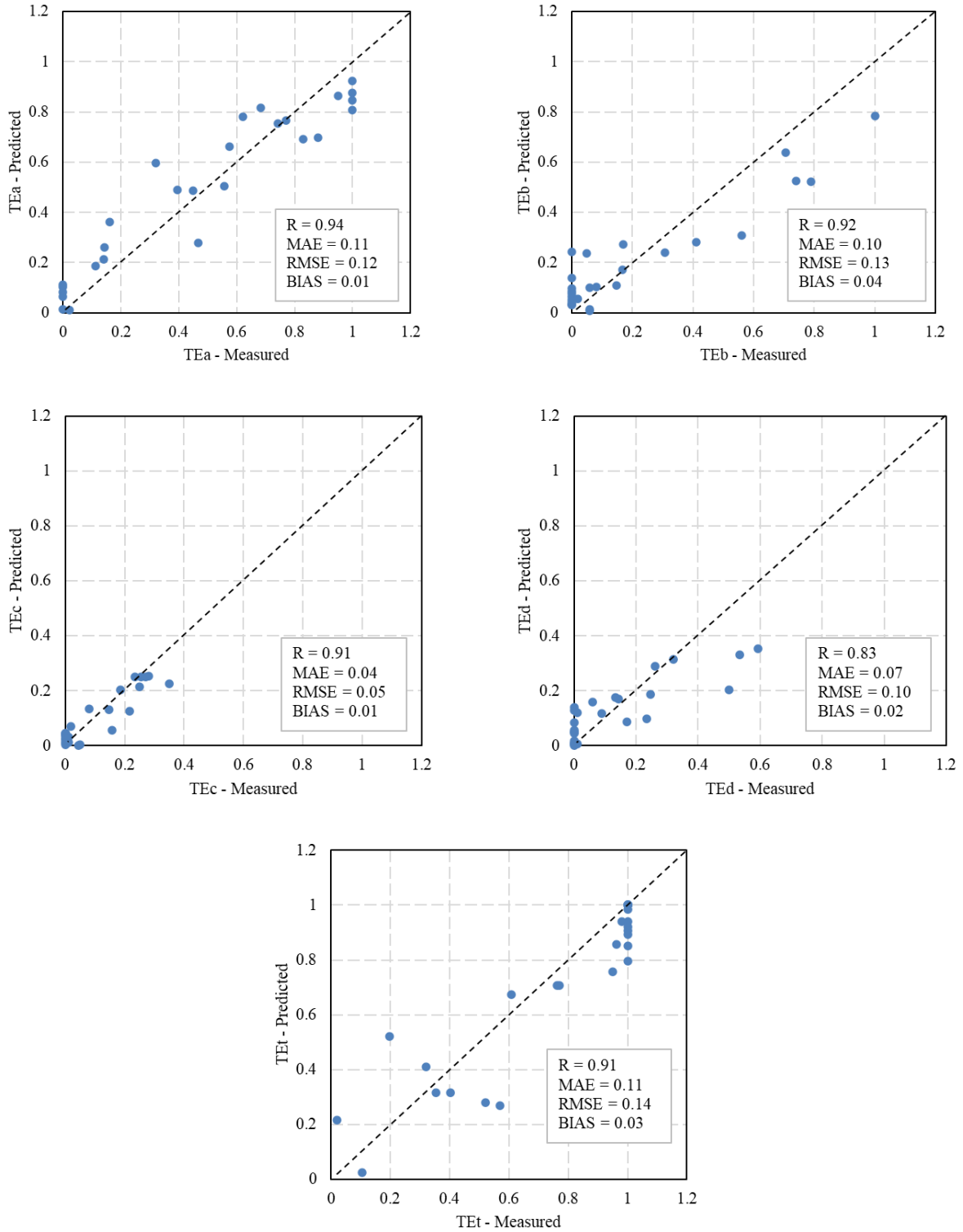


Figure 21. The prediction of cylindrical MPs deposition in different zones of FOV using ELM, Model 1.

2129 The analysis of the ELM performance was further continued by investigating the outcomes
2130 obtained from the other Models including different input combinations to determine the best
2131 Model. In other words, ELM training was continued with different combinations to find the best
2132 performance in terms of R , MAE , MSE , and $BIAS$. For this purpose, the average value of the
2133 statistical indices including all TE s values was calculated for all models, as shown in Table 13, to
2134 facilitate the comparison. In this regard, as mentioned before, 10 additional models (Model 2 to
2135 11) were analyzed alongside Model 1 by excluding one to two input variables. Among Models 2
2136 to 11, Model 10 had the poorest performance based on spherical ($R = 0.42$, $MAE =$
2137 0.19 , $RMSE = 0.23$, and $BIAS = 0.04$) and cylindrical ($R = 0.45$, $MAE = 0.16$, $RMSE =$
2138 0.20 , and $BIAS = 0.06$) data. In Model 10, $\pi_3 = \frac{u_f d_p}{v}$ and $\pi_4 = s$ were excluded from the input
2139 variables. This indicated that π_3 and π_4 were more important compared to the other parameters
2140 included in Eq. 10, when predicting TE s values. On the other hand, Model 8 and 11 showed weak
2141 performance. This indicated that excluding π_4 had the most significant impact on the ELM's
2142 performance, as confirmed by both spherical and cylindrical results.

2143 According to Table 13, Model 1 performed the best based on both spherical ($R = 0.94$, $MAE =$
2144 0.06 , $RMSE = 0.09$, and $BIAS = 0.01$) and cylindrical ($R = 0.90$, $MAE = 0.09$, $RMSE =$
2145 0.11 , and $BIAS = 0.02$) datasets, as stated before. This indicated an inevitable influence of each
2146 dimensionless parameter on TE values.

2147 In general, the best ELM model using spherical data exhibited slightly higher accuracy in
2148 predicting TE s values than the best model using cylindrical data, although Model 1 showed the
2149 performance in both datasets. It is important to note that ELM model's results are consistent with
2150 the fundamental behaviors and properties of MPs observed under real conditions (simulated by
2151 conducting experimental tests).

2152

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Table 13. The average values of statistical indices of all TEs obtained for Models 1 to 11

Model	R		MAE		RMSE		BIAS		Best Model	
	Sphr.*	Cyln.**	Sphr.	Cyln.	Sphr.	Cyln.	Sphr.	Cyln.	Sphr.	Cyln.
Model 1	0.94	0.90	0.06	0.09	0.09	0.11	0.01	0.02	☒	☒
Model 2	0.81	0.79	0.10	0.11	0.13	0.15	0.01	0.03	☒	☒
Model 3	0.80	0.82	0.11	0.11	0.15	0.14	0.01	0.02	☒	☒
Model 4	0.76	0.77	0.13	0.11	0.18	0.16	0.03	0.01	☒	☒
Model 5	0.67	0.57	0.15	0.14	0.19	0.19	0.02	0.01	☒	☒
Model 6	0.68	0.75	0.13	0.12	0.19	0.17	0.04	0.04	☒	☒
Model 7	0.65	0.62	0.15	0.14	0.21	0.18	0.06	0.02	☒	☒
Model 8	0.49	0.52	0.16	0.16	0.21	0.21	0.03	0.05	☒	☒
Model 9	0.63	0.68	0.16	0.14	0.20	0.18	0.05	0.04	☒	☒
Model 10	0.42	0.45	0.19	0.16	0.23	0.20	0.04	0.06	☒	☒
Model 11	0.51	0.53	0.16	0.16	0.21	0.20	0.03	0.02	☒	☒

2158 *Spherical

2159 **Cylindrical

2160 **4.4.4 Sensitivity analysis**

2161 In this study, sensitivity analysis was performed for each input variable included in the best model
 2162 derived from the optimal ELM model. The best model based on either spherical or cylindrical data
 2163 was Model 1. According to the PDSA technique, the matrices derived from the optimal ELM
 2164 model were substituted into Eq.11, and subsequently, the partial derivative of TE_t with respect to
 2165 each input variable included in the best model was calculated. The PDSA results of the ELM-
 2166 based equation for all input variables which were included in the best model were depicted in
 2167 Figure 22 and Figure 23, based on spherical and cylindrical data. It was observed that the partial
 2168 derivative (PD) obtained based on both spherical and cylindrical data in negative and positive
 2169 values resulted in a decrease and an increase in the target variable (TE_t), respectively.

2170 Among the input variables included in Model 1 (i.e., $\pi_1, \pi_2, \pi_3, \pi_4$), the values indicate the
 2171 sensitivity of TE_t , measured for spherical MPs, to π_2 were calculated as the highest (see Figure

2172 22). It is important to note that π_4 also had significant effect on the estimation of TE_t . The PD
 2173 values for π_2 varied from -15 to $+12$. It was observed that π_1 and π_3 had almost the same effect
 2174 on the prediction of TE_t .

2175 The PD values for four dimensionless input variables included in Model 1 and obtained based on
 2176 cylindrical data were shown in Figure 23. Similar to the results of sensitivity analysis based on
 2177 spherical data, π_2 had the highest PD values (when $\pi_2 = 0.7$) and the largest effect on the predicted
 2178 TE_t , although the variation is limited to a larger range (-12 to 20). Contrary to what was obtained
 2179 based on spherical data, π_1 also had an undeniable influence on predicting TE_t . In addition, π_3
 2180 and π_4 had almost the same impact on the prediction of TE_t . The methodology employed here
 2181 demonstrates how a sensitivity analysis can be used to understand the relative influence of different
 2182 parameters affecting MPs deposition in various applications.

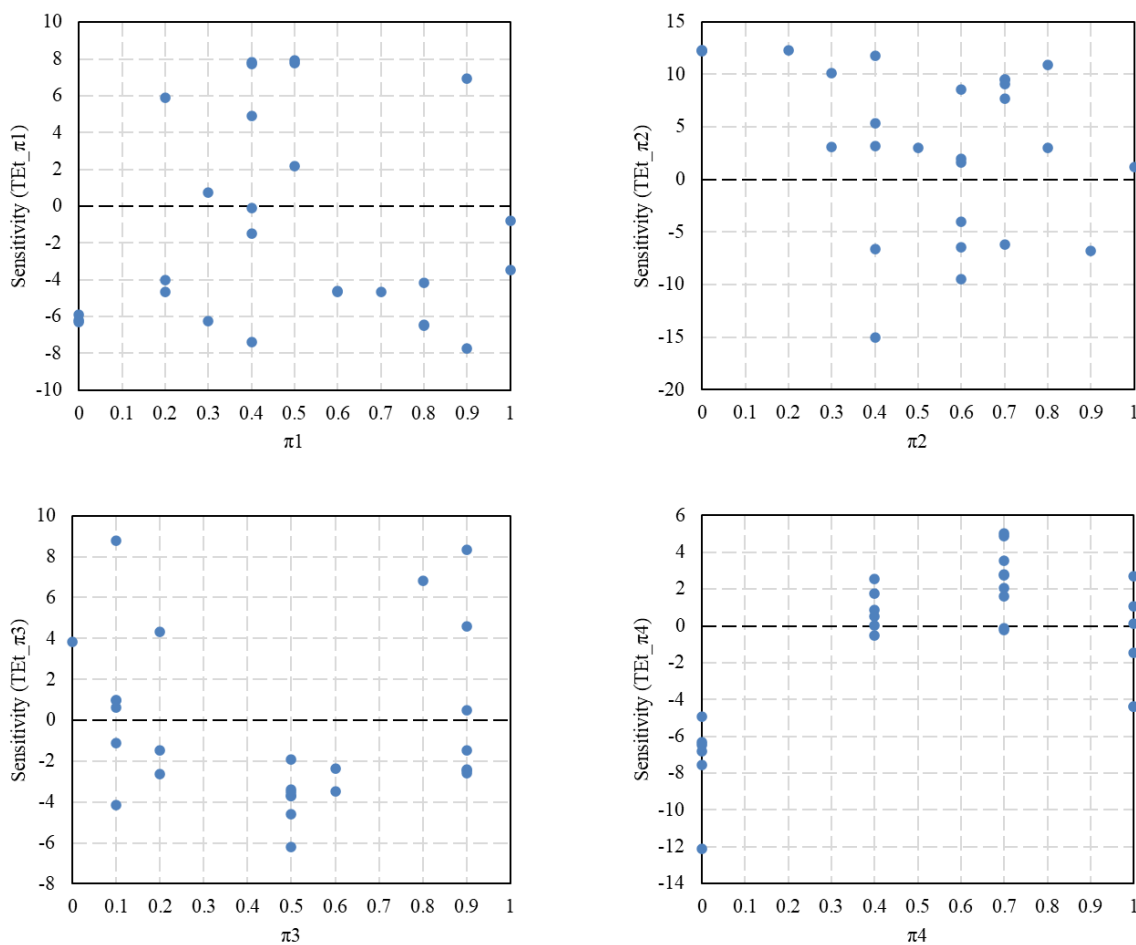


Figure 22. Sensitivity analysis results obtained from Model 1 based on spherical data.

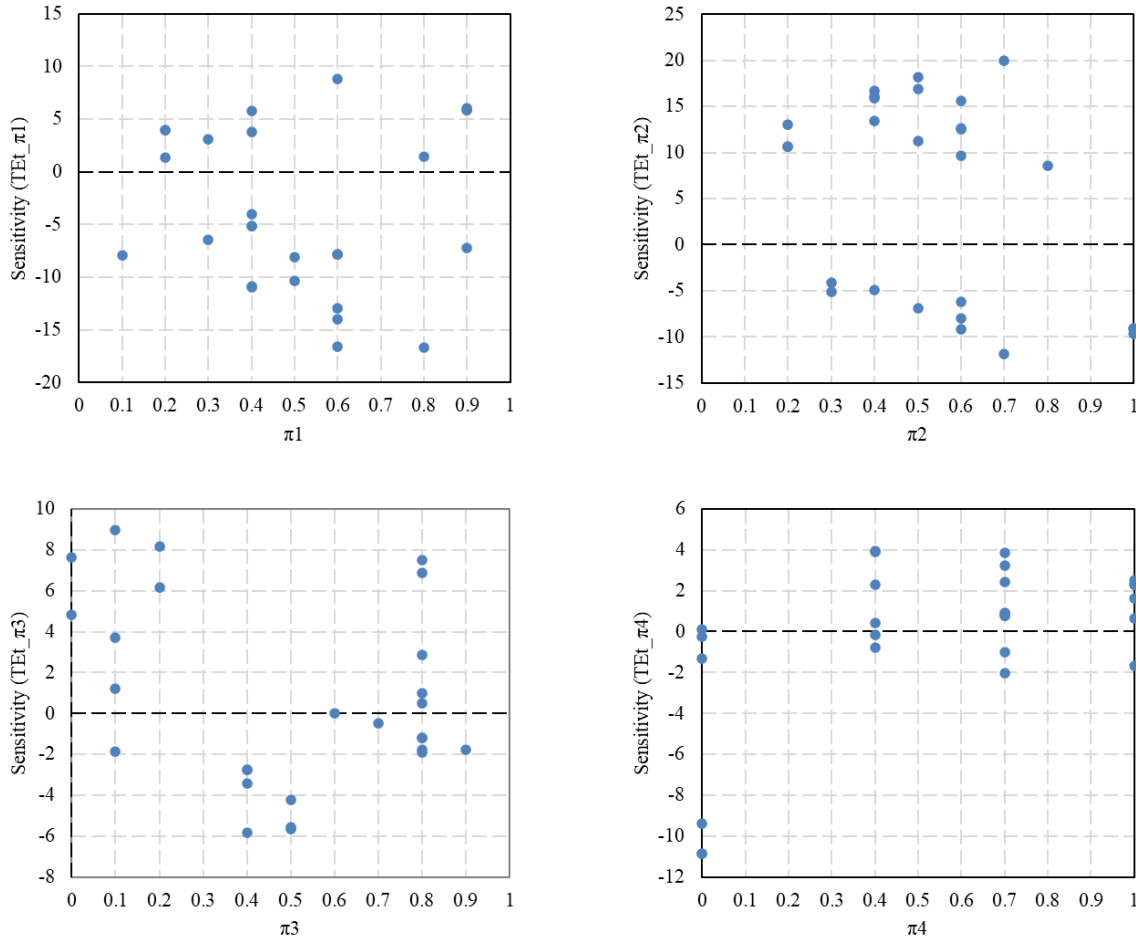


Figure 23. Sensitivity analysis results obtained from Model 1 based on cylindrical data.

2183 **4.4.5 Uncertainty analysis**

2184 In the present study, the uncertainties in predicting TE_t values, as the target variable, were provided
 2185 by evaluating the ELM model quantitatively. In this regard, the positive or negative \bar{e} value shows
 2186 whether the prediction model is overestimating or underestimating the target variable (TE_t).

2187 The confidence band was defined by \bar{e} and S_e values and using the Wilson score method without
 2188 continuity correction and a confidence level of 95%. It should be noted that since the sample size
 2189 was under 30, we had to switch to using a t-value instead of a Z-score to accommodate the smaller
 2190 sample size and dependence on sample standard deviation (LaMorte, 2021). Thus, a confidence
 2191 band of almost 95% was obtained for both spherical and cylindrical data using $\pm 1.706S_e$ values.
 2192 Accordingly, the width of uncertainty band (WUB) was calculated, which indicates the gap

2193 between the upper and lower uncertainty bands. The results of the uncertainty analysis, including
 2194 \bar{e} , S_e , WUB , and 95% prediction interval error (PEI) were summarized in Table 14. The observed
 2195 narrow uncertainty bands (± 0.139 and ± 0.155) and small 95% prediction error intervals, (0, 0.20)
 2196 and (0, 0.25), for spherical and cylindrical data, respectively, indicate the reliability of the ELM
 2197 model for the present application.

2198 **Table 14.** Results of the uncertainty analysis for ELM in testing phase

Method	Number of Test Samples	Dataset Source	Mean Error (\bar{e})	Standard Deviation of Absolute Error (S_e)	Width of uncertainty band (WUB)	95% prediction error interval (PEI)
ELM	27	Spherical	0.063	0.081	± 0.139	(0, 0.20)
		Cylindrical	0.096	0.091	± 0.155	(0, 0.25)

2199

2200 4.5 Conclusion

2201 This study aimed to evaluate the capability of machine learning methods in predicting of MPs
 2202 deposition, at the same particle diameter (d_p) and dynamic viscosity of water (ν) in a generic
 2203 experimental setup and included a comparative analysis of different potential input combinations.

2204 ELM was used as a subset of AI to evaluate the feasibility of predicting the deposition of spherical
 2205 and cylindrical MPs under different test scenarios (i.e., various water depths (w), velocities of the
 2206 water flow (u_f), water depths in the channel where it undergoes a deepening in the FOV (h), and
 2207 slopes applied to the channel's bed in FOV (s)) conducted experimentally in the laboratory. Four

2208 dimensionless input variables including $\pi_1 = \frac{w}{d_p}$, $\pi_2 = \frac{h}{d_p}$, $\pi_3 = \frac{u_f d_p}{\nu}$, and $\pi_4 = s$ were found to

2209 be the most effective parameters as a result of dimensional analysis on the deposition of MPs. The
 2210 trap efficiencies (TE s) of MPs, parameters for showing their deposition, were measured in the lab
 2211 for four zones including TE_a (trap efficiency of MPs in zone a), TE_b (trap efficiency of MPs in
 2212 zone b), TE_c (trap efficiency of MPs in zone c), TE_d (trap efficiency of MPs in zone d), and TE_t
 2213 (total trap efficiency of MPs in FOV). The ELM model was trained based on spherical and

2214 cylindrical datasets including TE_s values and different combinations of input variables.
2215 Accordingly, 11 different models were introduced including various combinations of
2216 dimensionless inputs. The best model for spherical and cylindrical data showed high accuracy in
2217 predicting MPs deposition in different zones of FOV.

2218 Model 1 performed best based on both spherical ($R = 0.94$ and $RMSE = 0.09$) and cylindrical
2219 ($R = 0.90$ and $RMSE = 0.11$) datasets. This is the first study to conduct an in-depth sensitivity
2220 analysis, using a proposed equation, on the trained ELM model for each input variable. In this
2221 regard, the partial derivative sensitivity analysis (PDSA) technique was employed for the first time
2222 for analyzing the sensitivity of the target variable (TE_t) to each input variable included in the best
2223 model using the ELM-based equation. Conclusively, the best model selected based on spherical
2224 and cylindrical data (i.e., Model 1) confirmed that π_2 had the largest effect on the predicted TE_t .
2225 In addition, π_1 and π_4 also had a noticeable impact on the prediction of TE_t , which were estimated
2226 based on cylindrical and spherical data, respectively. Moreover, π_3 had the most limited range of
2227 PD values, which had the least impact on the prediction of TE_t compared to the other
2228 dimensionless parameters in both spherical and cylindrical datasets. It is worth noting that the
2229 differences in the obtained ranges of PD values for the dimensionless inputs were not very
2230 significant, highlighting the influence of each input on the prediction of TE_t .

2231 Uncertainty analysis was also used to quantitatively assess the ELM model. Narrow uncertainty
2232 bands (± 0.139 and ± 0.155 for spherical and cylindrical data, respectively) and small 95%
2233 prediction error interval, (0, 0.20) and (0, 0.25) based on spherical and cylindrical data,
2234 respectively, were obtained for the ELM model.

2235 This study demonstrates the promising capability of ELM in predicting MPs deposition observed
2236 under different experimental scenarios in the presence of sufficient data and lays the foundation
2237 for further applications of the ELM method in MPs deposition predictions.

2238 **4.6 References**

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2330

2331

2332 **5 Numerical Modeling of Large Spherical Microplastics Transport and** 2333 **Deposition using CFD-DEM**

2334

2335 **5.1 Abstract**

2336 The transport and deposition patterns of microplastics (MPs) in water are significantly influenced
2337 by the hydraulic characteristics of the surrounding fluid flow, including the water level, flow
2338 velocity, and the bathymetric and geometric conditions of the water bodies. Accordingly, this study
2339 aims to employ an Eulerian-Lagrangian numerical modeling approach to illuminate spherical MPs
2340 transport and deposition patterns by comparing different scenarios applied to flow and channel
2341 geometry. Various test scenarios were simulated to assess the impact of various variables,
2342 including the flow velocity (u_f), depth of turbulent flows (w), deepening of the channel bed (d),
2343 and slopes applied to the bottom of the channel bed (s). The utilized numerical model is based on
2344 the unresolved computational fluid dynamics (CFD) model coupled with discrete element method
2345 (DEM) solver, which was validated against laboratory experimental measurements. The results
2346 demonstrated that the CFD-DEM could simulate MPs transport and deposition with an accuracy
2347 of over 61%. Additionally, the model effectively captured the effects of hydraulic parameters (w ,
2348 u_f , d , and s) on MPs deposition. It was concluded that CFD-DEM model is highly sensitive to the
2349 values of the sliding and rolling friction coefficients, and the number of particles considered for
2350 running simulations.

2351 **Keywords:** Numerical Modeling, Microplastics, Transport, Deposition, CFD-DEM

2352 **5.2 Introduction**

2353 The widespread distribution of microplastics (MPs), driven by their high production rates and rapid
2354 industrial growth, poses a significant threat to the sustainability of aquatic ecosystems and
2355 biodiversity, thereby contributing to environmental pollution. For example, large micro-sized (≤ 5
2356 mm) plastic particles (Rocha-Santos & Duarte, 2015) can be digested by aquatic fauna, enter the
2357 food chain, and persist without degradation (Alvarez-Zeferino et al., 2020). The dynamics of MPs
2358 in aquatic environments is influenced by various factors, including particle-specific characteristics
2359 (such as size, density, shape, and chemical composition), hydrodynamic conditions, and sediment
2360 characteristics of the ambient environment (Bigdeli et al., 2024; Bigdeli et al., 2022; Klaine et al.,
2361 2008; Masliyah, 1998; Wiesner & Bottero, 2017).

2362 Previous studies have provided significant insight into the abundance, shape, color, and size of
2363 MPs in the aquatic environments (Ballent et al., 2016; Bujaczek et al., 2021; Castañeda et al., 2014;
2364 Corcoran et al., 2019). However, their transport and deposition patterns remain largely unknown,
2365 primarily due to the complexity and limited understanding of the processes governing MPs
2366 movement in such dynamic environments. This knowledge gap hinders our understanding of their
2367 environmental impacts on aquatic ecosystems.

2368 Recent studies have demonstrated the effectiveness of numerical modeling tools, particularly
2369 Eulerian-Lagrangian models, in predicting the behavior of MPs in aquatic environments (Declerck
2370 et al., 2019; Díez-Minguito et al., 2020; Ghazizadeh et al., 2024; Isobe et al., 2019; Lebreton et
2371 al., 2012; Liubartseva et al., 2016; Maximenko et al., 2012; Pilechi et al., 2022; Van Sebille et al.,
2372 2015). Numerical modeling is a valuable tool for predicting the fate and transport of MPs, offering
2373 a cost-effective approach alternative to physical modelling and field surveys. It enables the
2374 prediction of pathways, identification of accumulation zones, and potential sources in aquatic
2375 environments, particularly in remote areas, where water quality monitoring is often impractical.
2376 However, it remains crucial to calibrate these with field or laboratory data to ensure their accuracy
2377 and reliability.

2378 Deposition is a critical process influencing the transport of MPs in aquatic environments. Previous
2379 mass balance analyses of plastic pollution have revealed that a substantial portion of plastic
2380 entering these systems remains unaccounted for (Bigdeli et al., 2024; Bondelind et al., 2020b;
2381 Domercq et al., 2022). This gap suggests that our understanding of MPs deposition is incomplete
2382 and that much of the 'missing' plastic may have accumulated in sediment substrates.

2383 Numerical simulations of deposition and transport of MPs in water bodies can be performed using
2384 Eulerian and Lagrangian models. Eulerian grid models operate within a fixed reference frame,
2385 unlike Lagrangian trajectory models, which employ a movable coordinate system in space (Van
2386 Utenhove, 2019). Lagrangian models, often referred to as particle-tracking methods (PTMs), focus
2387 on tracking individual particles, whereas Eulerian models address advection and diffusion
2388 processes at specific locations (computational nodes) (Van Sebille et al., 2018). PTMs are
2389 recognized as cost-effective and promising tools for interpreting MPs' behavior, offering valuable
2390 insights into their fate and transport in dynamic flow systems (Hardesty et al., 2017). Lagrangian
2391 models have the advantage of solving fewer equations per node since they do not compute

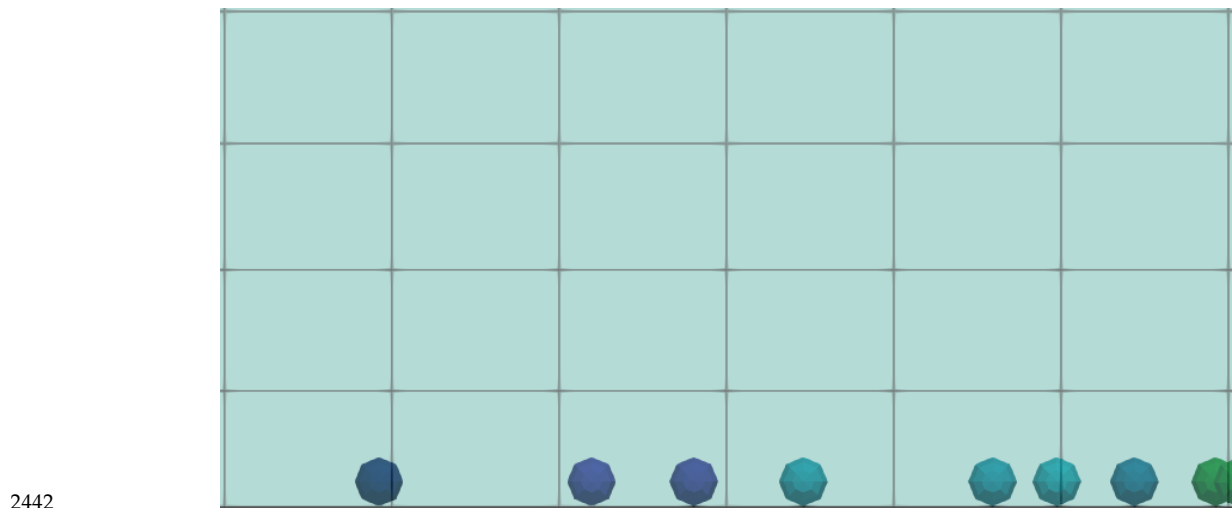
2392 equations for every location in the mesh. However, the strengths and limitations of both methods
2393 make them complementary. Thus, these two approaches can be integrated, with Eulerian
2394 hydrodynamic models providing essential hydrological parameters as input for PTMs. By coupling
2395 particle-tracking models with hydrodynamic models, researchers can analyze potential trajectories
2396 of plastic debris on scales ranging from local (Lammerts, 2016; Rodríguez-Díaz et al., 2020;
2397 Stuparu et al., 2015) to global (Lebreton et al., 2012). This combination enables the assessment of
2398 MPs transport, deposition, distribution, residence time, and fate under various environmental
2399 conditions (Carlson et al., 2017; Critchell & Lambrechts, 2016; Ebbesmeyer et al., 2012; Isobe et
2400 al., 2009; Lebreton et al., 2012; Liubartseva et al., 2018; Van Utenhove, 2019).

2401 There are two primary approaches for numerically modeling MPs particles. The first approach,
2402 known as the kinematic approach, treats particles as an passive entities that move with the
2403 surrounding flow field (Goniva et al., 2010; Pilechi et al., 2022). However, this method does not
2404 account for the local behavior of the individual particles. In contrast, the second modeling strategy
2405 (i.e., dynamic approach) simulates the motion of each particle individually. In other words, this
2406 approach treats particles as active entities whose motion is governed by a balance of forces such
2407 as gravity, friction, drag, and inertia, along with interactions with the surrounding medium. In the
2408 dynamic approach, the movement of particles is determined by the resultant external and internal
2409 forces acting on them, as described by the equation of motion (Pilechi et al., 2022). The most
2410 prominent models in this category are those based on the Discrete Element Method (DEM) and its
2411 variations. The DEM has the potential to capture all granular physical phenomena.

2412 The Eulerian-Lagrangian CFD-DEM method (Anderson & Jackson, 1967) suggests calculating
2413 fluid motion by employing the Navier–Stokes equations within a fluid domain that is partitioned
2414 into cells, as is typical in the finite volume method. The trajectory of particles is modelled in the
2415 DEM model using the hydrodynamic forces calculated by the CFD model at each time step
2416 (Jérémy et al., 2020). DEM can also incorporate the interactions between fluid and different phases
2417 of particulate matter within (e.g., sediment and MPs) based on their shapes and forces they undergo
2418 (Y. Wang et al., 2019; Zhong et al., 2016). Further details about the coupling process of CFD and
2419 DEM has been provided in the subsequent section. Thus, in simulating particle trajectories, the
2420 DEM approach has been employed as a powerful tool in previous studies (Valyrakis et al., 2008;
2421 Zhong et al., 2016).

2422 There are two well-known methods for modeling particle collisions that DEM is capable to
2423 employ: (1) soft particle and (2) hard particle, (also known as soft/hard sphere) approaches. In the
2424 soft particle approach, a pair of particles is allowed to deform or overlap during a collision that
2425 spans multiple time steps (Cundall & Strack, 1979). This makes it possible to calculate the elastic,
2426 plastic, and frictional forces between particles. Additionally, this approach allows for prolonged
2427 collisions between multiple particles, enabling the formation of chains or clusters, with all inter
2428 particle forces being accounted for during the collision. The soft particle approach has been
2429 considered the best balance between accuracy and computational effort in most cases (Aspherix[®],
2430 2024) such as dense granular flows (Gantt et al., 2006). In contrast, the hard particle method
2431 assumes that collisions are instantaneous and occur between only two particles at a time, meaning
2432 only one collision is considered per time step. Because of the instantaneous nature of these
2433 collisions, the forces between particles are often not explicitly calculated (Zhu et al., 2007).

2434 The size difference between the particles and the mesh resolution influences the selection of the
2435 coupling methods, including unresolved, resolved, and hybrid methods. In the unresolved
2436 approach (Figure 24), most particles are smaller than the CFD cell size. Rigid and noncohesive
2437 particles are mostly simulated using an unresolved method. If the particles and mesh cells are of
2438 similar sizes, neither the resolved nor unresolved method alone is ideal, but a hybrid approach can
2439 be applied. For deformable and cohesive particles, it is recommended to use the resolved approach,
2440 in which the particle size is larger than the CFD cell size (Aspherix[®], 2024). While the resolved
2441 method enhances accuracy, it also increases the computational effort required.



2442
2443 **Figure 24.** CFD-DEM coupling using the unresolved approach for simulating spherical MPs

2444 The objective of this study is to investigate the capability of the coupled CFD-DEM (also known
2445 as CFDEM) by employing unresolved coupling method and soft particle collision approach to
2446 simulate spherical MPs that are denser than the flow (i.e., water). The MPs deposition patterns in
2447 some generic cases observed through laboratory test scenarios in a flume setup that mimics a
2448 natural channel. The scenarios that were simulated numerically included various hydraulic
2449 variables, such as the flow velocity, depth of turbulent flows, deepening in the depth of the channel,
2450 and slopes applied to the bed of the channel. Thus, CFD-DEM numerical simulations were
2451 implemented accordingly to assess the performance of the model. The validation process included
2452 the assessment of the performance of the model against laboratory measurements. The numerical
2453 results agreed well with the experimental observations. The analysis of key parameters such as the
2454 static and rolling friction coefficients, the coupling interval, DEM time-step, and the number of
2455 particles considered for running simulations revealed that they have significant influence of on the
2456 model's accuracy and predictive capability.

2457 This paper is structured as follows: Section 5.3 outlines the methodology including experimental
2458 and numerical modeling, covering details such as the physical model setup, governing equations
2459 in the numerical model, and numerical model setup and configurations. Section 5.4 presents the
2460 results obtained from experimental observations and numerical simulations along with a discussion
2461 of the findings. Section 5.5 presents conclusions drawn from the implementation of the test
2462 scenarios numerically.

2463 **5.3 Methodology**

2464 **5.3.1 Physical model**

2465 In this study, different test scenarios were conducted in a tilting flume at the National Research
2466 Council Canada Ocean, Coastal, and River Engineering Research Centre (NRC-OCRE) facility.
2467 Test scenarios consisted of different water levels (w), velocities of the flow (u_f), slopes (s) and
2468 deepening levels (d) applied to the channel bed within the field of view (FOV). The physical model
2469 setup as well as the implementation of the experimental procedure and test scenarios were
2470 considered the same as in the study conducted by Bigdeli et al. (2025). Out of 57 unique
2471 experimental tests, six physical model setups, including 18 test scenarios, were selected for the
2472 numerical modeling in this study, as shown in Table 15. Two water levels ($w = 0.1$ and 0.2 m),
2473 two velocities of flow ($u_f = 0.1$ and 0.3 m/s), two channel's bed deepening in the FOV ($d =$

2474 0, 0.035 and 0.07 m), and three slopes applied to the channel's bed in the FOV ($s =$
 2475 $10^\circ, 16.25^\circ,$ and 22.5°). The total MPs deposition within the FOV, as discussed previously
 2476 (Bigdeli et al., 2025), was measured using trap efficiency (TE), which is the ratio of the number
 2477 of deposited MPs to the total number of MPs entering the FOV. The flow water had a density of
 2478 1000 kg/m³. A total of 1000 polyethylene terephthalate (PET) particles were used as spherical
 2479 MPs with a density of 1380 kg/m³. The spherical MPs had a diameter of 3 mm.

2480 **Table 15.** Different model setups selected for numerical modeling

Model Setup	w (m)	u_f (m/s)	d (cm)	s ($^\circ$)
1	0.1 and 0.2	0.1 and 0.3	3.5	10
2	0.1 and 0.2	0.1 and 0.3	3.5	16.25
3	0.1 and 0.2	0.1 and 0.3	3.5	22.5
4	0.1 and 0.2	0.1 and 0.3	7	10
5	0.1 and 0.2	0.1 and 0.3	7	16.25
6	0.1 and 0.2	0.1 and 0.3	7	22.5

2481

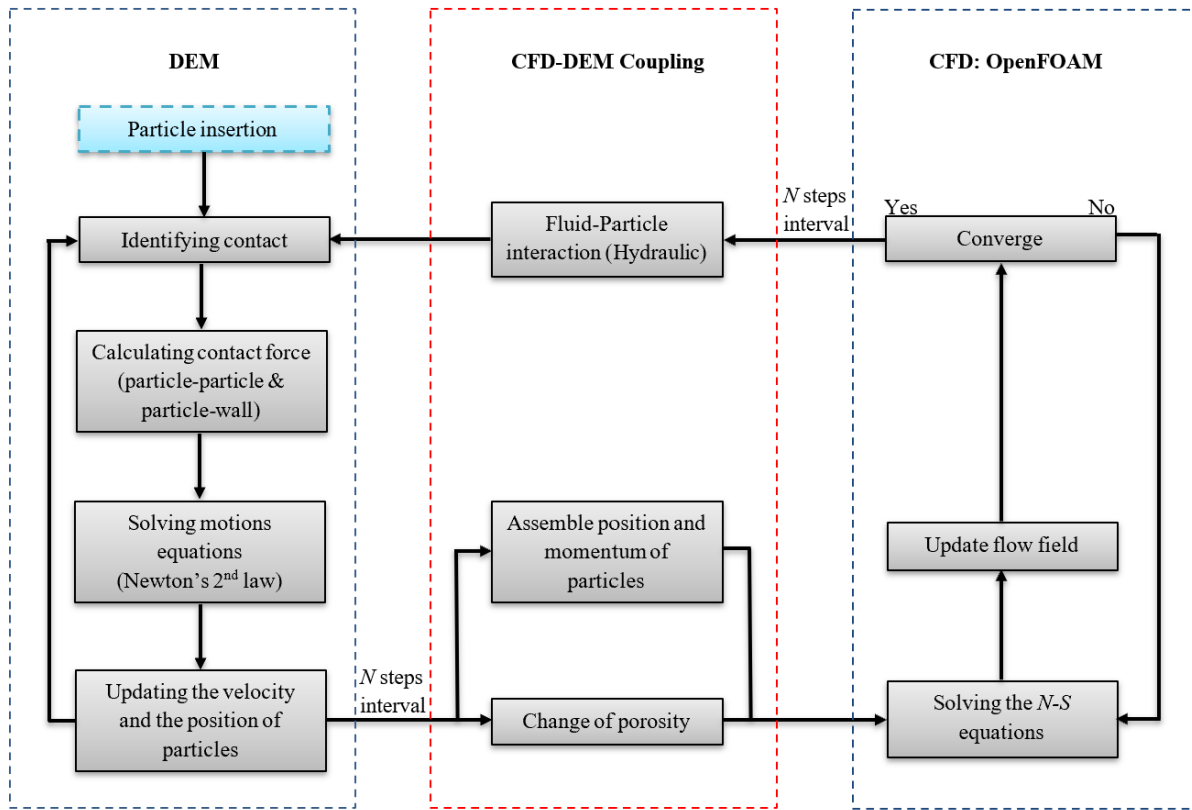
2482 **5.3.2 Numerical Modelling**

2483 As stated above, 18 test scenarios were simulated using the numerical method described below,
 2484 maintaining the same initial and boundary conditions of the generic channel employed in the
 2485 laboratory, as well as the same physical properties of the particles. Furthermore, the numerical
 2486 approach includes a four-way coupling method, which accounts for the interaction between the
 2487 fluid and particles (i.e., indicating a two-way simulation), as well as particle-particle collisions that
 2488 demonstrate two-way modeling. This enhances the accuracy of the numerical simulation, making
 2489 it a faithful representation of the real-world conditions.

2490 **5.3.2.1 CFD-DEM approach**

2491 The model used in this study comprises two primary phases: fluid (i.e., water) flow and spherical
 2492 MPs particles. The motion of the particles were simulated using Discrete Element Method (DEM)
 2493 by solving Newton's second law of motion for each individual particles, while the fluid flow was
 2494 simulated using an OpenFOAM[®] (Jasak, 2009) which is an open-source CFD toolbox. The CFD

2495 simulations were conducted using OpenFOAM’s PISO solver, which applies the Finite Volume
 2496 Method to solve the Navier–Stokes equations, calculating fluid pressure and velocity on an
 2497 Eulerian mesh grid. This information is then passed to the DEM model at each CFD-DEM coupling
 2498 interval to determine the forces (e.g., buoyancy, drag) acting on particles, and calculate their
 2499 trajectories. The DEM modeling of this study was performed using Aspherix® (Aspherix®, 2024)
 2500 which is coupled with OpenFOAM through the CFDEM® library (Goniva et al., 2010). Given the
 2501 above descriptions and to gain better insight into the CFD-DEM coupling algorithm, a flowchart
 2502 of the CFD-DEM coupling method is shown in Figure 25. Accordingly, the particulate system is
 2503 solved using the DEM package, which results in obtaining dynamic information about the
 2504 particles, such as their location and velocity in each fluid cell as well as forces affect their motion
 2505 trajectory. Subsequently, the hydrodynamic forces, including drag, pressure gradient, and viscous
 2506 forces, which are applied to the particles, are averaged over each fluid cell.



2507
 2508 **Figure 25.** Flowchart of the CFD-DEM coupling method adapted from (Hu et al., 2019)
 2509

2510 5.3.2.1.1 CFD

2511 As stated above, the CFD modeling in this study was performed using the incompressible fluid
2512 solver in OpenFOAM, which calculates hydrodynamic properties, such as velocity and pressure,
2513 at each computational mesh cell by solving the Navier-Stokes equations (Bigdeli et al., 2022; Yin
2514 Wang et al., 2019). The solver does not account for free surface changes, which were considered
2515 negligible for the case investigated in this study. Since the CFD simulation was performed in
2516 conjunction with DEM, the CFD solver used a modified version of the governing equations to
2517 reflect the impact of particles occupying the Eulerian mesh cells, as follows:

$$\frac{\partial(\varepsilon_f \rho_f)}{\partial t} + \nabla \cdot (\varepsilon_f \rho_f U^f) = 0 \quad (1)$$

$$\frac{\partial(\varepsilon_f \rho_f U^f)}{\partial t} + \nabla \cdot (\varepsilon_f \rho_f U^f U^f) - \varepsilon_f \nabla \cdot (\mu_f \nabla U^f) = -\nabla p + \varepsilon_f \rho_f g - f^{f,p} \quad (2)$$

2518 where $\varepsilon_f = V_{void}/V_c = 1 - V_p/V_c$ indicates the porosity, also known as the void fraction ($V_{void}, V_c,$
2519 and V_p are the total void volume in a fluid cell, the total volume of the fluid cell, and the occupied
2520 volume due to particles, respectively), ρ_f represents the fluid density, t is time, U^f denotes the
2521 velocity vector of the flow, μ_f indicates the viscosity of the fluid, p indicates the pressure, g is the
2522 gravitational acceleration vector, and $f^{f,p}$ represents the drag force exerted by the particles on the
2523 fluid in the cell.

2524 In this study, the standard $k - \varepsilon$ model was employed in the CFD modeling of the flow, which is
2525 one of the most widely used turbulence models in CFD simulations (Mohsenabadi et al., 2019).
2526 This method is based on the concept of eddy viscosity. Two partial differential equations known
2527 as transport equations can be solved using this turbulence model. The first is the turbulent kinetic
2528 energy (k) and the second is the turbulent eddy dissipation (ε), which is the dissipation rate of the
2529 turbulent kinetic energy. Further details can be found in Bigdeli et al. (2023) and Mohsenabadi et
2530 al. (2019).

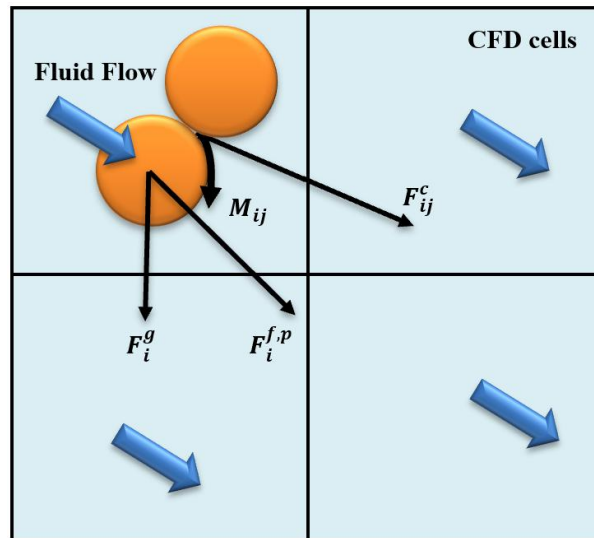
2531 5.3.2.1.2 DEM

2532 The calculation of a particle's trajectory in DEM model is performed based on standard discrete
2533 element theory accounting for force and torque balances (Fonceca et al., 2021; Yin Wang et al.,
2534 2019):

$$m_i \frac{dU_i^p}{dt} = \sum_{j=1}^{n_i^c} F_{ij}^c + F_i^{f,p} + F_i^g \quad (3)$$

$$I_i \frac{d\omega_i}{dt} = \sum_{j=1}^{n_i^c} M_{ij} \quad (4)$$

2535 where m_i denotes the mass of particle i , U_i^p is the translational velocity of particle i , F_{ij}^c is the
 2536 contact force applied by particle j or the wall(s) on particle i . $F_i^{f,p}$ and F_i^g are the particle-fluid
 2537 interaction (including drag force in this case) and gravitational forces, respectively, which act on
 2538 particle i in the fluid-phase cell, as shown in Figure 26. I_i represents the moment of inertia and ω_i
 2539 is the angular velocity of particle i , and M_{ij} is the torque exerted by particle j or the wall(s) on
 2540 particle i . It should be noted that there are other forces, such as viscosity, which should be
 2541 considered in the formulation of different fluid-particle interactions, as explained in depth by
 2542 Smuts (2015). Viscous force arises from the shear stress within the fluid caused by its velocity
 2543 gradient. Accordingly, this force drives particles towards areas of reduced shear (Smuts, 2015).
 2544 This force was considered to be small in this study since it becomes significant when particles are
 2545 large, the surrounding fluid is highly viscous, the relative velocities between particles are high
 2546 (such as during a potential collision), and the distance between particles is much smaller than their
 2547 radius.



2548

2549 **Figure 26.** Schematic diagram of CFD cells including fluid flow, and drag forces and torques
 2550 exerted on particles adapted from (Yin Wang et al., 2019)

2551 There are various equations in the literature (Smuts, 2015; Yin Wang et al., 2019; Zhu, 2019)
2552 which have been proposed as equations for calculating the drag force ($F_i^{f,p}$) (e.g., Di Felice (Di
2553 Felice, 1994) and Koch and Hill (Koch & Hill, 2001). It has been reported that the Di Felice model
2554 is a promising model for developing relations of the drag force around a single particle in
2555 multiphase systems that include non-spherical particles (Yin Wang et al., 2019).

2556 However, the more recent drag force model, i.e., the Koch and Hill correlation (Koch & Hill,
2557 2001), considers particle-particle interactions based on lattice-Boltzmann simulations (Smuts,
2558 2015; Van Buijtenen et al., 2011; Vangö, 2019; Yin Wang et al., 2019; Zhu, 2019). Therefore, in
2559 this study, the Koch and Hill model (Koch & Hill, 2001) was employed because it provides a better
2560 estimation of drag forces around particles based on the literature (Smuts, 2015) compared with
2561 other models such as Di Felice (Di Felice, 1994).

2562 The combination of the Hertz contact model with Coulomb's friction law was utilized to
2563 characterize the interaction between particles in both the normal and tangential directions (Ma et
2564 al., 2019; Smuts, 2015). The Hertz contact model, including the non-linear contact relationship,
2565 was employed to calculate the normal elastic force (Ma et al., 2019). Tangential shear forces can
2566 be generated during collisions, leading to the development of a torque. Many models apply a
2567 Coulomb-type friction law, occasionally incorporating damping, to calculate tangential contact
2568 force. It is important to note that the Coulomb's friction criterion limits the tangential force using
2569 the sliding friction coefficient (Ma et al., 2019).

2570 The model parameters were set to the values listed in Table 16. Key parameters selected during
2571 implementation of coupled CFD-DEM simulations. To reduce computational time, the stiffness of
2572 the particles is specifically reduced by lowering the Young's modulus value. This commonly used
2573 technique allows for the use of larger time steps in simulations. (Smuts, 2015; Xiong et al., 2021).
2574 Selecting the Young's modulus in DEM serves as a preliminary stage in simulation, given its
2575 critical role in determining the DEM's time step and assessing particle overlap (Xiong et al., 2021).
2576 A broad spectrum of Young's modulus values (ranging from 10^6 to 10^{10} Pa) has been employed in
2577 numerical simulations to analyze particulate systems (Akbarzadeh, 2014; Hu et al., 2019; Kong et
2578 al., 2021; Ma et al., 2019; Smuts, 2015; Xiong et al., 2021). Relaxing the stiffness of particles
2579 prolongs their contact duration, enabling the utilization of larger time steps while ensuring an
2580 adequate resolution of collision events. Furthermore, the literature indicates that reducing

2581 magnitude of the Young's modulus does not substantially impact the physical behavior of the
 2582 particulate system. In other words, studies have demonstrated that reducing the particle stiffness
 2583 by two to three orders of magnitude typically does not significantly alter the overall behavior of
 2584 simulations (Hoomans et al., 1996; Tsuji et al., 1993; Wassgren & Sarkar, 2008). In this study, the
 2585 Young's modulus was set to 5 MPa in the particulate phase for both particle-particle and particle-
 2586 wall contacts according to the literature (Ma et al., 2019). In addition, the restitution coefficient,
 2587 which represents the ratio of the rebound velocity of a particle to its velocity before the collision
 2588 with other particles, was set to 0.5 (Ma et al., 2019; Xiong et al., 2021) to consider energy
 2589 dissipation within the fluid-particulate system.

2590 Unlike sliding friction, which limits the tangential force, rolling resistance is frequently employed
 2591 as a practical method to estimate the influence of the particle shape, as observed in various studies
 2592 (Ai et al., 2011; Cheng et al., 2017; Estrada et al., 2011; Iwashita & Oda, 1998). It is quantified
 2593 through a torque that accounts for particle rotation which has been described in depth by Ma et al.
 2594 (2019). In this study, the sliding and rolling friction coefficients were set to 0.01 and 0,
 2595 respectively, in both particle-particle and particle-wall interactions to represent the closest results
 2596 to the real condition.

2597 **Table 16.** Key parameters selected during implementation of coupled CFD-DEM simulations

Fluid (CFD)/Particulate (DEM) phase items	Parameter	Selected value
Fluid phase (water flow at 20 °C and 1 atm)	Density (kg/m ³)	998.2
	Viscosity (Pa.s)	1×10^{-3}
	CFD cell size (m)	1×10^{-2}
Particulate phase (spherical particles)	Particle Number	500
	Particle Density (kg/m ³)	1300
	Radius (mm)	3
	Young's modulus (particle-particle contact) (MPa)	5
	Young's modulus (particle-wall contact) (MPa)	5
	Poisson's ratio	0.3

Fluid (CFD)/Particulate (DEM) phase items	Parameter	Selected value
	Restitution coefficient	0.5
	Interparticle sliding friction coefficient	0.01
	Particle-wall sliding friction coefficient	0.01
	Interparticle rolling friction coefficient	0
	Particle-wall rolling friction coefficient	0
Simulation timesteps	DEM Time step (s)	1×10^{-5} and 2×10^{-5}
	CFD Time step (s)	1×10^{-3} and 2×10^{-3}
	Coupling Interval	10 and 100
Global coupled CFD-DEM settings	Gravitational acceleration (m ² /s)	9.81

2598

2599 In Aspherix[®], the critical time step size is determined using two main criteria: the Rayleigh and
2600 Hertz criteria (Aspherix[®], 2024). Li et al. (2005) suggested that it is reasonable to assume that all
2601 energies are transferred via Rayleigh waves on the particle surface in the simulations. Kremmer
2602 and Favier (2001) observed that calculating the critical time step size based on the Rayleigh wave
2603 frequency is a common practice. Typically, the Rayleigh criterion is more significant, particularly
2604 for dense systems, while the Hertz criterion is more applicable to dilute systems (Smuts, 2015).
2605 Rayleigh time (T_R) refers to the duration required for a wave to travel through a solid particle
2606 (Aspherix[®], 2024; Li et al., 2005). This can be estimated as follows (Aspherix[®], 2024; Li et al.,
2607 2005):

$$T_R = \pi R \left(\frac{\rho}{G} \right)^{\frac{1}{2}} / (0.1631\nu + 0.8766) \quad (5)$$

$$G = \frac{Y}{2(\nu + 1)} \quad (6)$$

2608 where R and ρ denote particle radius and density, respectively. G is the shear modulus. ν is
2609 Poisson's ratio, and Y represents Young's modulus of the particle.

2610 It is advised to use a simulation DEM time-step between $0.1 T_R$ and $0.3 T_R$ (Aspherix[®], 2024).
 2611 However, a simulation time-step can be also selected according to the duration of particle contact
 2612 proposed in the Hertzian contact theory, which depends on the relative velocity between particles,
 2613 their mass, radius, Poisson's ratio, and Young's modulus (Y^*) as follows (Aspherix[®], 2024; Li et
 2614 al., 2005):

$$t_{contact,H} = 2.87 \left(\frac{m^{*2}}{r^* Y^{*2} v_{rel}} \right)^{1/5} \quad (7)$$

$$m^* = \left(\frac{1}{m_1} + \frac{1}{m_2} \right)^{-1}, \quad r^* = \left(\frac{1}{r_1} + \frac{1}{r_2} \right)^{-1} \quad (8)$$

$$Y^* = \left(\frac{1 - \nu_1^2}{Y_1} + \frac{1 - \nu_2^2}{Y_2} \right)^{-1} \quad (9)$$

2615 It is recommended that a small fraction of the Hertz contact time ($t_{constact,H}$), typically 5%-15%,
 2616 be used as the DEM time-step (Aspherix[®], 2024). This approach ensures that the contact is
 2617 accurately resolved, thereby preventing an excessive overlap of energy. In this study, T_R and
 2618 $t_{constact,H}$ were calculated to be 0.00014 s and 0.00039 s, respectively. Accordingly, the DEM
 2619 time step size considering both the Rayleigh and Hertz contact time criteria was calculated to be
 2620 2×10^{-5} s.

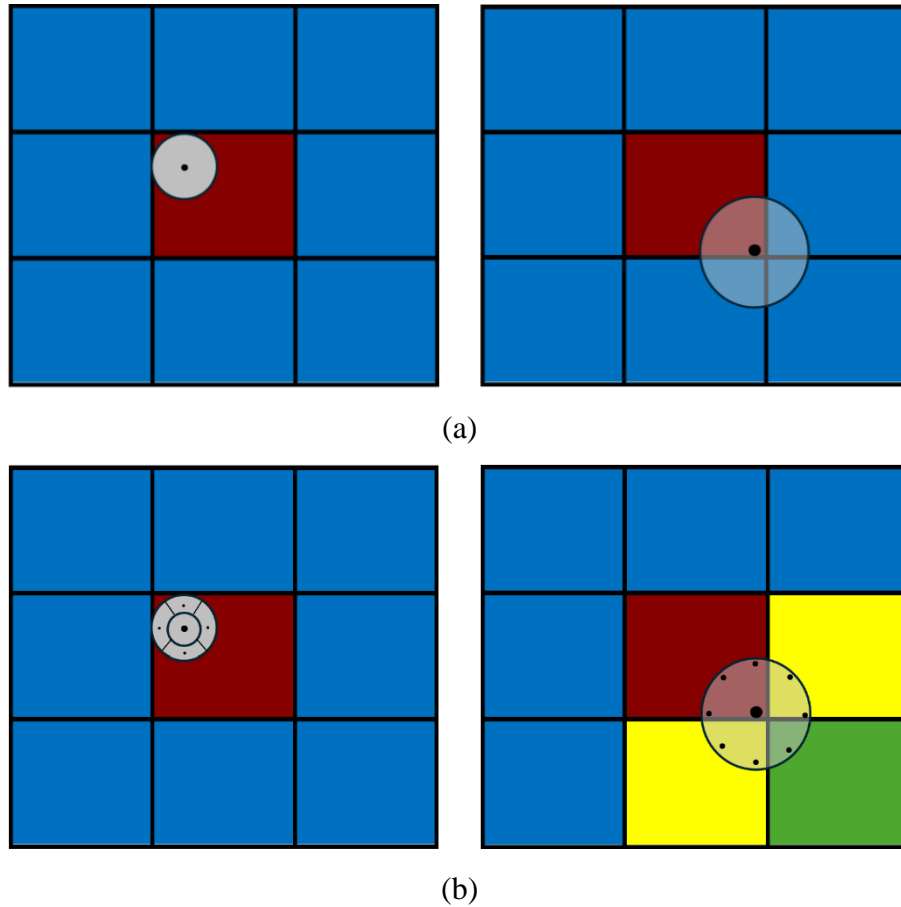
2621 5.3.2.1.3 CFD-DEM coupling configuration

2622 The contribution of particles to the fluid flow can be incorporated with two common formulations,
 2623 called the "Model A" and "Model B", which are employed in different fluid flow-particle
 2624 applications (Smuts, 2015). In both Model A and Model B, the total fluid-particle interaction force
 2625 which is applied on an every individual particle, is the sum of the drag, pressure gradient, and
 2626 viscous forces (typically considered as the dominant forces in interactions between the particle and
 2627 the fluid), and the combined effect of other less significant particle-fluid interaction forces,
 2628 including buoyancy, virtual mass, Saffman lift, Magnus force, and others (Smuts, 2015). These
 2629 two models differ primarily based on how they treat the pressure-source term in the formulations
 2630 (i.e., governing equations). Model B assumes that the pressure originates only from the fluid phase
 2631 while it is attributed to both the fluid and particle (solid) phases in Model A by including the void
 2632 fraction (Smuts, 2015; Zhou et al., 2010). In addition, the pressure gradient force, excluding the
 2633 buoyancy force, was combined with the drag force in Model B. While this approach works when

2634 the forces are aligned in the same direction. It may not be appropriate for all flow systems where
2635 this condition doesn't hold (Smuts, 2015). In this study, Model A was employed in the
2636 implementation of coupling to obtain a better estimation of pressure than Model B (Smuts, 2015;
2637 Zhou et al., 2010). On the other hand, `cfemSolverPiso` was employed, in this study, to couple
2638 CFD (i.e., $k - \varepsilon$ model) and DEM models, for handling the exchange of particle momentum
2639 (Smuts, 2015). The `cfemSolverPiso` solver (Smuts, 2015) is a developed version of OpenFOAM
2640 `pisoFoam` solver. It is recommended that the size of the fluid cell be several times larger than the
2641 particle diameter (Hu et al., 2019), which was confirmed in this study. In addition, to obtain more
2642 accurate results, as reported in the literature (Hu et al., 2019; Smuts, 2015), and given that CFD
2643 and DEM are run separately, it is recommended to consider a CFD time step 50–100 times larger
2644 than that of the DEM. In this study, the time interval was set to 100, after implementing a
2645 sensitivity analysis.

2646 The `voidFractionModel` determines the volume of the DEM particles within the CFD domain. It is
2647 advised to use the “divided” sub-model, as employed in this study, when particle sizes are
2648 comparable to that of a CFD cell (Smuts, 2015). This approach employed satellite points to
2649 distribute a particle's volume across the cells it intersects, making it more accurate than the “center”
2650 model, which assumed that the entire mass is concentrated in the cell containing the particle's
2651 center (Smuts, 2015), as shown in Figure 27. In other words, it is assumed that all particle volume
2652 is placed into the cell encompassing the particle center in the center sub-model (Figure 27a) while
2653 in the divided sub-model, the particle volume is distributed over several cells using splitting the
2654 particle into satellite points when the particle volume spans cell boundaries (Figure 27b). To use
2655 the `voidFractionModels`, the void fraction should be positive (i.e., $\varepsilon_f > 0$) and it is recommended
2656 $\varepsilon_f > 0.2$ (Aspherix[®], 2024).

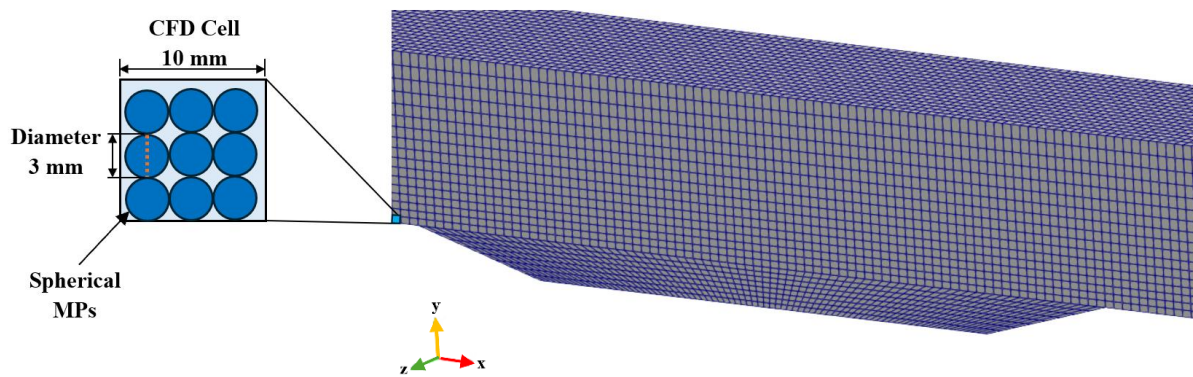
2657



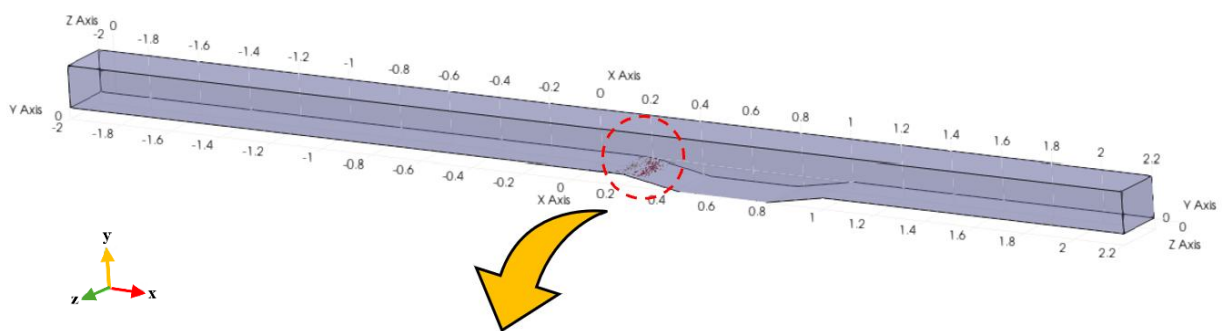
2658 **Figure 27.** Different approaches employed in the voidfractionModel, as the sub-model, for
 2659 calculating the void fraction: (a) center and (b) divided

2660 **5.3.2.2 Numerical model setup**

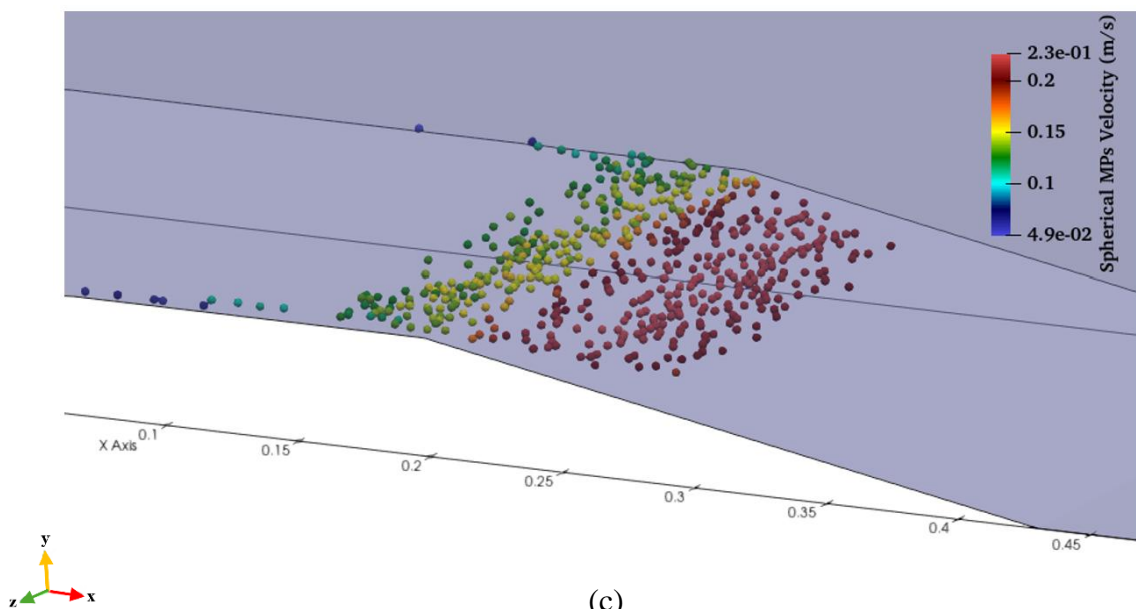
2661 The computational domains used of the CFD-DEM simulations was created with the same
 2662 dimensions and initial and boundary conditions as the physical model setups discussed above to
 2663 be able to compare the results with the experimental observations. The mesh was built with a
 2664 domain size of 4.2 and 0.2 meters in x and z-direction. The size of the domain in the y-direction
 2665 varied according to each test scenario. MPs particles were injected into the flow on the channel
 2666 bed, 0.2 m ahead of the FOV. Figure 28 depicts a section of the computational domain
 2667 corresponding to the field of view in the laboratory experiment. It is important to note that to use
 2668 the unresolved approach, the size of CFD cells should be larger than particles size. This is
 2669 necessary to ensure that there is a balance between the accuracy of the flow and accuracy of
 2670 particles while maintaining the feasible computational effort.



(a)



(b)



(c)

2671

2672 **Figure 28.** Snapshot of the computational domain: (a) mesh configuration, (b) full size, and (c) an
 2673 example of the simulated MPs distribution within the FOV.

2674 The mesh size was refined close to the bottom of the channel and corners several times to improve
 2675 the model stability (see Figure 28a). Mesh considerations play a crucial role in CFD-DEM method,
 2676 where it is commonly assumed that DEM particles should be considerably smaller than CFD cells.
 2677 However, there is a lack of established criteria regarding the extent of this size difference (Smuts,
 2678 2015). The sizing is often specific to each case, with the cell dimensions adjusted according to the
 2679 problem requirements. Nonetheless, empirical observations suggest that cell dimensions should be
 2680 approximately 3-4 times the particle diameter (Smuts, 2015). This range typically represents the
 2681 lower threshold because choosing cell size less than it leads to cells being excessively filled with
 2682 particles, leaving inadequate space for fluid.

2683 All surfaces of the channel were assumed to be smooth during numerical simulations. The lower
 2684 boundary and sidewalls were set as “wall”, and the inlet and outlet were set as a “patch”
 2685 representing open boundaries in the OpenFOAM model. Constant atmospheric pressure, zero shear
 2686 stress, and zero gradient for turbulence quantities were considered as boundary conditions to
 2687 account for the air-water interactions (Bigdeli et al., 2023). The detailed boundary conditions
 2688 employed for running the simulations are shown in Table 17. Boundary conditions selected in
 2689 OpenFOAM[®] for $k-\varepsilon$ model.

2690 **Table 17.** Boundary conditions selected in OpenFOAM[®] for $k - \varepsilon$ model

Boundary	U	P	k	ε	nut
Inlet	U.F.V ¹	Z.G ⁵	T.K.E.I ⁷	Z.G	C ¹⁰
Outlet	I.O ²	f.V ⁶	I.O	I.O	C
Top	S ³	Z.G	Z.G	Z.G	Z.G
Channel Bed	N.S ⁴	Z.G	K.R.W.F ⁸	E.W.F ⁹	N.K.W.F ¹¹
Side Walls	N.S	Z.G	K.R.W.F	E.W.F	N.K.W.F

2691 ¹ uniformFixedValue ⁵ zeroGradient ⁹ epsilonWallFunction

2692 ² inletOutlet ⁶ fixedValue ¹⁰ calculated

2693 ³ slip ⁷ turbulentIntensityKineticEnergyInlet ¹¹ nutkWallFunction

2694 ⁴ noSlip ⁸ kqRWallFunction

2695 To calibrate and validate the numerical simulations, test 2 was selected for calibration and the
 2696 other tests were considered for validation using four hydraulic parameters including w , u_f , d , and
 2697 s . In addition, three mesh sizes (i.e., fine, medium, and coarse) were considered for the sensitivity
 2698 analysis of simulations, as depicted in Table 18. Different configurations used for implementing
 2699 the sensitivity analysis of the mesh for Test 2.

2700 **Table 18.** Different configurations used for implementing the sensitivity analysis of the mesh for
 2701 Test 2

Mesh size	Total number of cells	Cell type	Minimum cell size (m)	Maximum cell size (m)
Fine	138,800	hexahedra	0.005	0.01
Medium	38,400	hexahedra	0.01	0.02
Coarse	4,720	hexahedra	0.01	0.03

2702 Accordingly, the size of the domain and the total number of cells in the computational domain for
 2703 each test scenario based on the fine mesh were listed in Table 19. Different test scenarios selected
 2704 for numerical modeling. Detailed characteristics of the 18 tests are provided in Section 5.4.1 (Table
 2705 20).

2706 **Table 19.** Different test scenarios selected for numerical modeling

Test Scenario	Model Setup	Size of the domain in each direction (x, y, z)	Total number of cells in the domain
1	1	4.2 m, 0.1668 m, 0.2 m	122,000
2	1	4.2 m, 0.189 m, 0.2 m	138,800
3	1	4.2 m, 0.2758 m, 0.2 m	220,800
4	2	4.2 m, 0.1671 m, 0.2 m	122,000
5	2	4.2 m, 0.2098 m, 0.2 m	155,600
6	2	4.2 m, 0.2768 m, 0.2 m	214,400
7	3	4.2 m, 0.1624 m, 0.2 m	113,600
8	3	4.2 m, 0.1934 m, 0.2 m	138,800
9	3	4.2 m, 0.2642 m, 0.2 m	197,600
10	4	4.2 m, 0.1679 m, 0.2 m	108,400

Test Scenario	Model Setup	Size of the domain in each direction (x, y, z)	Total number of cells in the domain
11	4	4.2 m, 0.211 m, 0.2 m	150,400
12	4	4.2 m, 0.2833 m, 0.2 m	209,200
13	5	4.2 m, 0.169 m, 0.2 m	116,800
14	5	4.2 m, 0.2027 m, 0.2 m	142,000
15	5	4.2 m, 0.2855 m, 0.2 m	217,600
16	6	4.2 m, 0.1726 m, 0.2 m	116,800
17	6	4.2 m, 0.1956 m, 0.2 m	142,000
18	6	4.2 m, 0.2871 m, 0.2 m	217,600

2707 5.3.2.3 Sensitivity analysis

2708 Several sensitivity analyses were performed to investigate the effects of different factors, such as
2709 the CFD mesh size, number of plastic particles, coupling interval, values of sliding and rolling
2710 friction coefficients, and the DEM time step size. Accordingly, 250, 500, and 1000 plastic particles
2711 were simulated separately to determine if there is a possibility, based on the obtained numerical
2712 results for TE, to execute the numerical model with fewer particles to decrease computational time
2713 (i.e., reducing computational costs). The simulation results described in section 5.4 showed that
2714 the optimal number for simulating MPs particles is 500 particles.

2715 Several mesh configurations were investigated, featuring a finer mesh near the channel's bed and
2716 corners and a coarser mesh towards the water surface. To determine the optimum mesh size, a
2717 sensitivity analysis was conducted using three different combination numbers in the x (110, 220,
2718 and 420), and z (8, 16, and 20) directions (Figure 28). The y-direction values varied with an
2719 expansion ratio of 1:2 depending on the characteristics of each test scenario. As stated in section
2720 5.4, it was concluded that grid number of 420 in the x-direction and 20 in the z-direction is optimal.

2721 Various combinations of sliding and rolling friction coefficients were also examined to identify
2722 the values that best align the numerical simulations to the experimental observations. The result
2723 indicated that the sliding (f_s) and rolling (f_r) friction coefficients of less than 0.1 provided the
2724 closest match to the experimental data. Accordingly, four combinations were used $f_s = 0.1$ and
2725 $f_r = 0.05$, $f_s = 0.05$ and $f_r = 0.1$, $f_s = 0.05$ and $f_r = 0.05$, and $f_s = 0.01$ and $f_r = 0$. The
2726 optimal combination was found to be $f_s = 0.01$ and $f_r = 0$, as described in section 5.4.

2727 Three coupling intervals, defined as the number of steps between the CFD and DEM models: 10,
2728 50, and 100. Additionally, two DEM time-step sizes of 1×10^{-5} s and 2×10^{-5} s, were evaluated
2729 for the stability and computational efficiency. The coupling intervals of 10 and 100, and combined
2730 with DEM time-steps of 1×10^{-5} s and 2×10^{-5} s was found to run stably, as further explained
2731 in section 3.3.

2732 **5.4 Results and discussion**

2733 In this section, we analyze the result from the laboratory experiments along with the CFD-DEM
2734 simulations for the various tested scenarios. The total deposition of MPs within the FOV is
2735 compared between the experimental and simulated data. Additionally, a sensitivity analysis is
2736 conducted on the numerical model performance, focusing on several key parameters: mesh
2737 configuration, DEM time step size, sliding and rolling friction coefficients, coupling interval
2738 values, and the total number of particles simulated by the numerical model. This comprehensive
2739 analysis aims to identify the optimal conditions for accurate and efficient simulations.

2740 **5.4.1 Laboratory observations**

2741 18 experimental cases, selected from Bigdeli et al. (2025), were investigated as shown in Table
2742 20. The cases were selected to ensure comparability in terms of how each hydraulic parameter
2743 influences the TE value. For example, cases with similar hydraulic parameters but varying flow
2744 velocities were included among the 18 selected cases in this study.

2745 Several findings can be derived from the comparison of TE values in the tests conducted. First,
2746 when water velocity is low (approximately 0.1 m/s), most MPs are deposited within the FOV. This
2747 is because the flow velocity near the bed is slower than in cases with higher flow velocities (e.g.,
2748 0.3 m/s), irrespective of hydraulic factors such as w , d , and s . This highlights the significant
2749 influence of water velocity on the TE value compared to other factors. Second, in higher flow
2750 velocities ($u_f = 0.3$ m/s), the deposition of MPs is notably affected by doubling the channel bed
2751 depth in tests with slopes of $s = 16.25^\circ$ and 22.5° . This contrasts with observations from tests with
2752 a slope of 10° , where steeper slopes resulted in greater water depth, reducing flow velocity due to
2753 the continuity equation. Third, TE values generally increase as the channel bed slope becomes
2754 steeper, which is expected. Steeper slopes reduce particle velocity, regardless of the values of
2755 w , u_f , and d .

Table 20. Characteristics of 18 experimental tests selected for numerical modeling

Test No.	w (m)	u_f (m/s)	d (cm)	s (°)	TE
1	0.1	0.1	3.5	10	0.95
2	0.1	0.3	3.5	10	0.07
3	0.2	0.3	3.5	10	0.118
4	0.1	0.1	3.5	16.25	0.99
5	0.1	0.3	3.5	16.25	0.088
6	0.2	0.3	3.5	16.25	0.11
7	0.1	0.1	3.5	22.5	1
8	0.1	0.3	3.5	22.5	0.10
9	0.2	0.3	3.5	22.5	0.16
10	0.1	0.1	7	10	0.98
11	0.1	0.3	7	10	0.13
12	0.2	0.3	7	10	0.17
13	0.1	0.1	7	16.25	1
14	0.1	0.3	7	16.25	0.99
15	0.2	0.3	7	16.25	1
16	0.1	0.1	7	22.5	1
17	0.1	0.3	7	22.5	1
18	0.2	0.3	7	22.5	1

2757

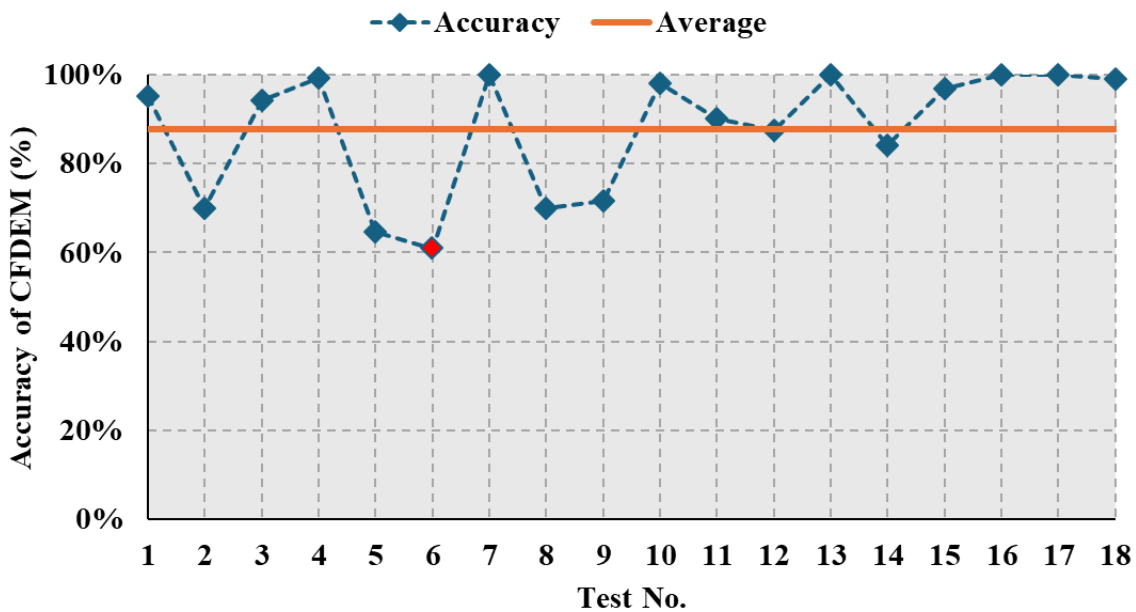
2758 **5.4.2 Numerical simulation results**

2759 To evaluate the capability of CFD-DEM model, digital replications of the scenarios tested in the
 2760 experimental work were developed. The predicted trap efficiency (TE) from the model was
 2761 compared with the lab measurements. The model's performance in terms of the accuracy was
 2762 reported using the relative accuracy (RA) parameter, which was defined by the formula below,
 2763 with the results presented in Figure 29:

$$RA = 1 - \frac{(|x_i - y_i|)}{|x_i|} \quad (10)$$

2764 where x_i and y_i are experimental and numerical TE values, respectively.

2765 The numerical model showed an accuracy of over 61% in predicting the TE values for the selected
 2766 test scenarios (shown in Table 20). The lowest accuracy was observed for Test 6, in which the
 2767 hydraulic parameters were set to $w = 0.2 \text{ m}$, $u_f = 0.3 \text{ m/s}$, $d = 3.5 \text{ cm}$, and $s = 16.25^\circ$. For this
 2768 test, the model overestimated the TE value (15.3%) during the numerical simulation compared to
 2769 the observed TE value (11%) in the laboratory (indicating approximately 61% accuracy).



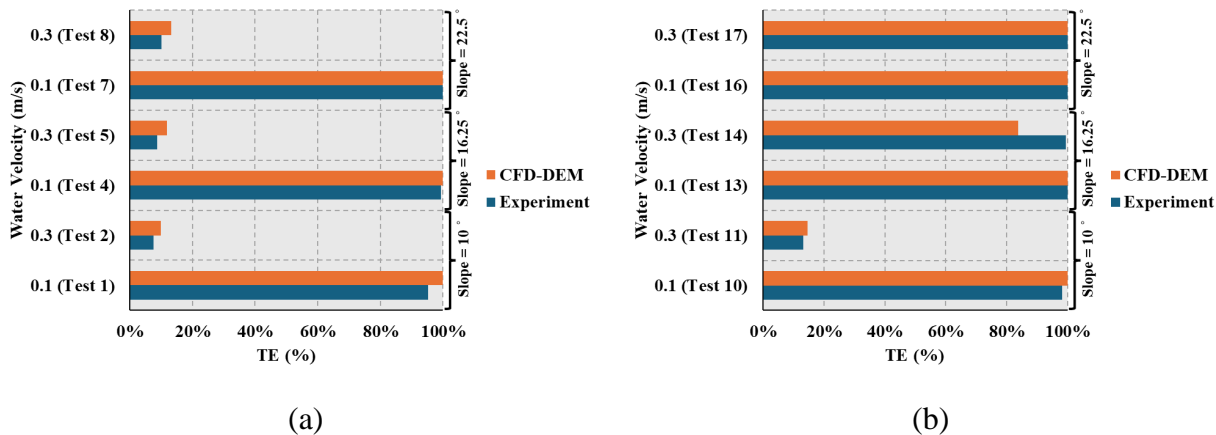
2770 **Figure 29.** The accuracy of CFD-DEM in simulating TE values based on experimental
 2771 observations

2772 To evaluate the impact of various parameters on particle behaviour and CFD-DEM performance
 2773 the simulated TE values were compared with experimental data, as shown in Figure 30, Figure 31,
 2774 and Figure 32. Each figure depicts the variation of a single parameter – bed deepening depth (d),
 2775 water velocity (u_f), and water level (w) – allowing for the isolated analysis of the individual effect
 2776 of different parameters on TE values. In each figure, the individual TE graphs were also grouped
 2777 for pairs of tests with the same bed slope within the FOV, allowing for a clearer comparison of
 2778 results under similar conditions. The cases with steeper slopes were associated with greater water
 2779 depth at the bed deepening, which resulted in slower velocities and increased deposition. These

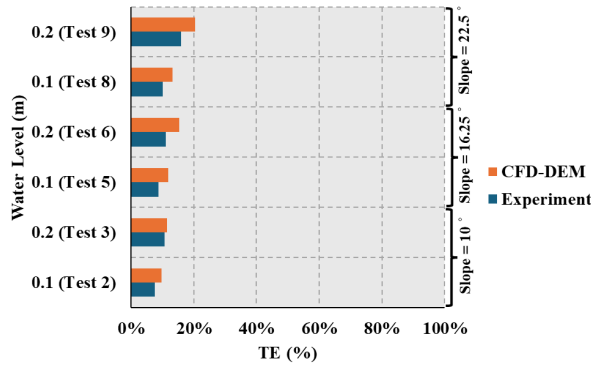
2780 trends were also captured by the CFD-DEM model as shown in Figure 30. In tests 2, 5, and 8,
 2781 tripling the flow velocity resulted in a significant reduction in TE values (by more than 80%)
 2782 compared to the corresponding tests 1, 4, and 7, as shown in Figure 30a. However, this effect was
 2783 not observed when the bed deepening depth was doubled in steep-slope scenarios, such as in tests
 2784 14 and 17 (Figure 30b).

2785 On the other hand, by providing results close to experimental observations, CFD-DEM was
 2786 reasonably capable of simulating tests in which the water level changed. As depicted in Figure 31,
 2787 the simulations captured the effect of the water level on TE values despite the existing error
 2788 between the numerical and experimental results. In this regard, the higher water depth leads to
 2789 lower particles velocity on the bed, resulting in larger deposition rate (i.e., TE). The TE values of
 2790 tests with doubled water depth showed a slight increase of less than 6%, regardless of the doubling
 2791 of d and the value of s (Figure 31a and Figure 31b).

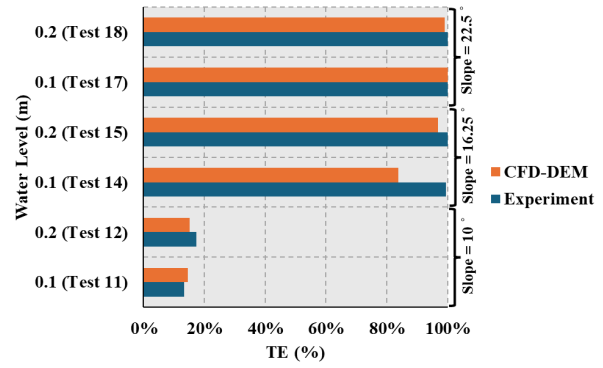
2792 The model also performed well in capturing the effect of changes in deepening of the channel bed
 2793 within the FOV area, as shown in Figure 32, when comparing the numerical and experimental
 2794 results in terms of increasing the depth of the channel bed. This leads to increase in the water depth
 2795 and results in higher values of TE. Doubling d led to a significant increase (over 80%) in TE values,
 2796 in steep-slope scenarios (Figure 32a and Figure 32c). However, this effect was not observed in
 2797 cases where MPs experienced low water depth and flow velocity (i.e., $w = 0.1\text{ m}$ and $u_f =$
 2798 0.1 m/s), as illustrated in Figure 32b.



2799 **Figure 30.** CFD-DEM results for tests in which the water velocity is varying while the other
 2800 parameters maintained constant: (a) tests in which $d = 3.5\text{ cm}$, (b) tests in which $d = 7\text{ cm}$

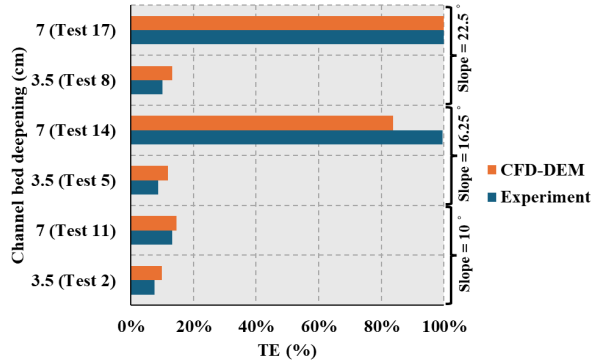


(a)

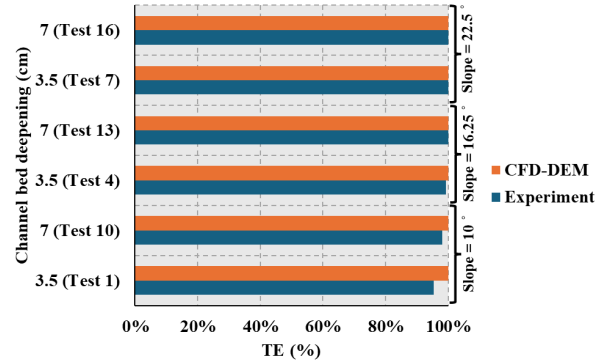


(b)

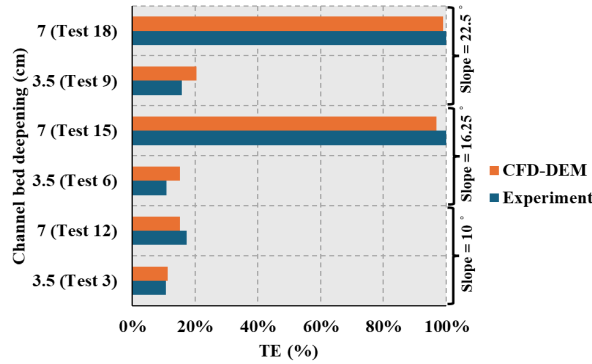
2801 **Figure 31.** CFD-DEM results for tests in which the water level is varying while the other
 2802 parameters maintained constant: (a) tests in which $d = 3.5$ cm, (b) tests in which $d = 7$ cm



(a)



(b)



(c)

2803 **Figure 32.** CFD-DEM results for tests in which the channel bed deepening is varying while the
 2804 other parameters maintained constant: (a) tests in which $w = 0.1$ m and $u_f = 0.3$ m/s, (b) tests
 2805 in which $w = 0.1$ m and $u_f = 0.1$ m/s, and (c) tests in which $w = 0.2$ m and $u_f = 0.3$ m/s

2806 To investigate the pattern of MPs deposition within the FOV in greater detail, tests 2, 5, and 8 were
2807 chosen for further analysis based on the results of the numerical model for 10 s of simulation, as
2808 depicted in Figure 33. The selected cases are completely similar to each other in terms of the
2809 hydraulic parameters except for the slope applied to the channel bed. In addition, they were among
2810 the cases in which CFD-DEM simulated MPs deposition and transport with the lowest accuracy
2811 (see Figure 29). The streamwise flow velocities in tests 2, 5, and 8, were simulated 0.37, 0.33, and
2812 0.35 m/s that were very closed to the experimental observations.

2813 In general, particles located near the center line within the FOV are transported faster compared
2814 to those on the sides due to the higher value of the flow velocity in the centerline compared to
2815 edges. Conversely, particles located close to the sides of the channel are more likely to deposit due
2816 to lower velocity in these regions, which is influenced by the wall effects. In addition, an increase
2817 in the slope results in a higher water depth and lower flow velocity (due to the continuity equation),
2818 leading to reduced particle velocity. Consequently, more MP particles become trapped, which
2819 aligns with the experimental results.

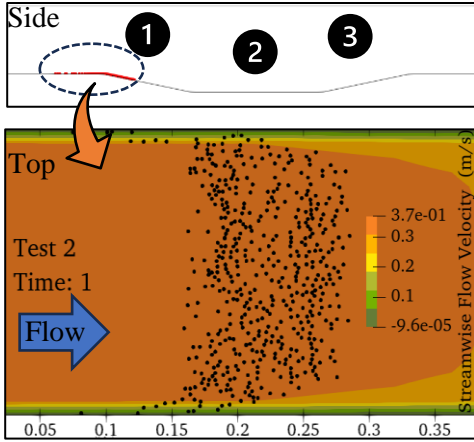
2820 In tests 2 and 5, the transport and deposition pattern of MP particles initially formed a semi-circular
2821 shape. Over time, this shape stretched out as most particles exited the FOV, leaving a few particles
2822 deposit along the channel sides. However, in test 8, MP particles showed a different transport and
2823 deposition pattern. They became trapped in circulations zone that pushed most particles further
2824 downstream, with a few being directed towards the sides of the channel. It can be concluded that
2825 a few plastic particles become trapped within the FOV over the time, and this is because of the
2826 effect of the sliding friction applied to MPs from the walls on the sides.

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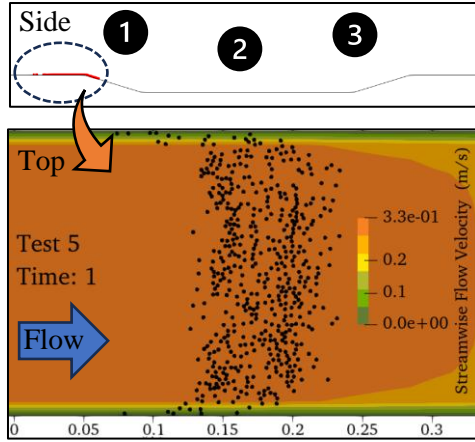
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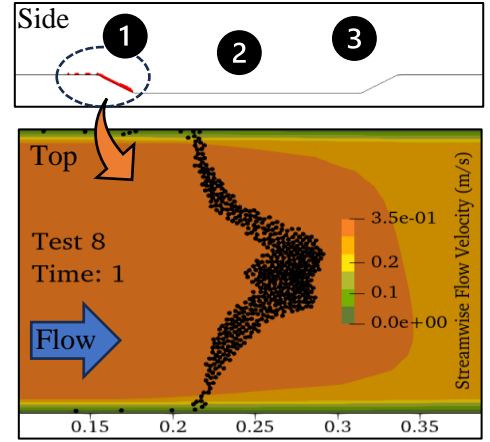
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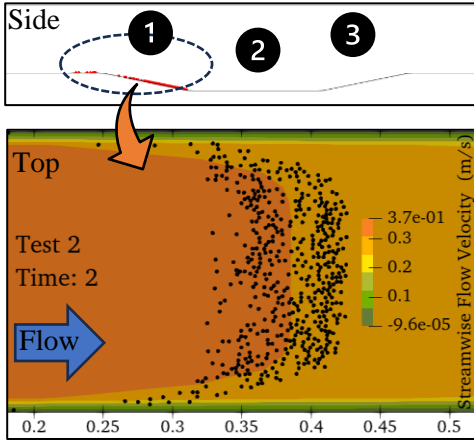
Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



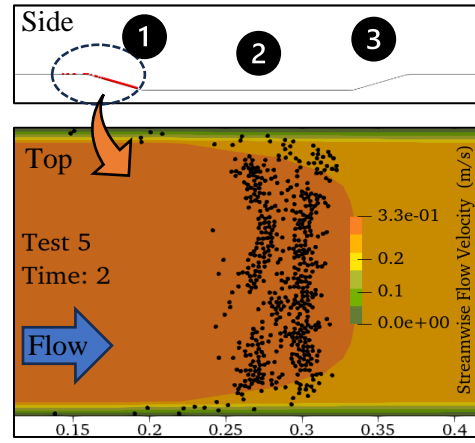
Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



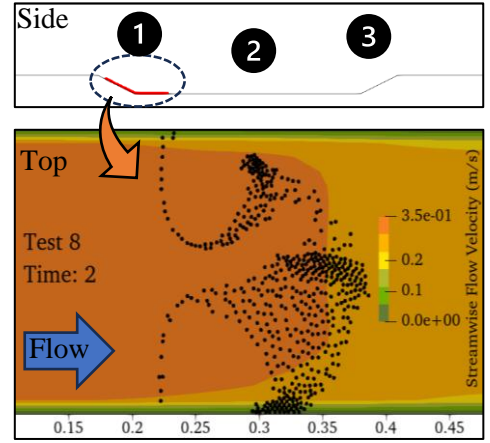
Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



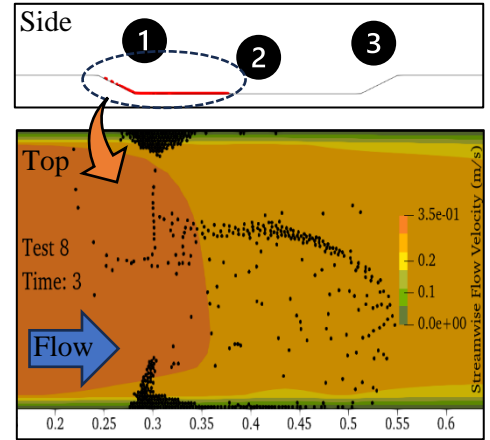
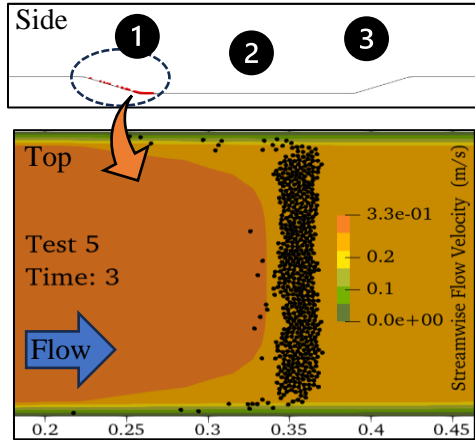
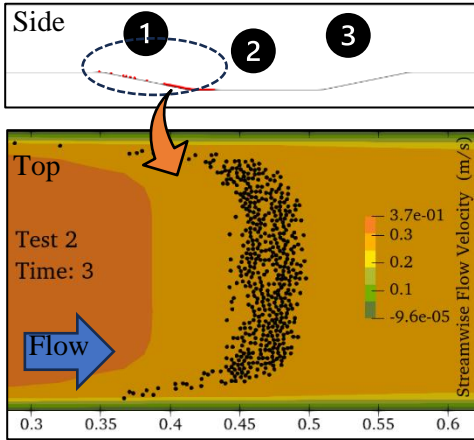
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Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



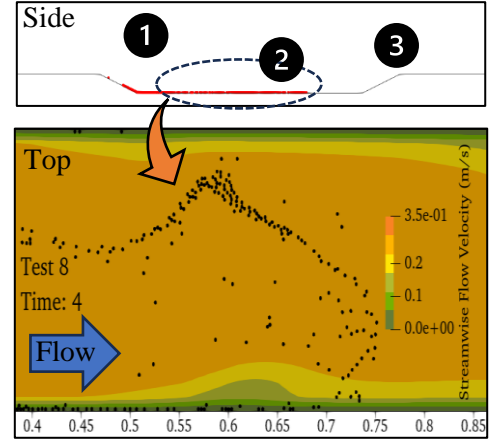
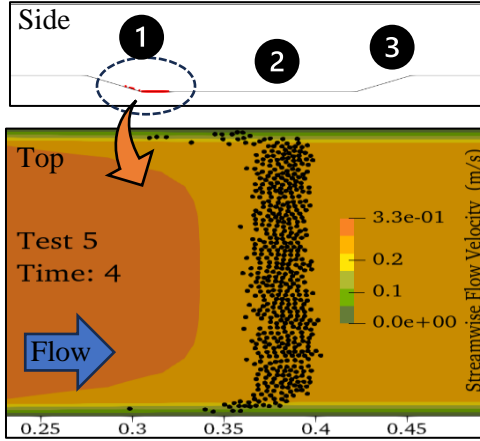
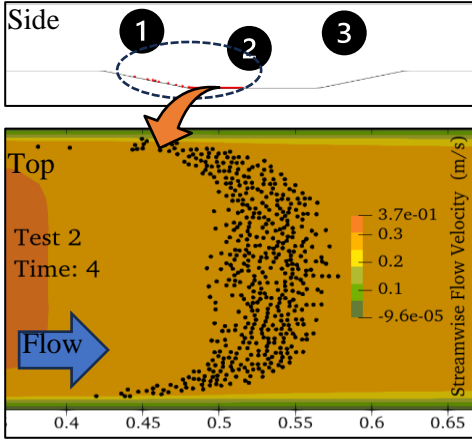
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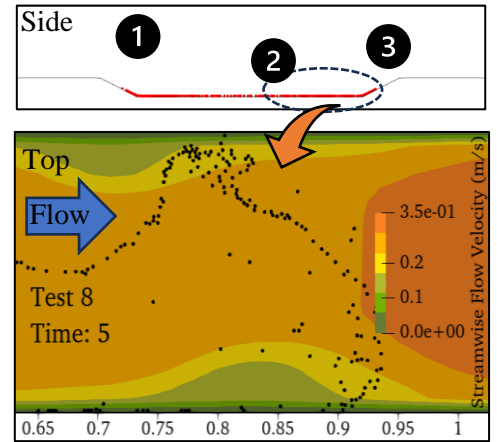
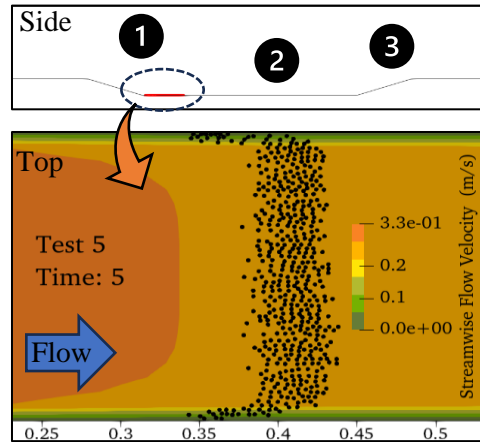
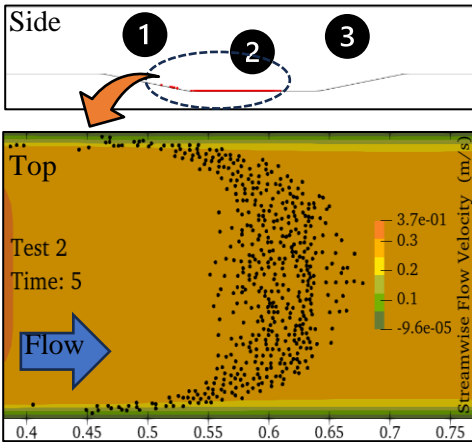
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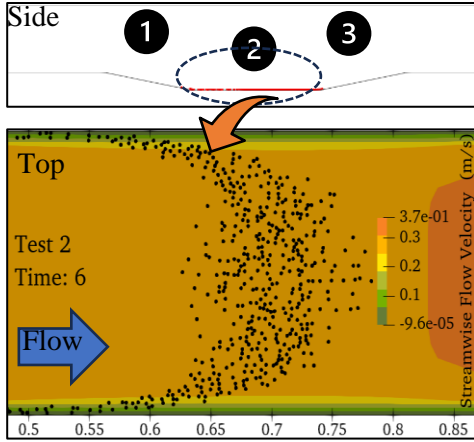
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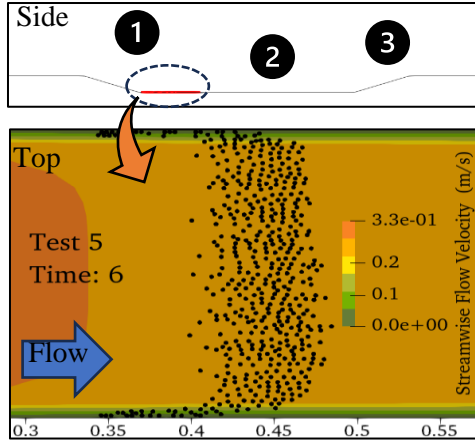
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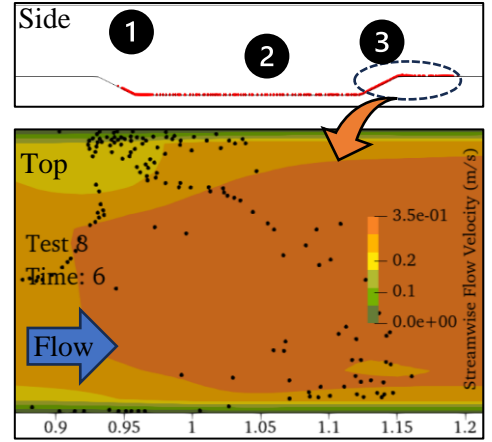
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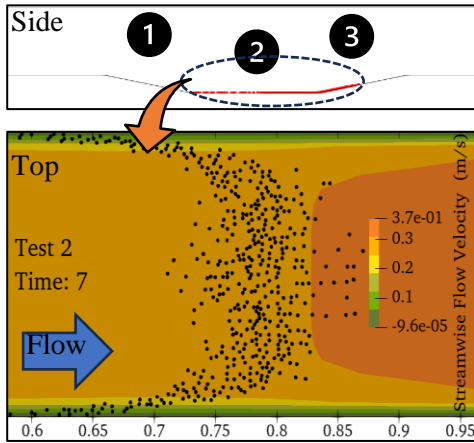
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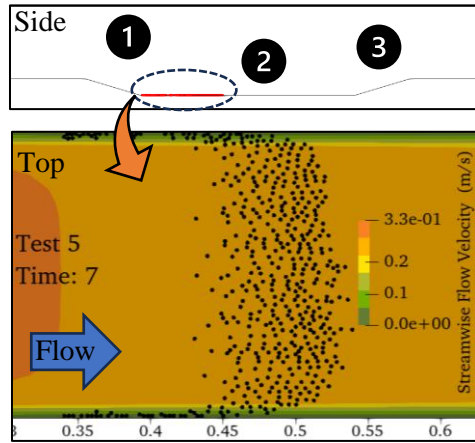
Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



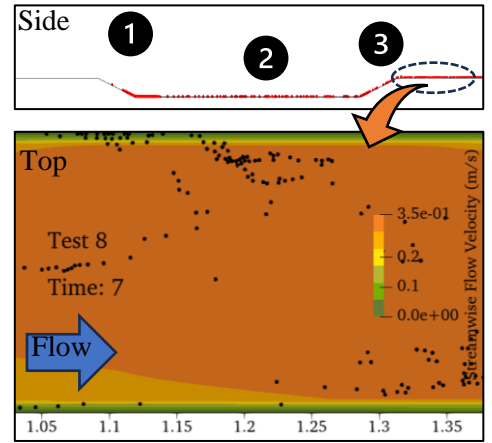
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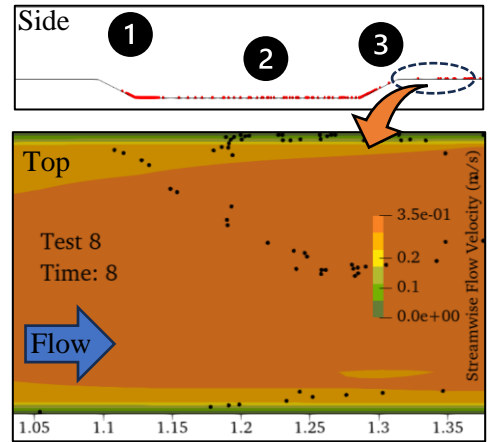
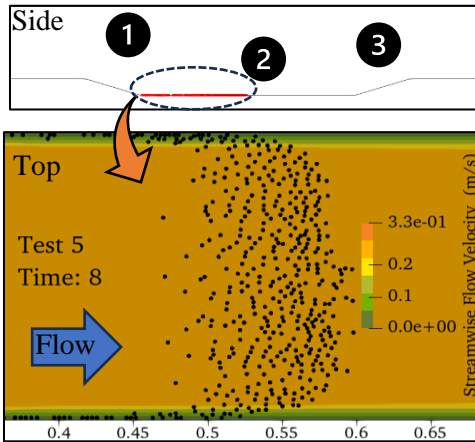
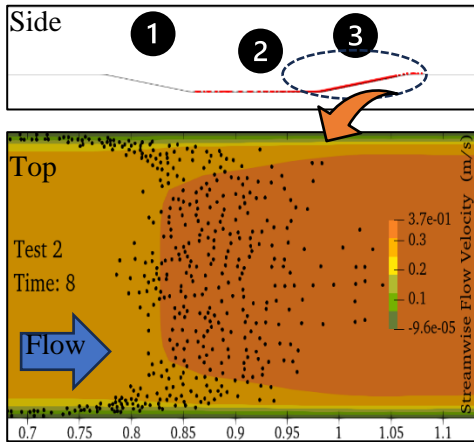
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Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)



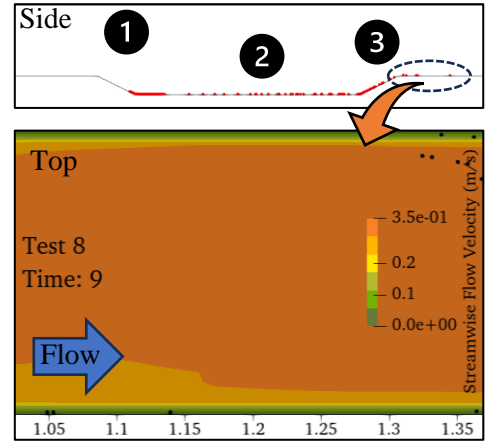
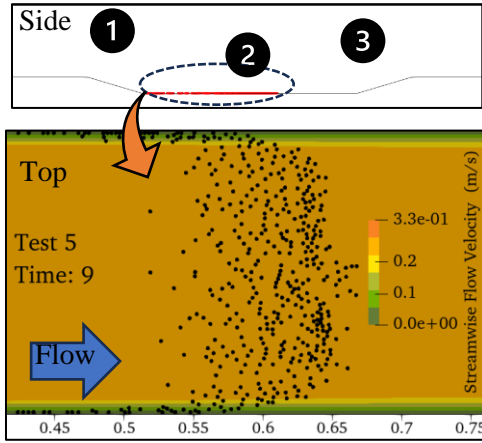
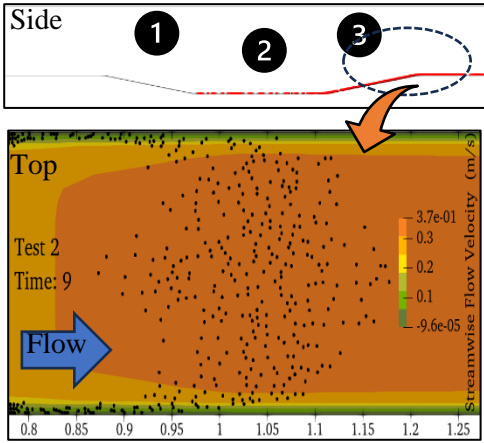
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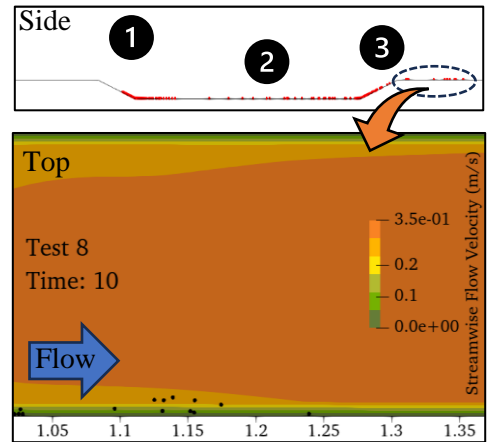
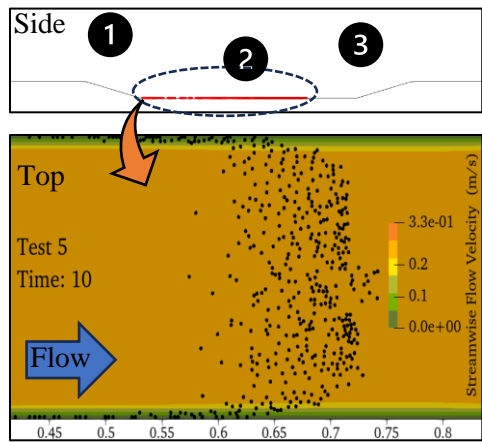
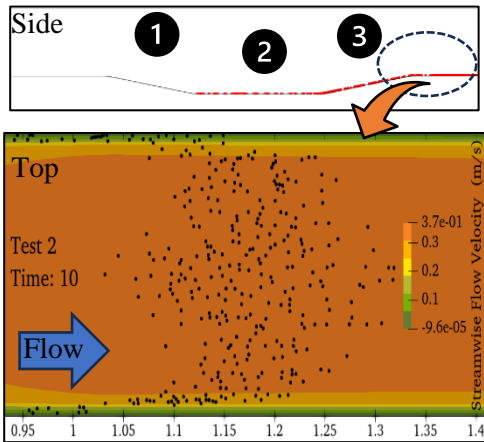
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Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)

Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)

Longitudinal distance from the injection point (at $x = 0$) in x -direction (m)

(a)

(b)

(c)

2831 **Figure 33.** Typical pattern of spherical MPs deposition within the channel around FOV ($x = 0.2 -$
 2832 1 m) from both top and side views based on CFD-DEM simulation for (a) Test 2, (b) Test 5, and

2833 (c) Test 8. The locations 1, 2, and 3 show the first slope, the flat bed, and the second slope applied
2834 to the channel bed, respectively.

2835 MPs particles exhibited circulation, particularly in test 8 (see Figure 33c), due to the influence of
2836 the increased slope of the channel bed, which induced significant flow circulation. To investigate
2837 this further, the turbulent kinetic energy (k) contour lines for all three tests were depicted from the
2838 side view (Figure 34) at time-step equal to two when the first circulation was observed (in Test 8).
2839 It can be seen that the maximum value of k in Test 8 was almost doubled the k value in tests 2 and
2840 5 at the same time across the channel bed including sloped and flat areas. Turbulent kinetic energy
2841 rises in areas with high local velocity gradients. This means that these areas are sources of turbulent
2842 kinetic energy where local velocity gradients are crucial (Pilechi et al., 2011). In this regard, it can
2843 be concluded that by increasing the slope of the channel bed, higher local velocities were generated
2844 resulting in larger k values and more intense circulations. As shown in Fig. 10, in all three tests,
2845 the maximum turbulent kinetic energy was reduced through diffusion and dissipation.

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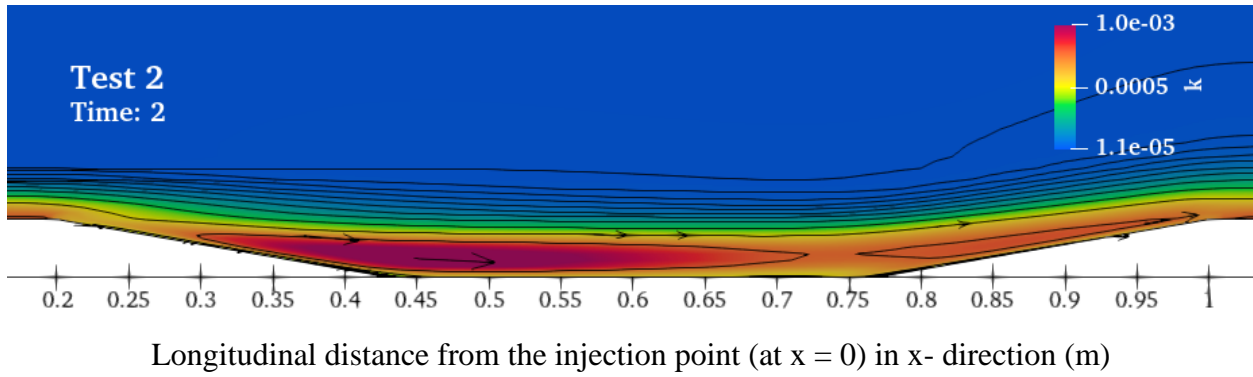
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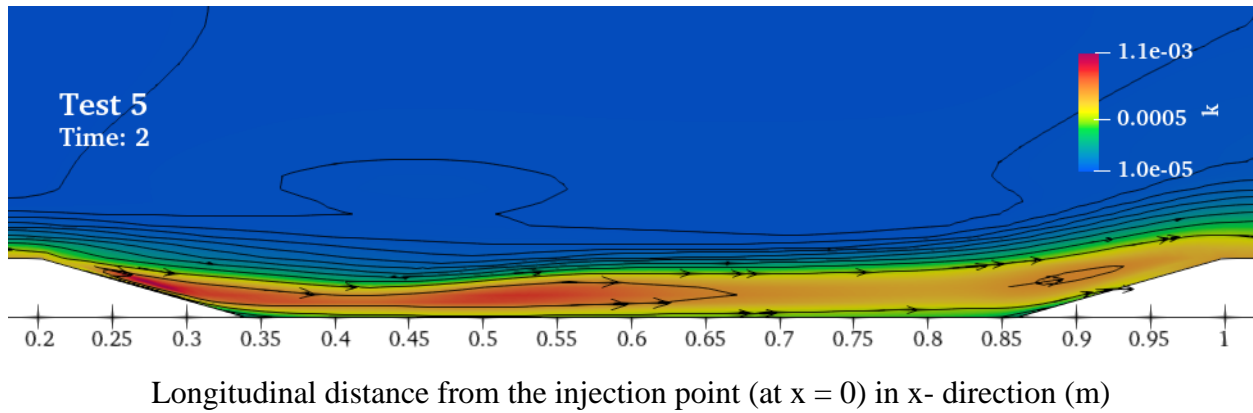
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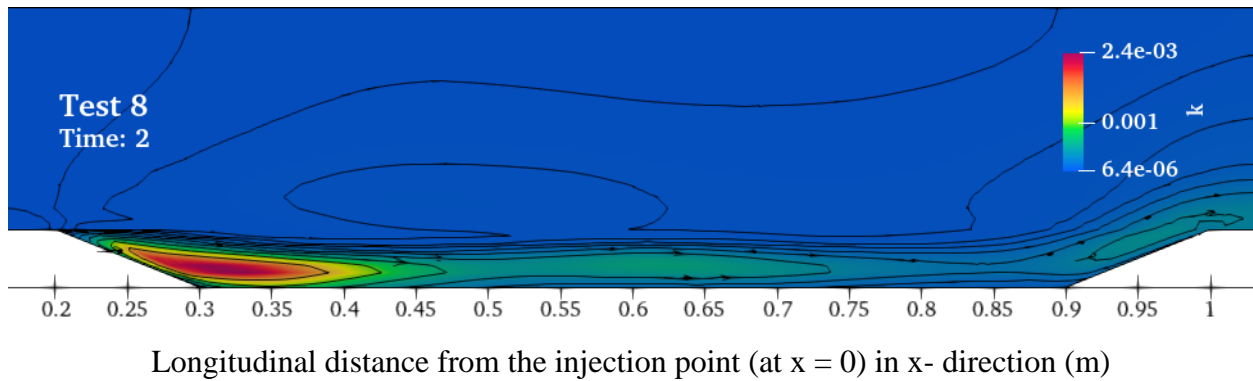
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(a)



(b)



(c)

2855 **Figure 34.** Longitudinal cross-section contours of k within FOV ($x = 0.2 - 1$ m) at time = 2 s,
 2856 from side view based on CFD-DEM simulation for (a) Test 2, (b) Test 5, and (c) Test 8.

2857 **5.4.3 Sensitivity analysis**

2858 To further evaluate the model's performance, a sensitivity analysis was conducted, as described in
 2859 section 5.3.2.3, on the simulation results with respect to various parameters identified as potential
 2860 impactful as shown in Table 21. The first parameter examined was the number of particles used is

2861 the simulation. No significant difference in TE values were observed between simulations using
2862 500 and 1000 particles. However, the TE values differed considerably when only 250 particles
2863 were used. This indicates that considering a very small number of particles compared to the number
2864 used in experiments or present in real conditions may lead to inaccurate results regarding the
2865 number of MPs accumulated (or deposited). On the other hand, considering the same number of
2866 MPs (used in experiments) in simulations highly increases the computational efforts without
2867 significantly improving result accuracy. Based on these findings, 500 particles were selected for
2868 the CFD-DEM coupling simulations across different test scenarios.

2869 No significant difference was observed in the simulated TE for CFD mesh grid sizes smaller than
2870 10 mm. Therefore, a grid resolution of 10 mm was chosen for the CFD mesh in all directions of
2871 simulations in this study, to balance accuracy and computational efficiency. It is important to note
2872 that the reported *RA* values are the best values obtained based on the measured and simulated TE
2873 values using the fine mesh.

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2875 **Table 21.** The performance (i.e., *RA* values) of the model in simulating MPs for Test 2 based on different sensitivity analysis
 2876 parameters

Relative Accuracy	Number of particles			Mesh size			Friction coefficient (f_s, f_r)				Coupling interval		DEM time-step (s)	
	250	500	1000	Fine	Medium	Coarse	(0.1,0.05)	(0.05,0.1)	(0.05,0.05)	(0.01,0)	10	100	1e-5	2e-5
<i>RA</i>	0.56	0.70	0.72	0.70	0.65	0.64	0.23	0.19	0.51	0.70	0.70	0.70	0.70	0.70

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Among the various numerical parameters for which sensitivity analysis was performed, the friction coefficient had the most significant impact on TE values across all 18 tested scenarios. Therefore, a detailed sensitivity analysis on the combination of sliding and rolling friction coefficients was conducted for all tests, as described in section 5.3.2.3. This analysis revealed that rolling friction (f_r) has a negligible effect compared to sliding friction (f_s). The combination of sliding friction $f_s = 0.01$ and rolling friction $f_r = 0$ showed the closest prediction to the experimental results.

Also, two coupling intervals of 10 and 100 with two DEM time-steps (1×10^{-5} s and 2×10^{-5} s) were analysed to assess their impact on computational efficiency and numerical stability. The coupling interval refers to the number of DEM time-steps between exchanges between the CFD and DEM models. Most simulations were successfully performed with a coupling interval of 100 and a DEM time-step of 2×10^{-5} s, which were more efficient in terms of computational cost. However, the model became unstable with this configuration in tests 4, 7, 13, and 16, which led to the use of smaller coupling intervals and DEM time-steps for these test cases. Despite these adjustments, no significant impact on the number of deposited MPs was observed with varying coupling intervals and DEM time-steps.

5.5 Conclusion

This study utilized a two-phase model combining fluid flow and spherical MPs dynamics, integrating CFD and DEM simulations. The study assessed the capability of a CFD-DEM model by creating digital replicas of experimental scenarios and comparing the model's predicted TE with laboratory measurements. Model accuracy was evaluated using the relative accuracy (RA) parameter. The numerical approach employed a four-way coupling method, incorporating fluid-particle interactions and particle-particle collisions, ensuring a two-way simulation for both dynamics. Fluid flow was modeled with OpenFOAM® using the PISO solver, while particle motion was simulated with Aspherix®. The CFD-DEM simulations were implemented using computational domains designed to match the dimensions, initial conditions, and boundary conditions of the physical model setups for comparison with experimental observations. Overall, the CFD-DEM model demonstrated an accuracy of over 61% in predicting TE values across the test scenarios.

The test with steepest slope (i.e., test 8) showed significant particle circulation due to increased bed slope. It was observed that the turbulent kinetic energy (k) value in this test was the largest compared to the k value obtained in the other tests. Therefore, it can be concluded that the larger

circulations occurred in steeper slope tests with higher k values. The steeper slope also resulted in greater water depth and lower flow velocity, causing reduced particle velocity and increased trapping of MPs, consistent with experimental observations. In milder slope tests (i.e., tests 2 and 5), MPs particles initially formed a semi-circular pattern, which gradually stretched out as most particles exited the FOV, leaving some deposited along the channel sides. However, in test 8, the particles behaved differently, becoming trapped in circulation zones that pushed them downstream, with some directed towards the channel sides. Over time, a few particles remained within the FOV due to the sliding friction from the channel walls.

A sensitivity analysis was conducted to evaluate the model's performance by examining the impact of various parameters on simulation results. Although the number of particles used in the simulation was found to significantly affect the results, the friction coefficients had the highest impact on TE values across all simulated test scenarios.

Given the findings provided in this study and to investigate more in-depth the capability of CFD-DEM in simulating MPs transport and deposition in water bodies, it is recommended to consider the following suggestions in future studies. First, it could be of great importance to include MPs particles with different physical properties (i.e., size, density, and shape) in numerical simulations performed using CFD-DEM. Second, it is crucial to include various types of water bodies with different hydraulic and geometric conditions in the numerical modeling of MPs transport and deposition using CFD-DEM, where a comprehensive physical modelling is required. Third, it is necessary to further investigate the effects of each parameter selected, through the sensitivity analysis, in this study on transport and deposition of MPs to obtain a better understanding of their impact on the accuracy of simulations.

5.6 References

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6 Summary, Concluding Remarks, and Recommendations

6.1 Summary and Concluding Remarks

The water flow characteristics such as velocity and water depth, as well as physical properties of particles (including density, size, and shape), and bathymetric changes in water bodies affect the deposition patterns of marine MPs. As stated above, the main objective of this thesis is to utilize both physical and numerical modeling techniques to shed light on how spherical and cylindrical MPs are deposited or transported in water bodies, particularly considering changes in water bodies' bathymetry. Physical modeling provides not only valuable results regarding the sensitivity of MPs transport and deposition to different hydraulic and bathymetric conditions of water bodies but is also considered a reliable reference for numerical results. To achieve this goal, a series of experiments were carried out by employing a water body with a generic geometry to investigate the influence of various factors, such as the initial velocity of particles and depth of turbulent flows, changes in bathymetry, and shape of the MPs particles on their transport and deposition. This research was primarily focused on the use of two numerical approaches, the coupled CFD-DEM and ELM methods, to better understand and predict the transport and deposition of marine MPs in aquatic environments.

There are some concluding remarks obtained during this research that can be summarized as below:

- 1- This research was strongly founded on the review of four reliable lagrangian PTM models (i.e., D-WAQ PART, Ichthyop, TrackMPD, and CaMPSim-3D) to deliver a clear picture of the existing models in simulating the fate and transport of MPs. The comparison highlighted variations in how these models account for physical transport and transformation processes, as well as the effects of MPs properties such as size, shape, and density. Key findings indicated that TrackMPD and CaMPSim-3D incorporate more accurate equations for settling velocity, considering particle shape, whereas D-WAQ PART and Ichthyop use simplified formulations. Regarding transformation processes, only TrackMPD and CaMPSim-3D can simulate biofouling and degradation. While all models implement advection-diffusion equations, their approaches to parameterizing physical processes differ, with D-WAQ PART, TrackMPD, and CaMPSim-3D considering wind-

induced drift. In terms of mesh structure, CaMPSim-3D uniquely supports unstructured meshes, whereas the others rely on structured mesh systems. Despite these advancements, limitations remain. None of the reviewed models simulate homo- and hetero-aggregation, agglomeration, photodegradation, or chemical and biological degradation—key processes affecting MPs' long-term fate.

- 2- This study highlighted the importance of understanding how different shapes of MPs are transported and deposited in channels with generic geometry. The findings demonstrated that flow velocity and bathymetric changes significantly influence MPs behavior, with cylindrical particles exhibiting higher trapping tendencies than spherical ones. The dataset generated in this study served as a valuable resource for calibrating and validating numerical models, particularly CFD simulations, ensuring their reliability in cases where experimental data is unavailable.
- 3- This study highlighted the effectiveness of ML, specifically the ELM model, in predicting MPs deposition under various experimental conditions tested during physical modeling. The results confirmed that the best-performing model, including all dimensionless inputs that consisted of flow characteristics and the channel specifications, demonstrated high accuracy for both spherical and cylindrical datasets. The uncertainty analysis further validated the reliability of the model, with narrow uncertainty bands and small prediction error intervals.
- 4- This research successfully integrated CFD and DEM simulations to model fluid flow and spherical MPs dynamics, demonstrating the effectiveness of the CFD-DEM approach in replicating experimental conditions. The model achieved an accuracy of over 61% in predicting trapping tendencies of MPs, with results indicating that steeper slopes led to increased particle circulation, higher turbulent kinetic energy, and greater MPs trapping. Sensitivity analysis revealed that friction coefficients had the most significant influence on MPs trapping.

6.2 Recommendations for Future Studies

To enhance our understanding of MPs behavior and their environmental impact, further research should focus on either refining or employing the other numerical models to capture complex transport and transformation mechanisms with greater accuracy. Future studies should address critical gaps such as investigation of transformational processes, which impact MPs transport and

deposition, to develop a more comprehensive and reliable modeling framework for MPs pollution in aquatic systems. Limitations such as the consideration of only specific MP shapes, sizes, and densities, as well as the exclusion of varied sediment types, indicate the need for further research. Future studies should explore a broader range of MP characteristics and environmental conditions to enhance the understanding of MP transport and deposition dynamics. This research establishes a solid foundation for the future application of ML models in predicting MPs deposition, emphasizing their potential for broader environmental studies. It is important to note that to improve the understanding and accuracy of CFD-DEM simulations in MP transport and deposition, future studies should consider incorporating MPs with varying physical properties, modeling different hydraulic and geometric conditions, and further investigating the impact of key parameters identified during this research.