

Morphological and Numerical Modeling of a Highly Dynamic Tidal Inlet at Shippagan Gully, New Brunswick

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My fascination with coastal engineering began in 2008 upon being introduced to the field in my final year of an undergraduate degree at Queen's University. Perhaps growing up a few short kilometres from the Atlantic Ocean played a role, but something about the sheer power and magnitude of our oceans and seas intrigued me. Being a fairly creative minded person, I was also attracted by the fact that no design code exists for coastal engineers. Instead, every project is unique, and therefore requires an equally unique and creative solution.

In response to my fascination, I quickly sought ways in which I could pursue a graduate degree in coastal engineering. After a series of set-backs and rather unfortunate events, I was extremely lucky to be led to Dr. Ioan Nistor and the University of Ottawa. As a graduate student, one could not ask for a better supervisor than Dr. Ioan Nistor. His level of care and interest in each and every one of his student's endeavours is un-surpassed by anyone I have ever met in academia. For that, I owe him many thanks.

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ABSTRACT

Shippagan Gully is a highly engineered, tidal inlet located near Shippagan, New Brunswick, on the Gulf of St. Lawrence. It is a particularly complex tidal inlet due to the fact that its tidal lagoon transects the Acadian Peninsula and is open to the Bay des Chaleurs at its opposite end. As such, two open boundaries with phase lagged tidal cycles drive the flow through the inlet, alternating direction with each tide and reaching velocities which exceed 2 m/s. Over the past century, despite various engineered interventions, shipping activities through the inlet have been threatened due to the increasing encroachment from sediment which has accumulated on the east side of the navigation channel and immediately offshore from the inlet. In addition to these highly asymmetrical sediment deposition patterns, severe down-drift erosion suggests the presence of a prominent westward net longshore transport, further complicating the coastal morphology at the inlet.

Due to the overwhelming requirement for constant maintenance dredging and the long-term degradation of existing coastal structures, a numerical model study of Shippagan Gully has been undertaken in order to identify principal morphology mechanisms and to provide guidance for future coastal works. The numerical models CMS Wave and CMS Flow (USACE) have been applied in a coupled manner to simulate the hydrodynamics, coastal processes and morphology changes at the inlet due to the combined effects of waves and tides. Once calibrated to historical morphologic evolution, the numerical modeling system was used as an analytical tool to identify and understand the hydrodynamic processes responsible for the morphological changes observed at Shippagan Gully. The numerical model was further used to predict future morphologic trends, in a qualitative manner, for both the *status-quo* and a number of engineered alternatives proposed by the author. The methodology and results of this study are presented herein.

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LIST OF SYMBOLS

A	= empirical coefficient typically ranging from 0.1 to 2 (non-dimensional)
A_{cr}	= empirical coefficient (non-dimensional)
a	= empirical coefficient (non-dimensional)
a_i	= amplitude of tidal constituent (m)
b	= sediment transport coefficient dependent on Reynolds number (non-dimensional) = distance between adjacent wave rays (m) = empirical coefficient (non-dimensional)
C	= wave celerity (m/s) = depth-averaged sediment concentration (volume of sediment over total volume) = Chezy friction coefficient (non-dimensional)
C_b	= empirical bottom-stress (friction) coefficient
C_d	= wind drag coefficient (non-dimensional)
C_ε	= coefficient determined from an empirical curve (non-dimensional)
C_x	= characteristic velocity in wave-action density equation (m/s)
C_y	= characteristic velocity in wave-action density equation (m/s)
C_θ	= characteristic velocity in wave-action density equation (m/s)
c_f	= Darcy-Weisbach friction coefficient (non-dimensional)
c_R	= reference concentration in Lund-CIRP formula for suspended load (kg/m ³)
D	= characteristic particle diameter (m) = sediment deposition rate (downward sediment flux) (kg/s)
D_x	= diffusion coefficient for the x direction (non-dimensional)
D_y	= diffusion coefficient for the y direction (non-dimensional)
d	= water depth (m)
d_b	= water depth at wave breaking (m)
d_{50}	= mean grain diameter (m)
d_*	= non-dimensionalized grain size (non-dimensional)
E_{total}	= total energy density (J/m ²)
f	= Coriolis coefficient (rad/s)
g	= gravitational acceleration (m/s ²)
H	= nearshore wave height (m)
H_b	= breaking wave height (m)
H_o	= offshore wave height (m)
H_{sb}	= significant breaking wave height (m)
h	= total water depth (m)
K_r	= refraction coefficient (non-dimensional)
K_s	= shoaling coefficient (non-dimensional)

K_x	= sediment diffusion coefficient for the x-direction (non-dimensional)
K_y	= sediment diffusion coefficient for the y-direction (non-dimensional)
k	= wave number (m^{-1})
L	= wavelength (m)
MWL	= mean water level (m)
MSL	= mean sea level (m)
m_b	= beach slope (m/m)
N	= wave-action density ($\text{J}\cdot\text{s}/\text{m}^2$)
n	= sediment porosity (non-dimensional) = Manning's roughness coefficient (non-dimensional)
P	= sediment pick-up rate (upward sediment flux) (kg/s) = material placed in sediment budget cell (m^3)
p	= sediment porosity (non-dimensional)
Q_k	= potential longshore sediment transport (Kamphuis formula) (m^3/yr)
Q_s	= longshore sediment transport rate (CERC formula) (m^3/yr)
Q_{sink}	= sink term to a sediment budget (m^3)
Q_{source}	= source term to a sediment budget (m^3)
q_{bw}	= bed load parallel to wave direction ($\text{m}^3/\text{m}\cdot\text{s}$)
q_{bn}	= bed load normal to the wave direction ($\text{m}^3/\text{m}\cdot\text{s}$)
q_{sb}	= bed load rate ($\text{m}^3/\text{m}\cdot\text{s}$)
$\bar{q}_{sb}$	= gravimetric specific bed-load rate ($\text{N}/\text{m}\cdot\text{s}$)
q_x	= flow per unit width parallel to the x-axis (m^2/s)
q_y	= flow per unit width parallel to the y-axis (m^2/s)
q_{tot}	= total load (both suspended and bed load) ($\text{m}^3/\text{m}\cdot\text{s}$)
R	= material removed from the sediment budget cell (m) = hydraulic radius (m)
S_{xx}	= wave-driven radiation stress (N/m^2)
S_{xy}	= wave-driven radiation stress (N/m^2)
S_{yy}	= wave-driven radiation stress (N/m^2)
SWL	= still water level (m)
T	= wave period

V_*	= shear velocity (m/s)
v	= depth-averaged current velocity parallel to the y-axis (m/s)
W	= wind speed (m/s)
w	= density of bed material divided by density of water (non-dimensional)
Y	= Y-axis in Shields curve (non-dimensional)
Y_{CR}	= critical parameter determined from the Y-axis of Shields curve (non-dimensional)
z	= water elevation relative to the bed (m)
α	= wave crest angle relative to the shoreline (0)
α_b	= breaking wave angle relative to the shoreline (0)
α_i	= phase angle of tidal constituent (0 or rad)
γ_s	= specific gravity (N/m ³)
ΔV	= net change in volume within the cell (m ³)
κ	= wave diffraction term (non-dimensional)
η	= water surface elevation relative to still water level (m)
	= relative water depth (z/h)
ρ	= fluid density (kg/m ³)
ρ_a	= density of air (kg/m ³)
ρ_s	= sediment density (kg/m ³)
ρ_w	= density of seawater (kg/m ³)
θ	= wind direction (0)
$\theta_{cw,m}$	= mean Shields parameter for combined waves and currents (non-dimensional)
θ_{cw}	= maximum Shields parameter for combined waves and currents (non-dimensional)
τ	= shear stress (N/m ²)
τ_b^*	= dimensionless shear stress (non-dimensional)
τ_{bmax}^*	= maximum shear stress at the bed (non-dimensional)
τ_{cr}	= shear stress at incipient sediment motion (N/m ²)
τ_c^*	= dimensionless critical shear stress (non-dimensional)
τ_{bx}	= bottom stress parallel to the x-axis (N/m ²)
τ_{by}	= bottom stress parallel to the y-axis (N/m ²)
τ_{wx}	= surface stress parallel to the x-axis (N/m ²)
τ_{wy}	= surface stress parallel to the y-axis (N/m ²)
τ_{Sx}	= wave stress parallel to the x-axis (N/m ²)
τ_{Sy}	= wave stress parallel to the y-axis (N/m

1.0 INTRODUCTION

Shippagan Gully is a coastal inlet located on the Gulf of St. Lawrence near Shippagan, New Brunswick. The inlet marks the south-eastern limit of a natural waterway which transects the Acadian Peninsula, providing an alternate route between the Bay de Chaleurs and the Gulf of St. Lawrence (refer to Figure 1.1). The inlet, which is shown in Figure 1.2, has been maintained using man-made coastal structures since the late 1800's as its navigability has proven crucial to the local economy, which is influenced in large part by the fishing industry. Since the initial construction of two 300 metre jetties in the 1880's, sediment accumulation both within and immediately offshore from the inlet has been observed. As such, a number of structural alterations, rehabilitations and dredging activities were undertaken throughout the 20th century in an attempt to stabilize the inlet and promote navigable depths. Regardless of these human interventions however, sediment has continued to deposit both within Shippagan Gully and offshore, greatly limiting its use as a navigable passage between the Bay de Chaleurs and the Gulf of St. Lawrence.

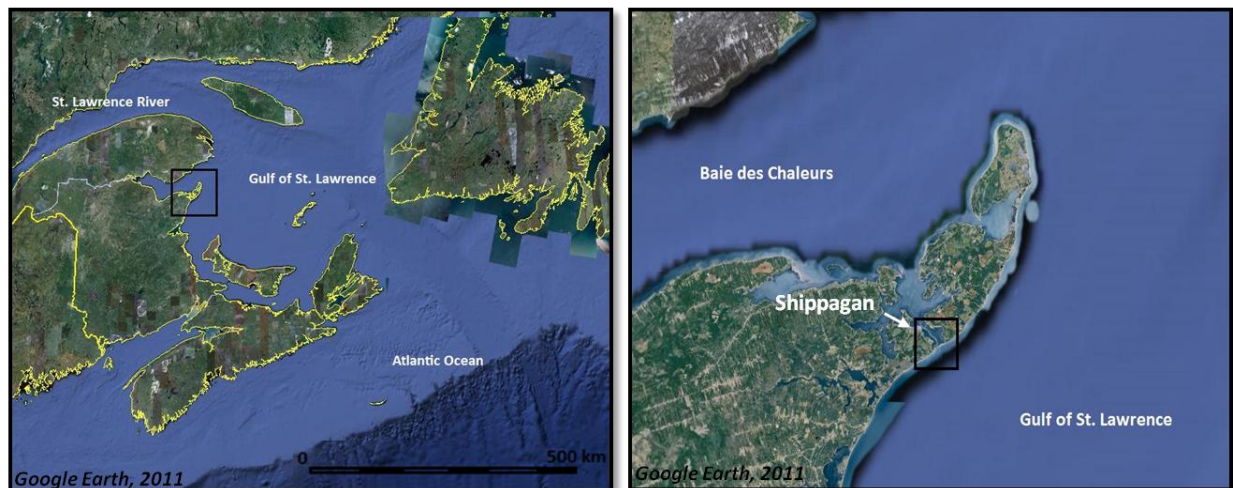


Figure 1.1: Satellite imagery showing the location of Shippagan Gully within Atlantic Canada (left) and the Acadian Peninsula (right) (Google Earth, 2011).



Figure 1.2: Aerial photograph of Shippagan Gully taken in 2007 (PWGSC, 2010).

In the absence of human intervention in recent years, natural morphological changes have taken hold of Shippagan Gully. Safe navigation through the inlet has become increasingly limited due to the continual growth and migration of morphologic features at the site. As such, the Canadian Hydraulics Centre (CHC) of the National Research Council of Canada (NRC) was retained by Public Works & Government Services Canada (PWGSC) to investigate the hydrodynamic and sedimentary processes at Shippagan Gully. This work was performed principally by the present author and in partnership with the University of Ottawa, thereby forming the basis of this thesis.

1.1 Significance of the Study

The contributions of the present thesis are significant in both a practical and technical sense. The practical significance of this study stems from the fact that Shippagan Gully is an important navigation channel which has a long history of being poorly understood and subsequently underutilized. Past coastal works at the site have proven unsuccessful in providing

maintainable and navigable depths through the inlet, thereby resulting in great losses to the local economy and community. By developing a proper understanding of the hydrodynamic and sedimentary processes present at the inlet, guidance can be provided from which future decisions regarding coastal works and sediment management can be confidently made. In a practical sense, this work represents a significant contribution to the development of a safe and navigable water way for future generations.

In a technical sense, this thesis is considered to be both significant and novel for a number of important reasons. First, numerical modeling studies of tidal inlet morphology are not a common occurrence in Canada. USA has begun to recognize the importance of such work and has subsequently formed a division within the United States Army Corps of Engineers (USACE) to study such topics, aptly named the Coastal Inlets Research Program (or CIRP). Canada however does not have such a task force to monitor and study the various coastal inlets which scatter our coastlines. As such, a Canadian study regarding tidal inlet morphology is novel, simply because it's Canadian.

A second point which makes the present study novel is the relative complexity of the morphological features present at Shippagan Gully. As is discussed later in this thesis, many morphological characteristics at Shippagan Gully do not fit the typical description of a 'normal' tidal inlet, nor are their dominant processes so easily identified. This is in large part due to the long history of engineering activities at Shippagan Gully, many of which have likely left the site worse off than it would otherwise have been. These historical human interventions further complicate the situation at Shippagan Gully, making it a particularly difficult inlet to fully understand.

Lastly, the most defining feature of Shippagan Gully is that its extensive tidal lagoon transects a peninsula, and is therefore open to tidal forcings from both sides. Due to this fact, much of our conventional knowledge regarding tidal inlet hydrodynamics is no longer applicable, as it is the phase-lagged differential between two water levels which governs the tidal hydrodynamics at the site. As a result, the ebb and flood tides may no longer be approximately equal in length or magnitude, the timing of maximum and minimum flows are likely to be drastically altered, and

the relationships which are built on the concept of a tidal prism are no longer valid. These topics are discussed in detail herein. As a result, Shippagan Gully cannot be studied, nor can it be properly understood, based solely upon experience and knowledge of previous tidal inlet works. Shippagan Gully is an example of an inlet which is not typical and contains key features which have never before been extensively studied, nor understood. As such, much of the findings presented herein are the first of their kind, making this study a truly novel one.

1.2 Study Objective and General Approach

The main objective of the present study is to provide a scientific basis to inform and guide decisions concerning the maintenance and future development of the existing infrastructure and navigation channel at Shippagan Gully. Under this main objective, the following sub-objectives were identified:

- Identify the governing hydrodynamic processes responsible for the complex morphological features present at Shippagan Gully;
- Predict future morphological evolution if the status quo is maintained;
- Study and understand the hydrodynamics and sediment transport for alternative arrangements of structures; and
- Predict future morphological evolution for alternative structural arrangements;

A numerical modeling approach was taken in order to fulfil the requirements and general objectives of the study. Two numerical models, CMS-Wave (Lin *et al.*, 2008) and CMS-Flow (Buttolph *et al.*, 2006) (both developed by the United States Army Corps of Engineers; USACE), were selected and used as together they represent state of the art technology for the simulation of tidal currents, wave-induced currents, wave-induced longshore sediment transport, current induced sediment transport, and morphological change. The models were calibrated to observed hydrodynamics and to historical morphological evolution of the inlet, as documented from historical aerial photographs, dredging records and hydrographic surveys. Once calibrated, the modeling system was applied to study the hydrodynamic and sedimentary

processes at Shippagan Gully, ultimately providing a tool with which the above listed study objectives could be fulfilled.

A number of preparatory tasks were undertaken in order to be able to model the hydrodynamics and morphological processes present at Shippagan Gully. These preparatory steps included:

- Site visit and field data collection;
- Meetings with local stakeholders, as well as fishermen and their representatives;
- Study of historical documents including aerial photographs and dredging records;
- Analysis of local sediment properties;
- Analysis of past morphologic change, including erosion and deposition rates;
- Assessment of local water levels;
- Assessment of nearshore wave climate;
- Assessment of longshore sediment transport potential;
- Numerical modeling of combined wave and tidal hydrodynamics;
- Numerical modeling of sediment transport and morphology change;
- Calibration of the numerical model to observed hydrodynamics; and
- Calibration of the numerical model to observed morphology change.

Once these above listed steps were successfully completed, the calibrated numerical model was used as a tool to identify the governing processes responsible for the complex morphology present at Shippagan Gully. Based on the findings of this task, a number of alternative design scenarios were developed in collaboration with colleagues from PWGSC and the Department of Fisheries and Oceans (DFO) to address the navigable requirements of the inlet. The calibrated numerical model was then applied to determine the expected outcome of each of the alternative scenarios including the status quo option, with regards to morphology at Shippagan Gully. Results of these simulations are presented in Section 6.2.

Based on the findings of this numerical model study, guidance was provided to PWGSC concerning the future development of coastal infrastructure and sediment management

programs at Shippagan Gully. This thesis thereby serves as a scientific basis on which future engineering and economic decisions can be made concerning coastal works at this highly dynamic tidal inlet.

1.3 Organization of Dissertation

The present thesis has been organized into seven main sections. The first section provides a brief introduction to the topic of the study, including the main objectives, the proposed methodology and the scientific and practical significance of the work presented herein. The second section provides a literature review of coastal-related topics which together form the scientific basis required to perform this study. This section also includes examples of past research and case-studies regarding tidal inlets which together serve to illustrate many of the topics presented in the literature review in a practical context. The third section provides a site description, which summarizes the important geographic and historical features of Shippagan Gully and the surrounding area, both of the built and natural environments. The fourth section of the present thesis includes all non-numerical modeling based analyses completed throughout the course of the study. This includes the collection of pertinent data, the assessment of past and present morphological change, the study of local environmental conditions and the preparatory tasks required for the numerical modeling component of the study. The fifth section documents the numerical modeling of hydrodynamics and morphology at Shippagan Gully. This section includes descriptions of the numerical models and supporting software which were consulted throughout the course of the study. It also includes details of the set-up, simulation and calibration of the numerical modeling system. The sixth section presents the application of the numerical model to fulfill the objectives of the study, thereby the identification of governing morphological processes and the assessment of alternative design scenarios. This section has been referred to as *Analysis and Discussion* as opposed to *Results and Discussion*. This is due to the fact that the numerical model was used as an analytical tool with which coastal processes at Shippagan Gully could be studied as opposed to simply being a numerical system which provides results. The numerical analyses are discussed within the same section such that the individual tasks and modeling scenarios could be analysed and

discussed sequentially. The seventh and final main section of this thesis presents conclusions based on the findings of the study as they pertain to the objectives outlined in Section 1.2. Guidance which was provided to PWGSC concerning future development and maintenance at Shippagan Gully is summarized in this section and potential areas for future work are highlighted.

2.0 LITERATURE REVIEW

Coastal inlets are common features of the coastal zone which connect the ocean or sea to an inland body of water. Natural coastal inlets are found all over the globe and are highly dynamic and variable in their nature. Several types of coastal inlets exist, categorized primarily by the hydrodynamic processes responsible for their formation and maintenance. For example, tidal inlets are inlets which are formed and maintained by ebb and flood flows entering and exiting a tidal lagoon. Inlets can also be created by waves, river outflow, or are man-made to serve a navigational or environmental purpose. Another defining feature of coastal inlets is the nature of the connected bodies of water. For example, a coastal inlet could connect the ocean to a tidal lagoon, a river estuary, or a man-made harbour, each having very different hydrodynamic and morphological characteristics.

An important differentiation to make with regards to coastal inlets is whether or not the inlet is serviced by land based flow, or in other words is part of a river system. When no land based flow exists at a coastal inlet, the source of all flow in and out of the inlet is the ocean or the sea. This type of inlet is typically referred to as a tidal inlet, as it is the tides which are responsible for the exchange of water through the inlet. Tidal inlets are common in North America, where vast expanses of sandy shorelines are present in conjunction with significant tidal ranges. Shippagan Gully can be categorized as a tidal inlet and, as a result, the present study and literature review focuses on this topic.

2.1 Tidal Inlets

A typical tidal inlet system is composed of a tidal lagoon and a short flow pathway (the inlet) which connects the lagoon to the open ocean. A tidal lagoon is a protected body of water which fills and drains (partially or fully) with each tidal cycle, via the tidal inlet. Tidal lagoons and inlets are naturally formed in many different environments and for many different reasons. In locations such as North America's eastern seaboard, long expanses of shoreline are subjected to continual coastal sediment migration (referred to as longshore transport). In these environments, tidal lagoons are commonly formed behind morphological features called barrier

islands. In this scenario, the tidal inlet transects the barrier island(s) providing a pathway between the ocean and the sheltered lagoon. Shippagan Gully is an example of such an inlet. As such, it is useful to discuss tidal inlets in this context for the present study.

Barrier islands are long spits of sediment formed by years of wave induced longshore coastal sediment transport (a term given to the lateral migration of sediment along a coastline). In the absence of hydrodynamic forcing such as tides, this process could eventually result in the barrier island re-joining the mainland and thus creating a separate, protected body of water in its lee. In the presence of tides however, flow is driven into the back-bay system (behind the barrier island) with every flood tide and returned to the ocean with every ebb tide. Depending on the local tidal range and the volume of the back-bay system, this perpetual inflow and outflow of seawater is often powerful enough to sustain a self scouring inlet, thus halting the barrier island from reaching complete closure. An example of a typical tidal inlet - barrier island system is presented in Figure 2.1.



Figure 2.1: An example of a classical barrier island transected by a natural tidal inlet. Oregon Inlet, North Carolina (Google Earth, 2011).

The width and depth of a tidal inlet depend greatly on the local coastal processes as well as the volume of the back-bay system. In general, large (cross-sectional) inlets are found where either the tidal range is large or where the back-bay system is expansive in volume (or both). Both of these scenarios result in large volumes of water (referred to as a *tidal prism*) entering and exiting the tidal lagoon via the inlet with each tidal cycle. This process can result in a large flow rate through the tidal inlet which in turn scours a larger natural channel. Conversely, narrow and/or shallow tidal inlets are typically found where the tidal lagoon is smaller in volume or the tidal range is minimal.

The previous statements of course exclude the effects of waves, which are extremely important when considering tidal inlet morphology. Waves tend to mobilize coastal sediment and facilitate sediment migration in both the cross-shore and longshore directions. As such, waves are capable of transporting a great deal of sediment to and from the nearshore zone surrounding a tidal inlet and thus further complicating the local morphology. Furthermore, energetic storm waves are capable of producing coastal erosion, particularly when accompanied by a storm surge. In fact, it is often the enormous erosive power of waves which trigger the initial formation of a tidal inlet through a barrier island, a process which is known as a *break-through* (Dean and Dalrymple, 2002).

2.2 Tidal Inlet Hydrodynamics

There are two principal contributors to the hydrodynamics at most tidal inlets; astronomical tides and water waves. Both processes strongly influence flow magnitudes and flow patterns in and around the inlets and their combined effects can be extremely complex, particularly with respect to inlet morphology. Prior to examining these complexities however, it is essential to have a basic understanding of each hydrodynamic mechanism separately.

2.2.1 Astronomical tides

Astronomical tides are caused by a combination of forces acting on individual water particles. These forces include the gravitational attraction of Earth, the centrifugal force generated by the rotation of the Earth and the gravitational pull from other celestial bodies such as the Moon and

the Sun (Kamphuis, 2000). Of these factors, the gravitational pull of the Moon and the Sun are the most significant. These two primary celestial tidal forcings create what are often referred to as equilibrium tides, since they result from the assumption that they act on water particles for a very long time such that equilibrium is achieved between their resulting force and Earth's gravity (Kamphuis, 2000).

The gravitational pull from the Moon results in a small horizontal force on water particles. On the side of the planet which faces the Moon, water particles are pulled in its direction. At this, the location on Earth which is closest to the Moon, the gravitational pull is at its maximum. Conversely, on the opposite side of Earth, water particles are pushed away from the Moon. This occurs due to the fact that the gravitational pull induced by the moon is the weakest at this location. The end result is two bulges in the water surface profile on opposite sides of the planet and in line with the position of the Moon. If the whole Earth were covered with water, this phenomenon would result in an elliptical water surface profile, with the long axis of the ellipse in line with the location of the Moon (illustrated in Figure 2.2). The combined effect of Earth's rotation and the orbit of the Moon results in a relative frequency of 24.84 hours. Since the water surface profile is elliptical and thus possesses two bulges with respect to a spherical Earth, the resulting tidal period is in fact 12.42 hours.

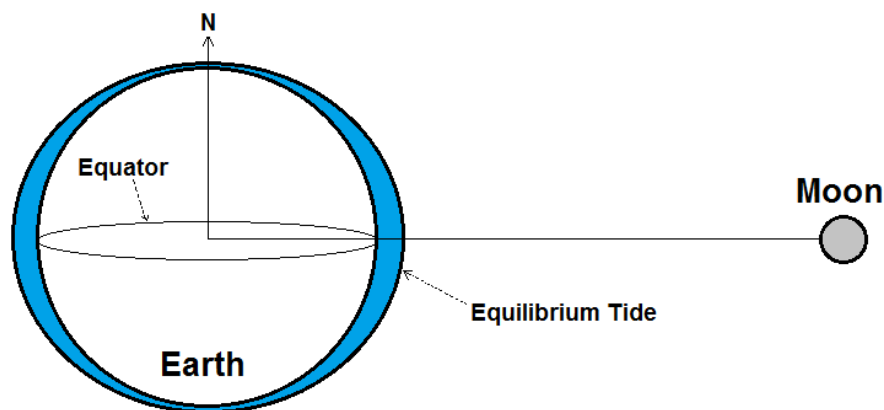


Figure 2.2: Lunar equilibrium tide, creating an elliptical water surface profile over a spherical, water-covered Earth with a period of 12.42 hours.

The Sun creates a similar gravitational pull on the Earth resulting in a second, smaller set of bulges on opposite sides of the conceptual water-covered Earth. Since our day is measured with respect to the Sun, the period of the tidal cycle generated by the Sun is 12 hours.

Since both of these tidal forcings act continually, their combined effects are important. When the Moon and the Sun are aligned (at new Moon and full Moon) the tides will be higher than usual as both gravitational pulls will operate together. This occurrence is referred to as a spring tide. Conversely, at quarter Moon the forces of the Sun and Moon will act at 90° to each other, thus limiting each other's gravitational pull and subsequently resulting in lower than normal tides. This occurrence is called a neap tide. Since the Moon orbits the Earth monthly (29.3 day period), two springs and two neaps are incurred per month.

These two recurring tidal forcings are referred to as tidal constituents. A tidal constituent is an independent tidal forcing which can be attributed a frequency (12.42 hours for the Moon and 12 hours for the Sun) and amplitude (Dronkers, 1964). If the planet were in fact covered in water of uniform depth, and the Earth's rotation and celestial orbits were strictly circular and two-dimensional (in plane), the amplitude of the equilibrium tidal constituents would be consistent along any latitude. This is not the case however. Staying with the assumption that the Earth is covered in water of uniform depth, the Sun and the Moon are seldom in plane with the Earth's equator, nor is their motion strictly two-dimensional or circular. In fact, the Moon and Sun generally have a north or south declination with respect to the equator as pictured in Figure 2.3. Thus, an observer traveling at constant latitude would experience two tides per day of unequal height. This phenomenon is referred to as daily inequality and is represented as its own constituent (one for solar and one for lunar). Daily inequality is most dramatic when the Sun or Moon is furthest north or south of the equator. Both lunar and solar inequality are harmonic processes with the lunar cycle repeating itself every lunar month (29.3 days) and the solar cycle every calendar year. In addition to daily inequality, there are dozens of other secondary effects which are each attributed a tidal constituent of their own (Dean and Dalrymple, 2002).

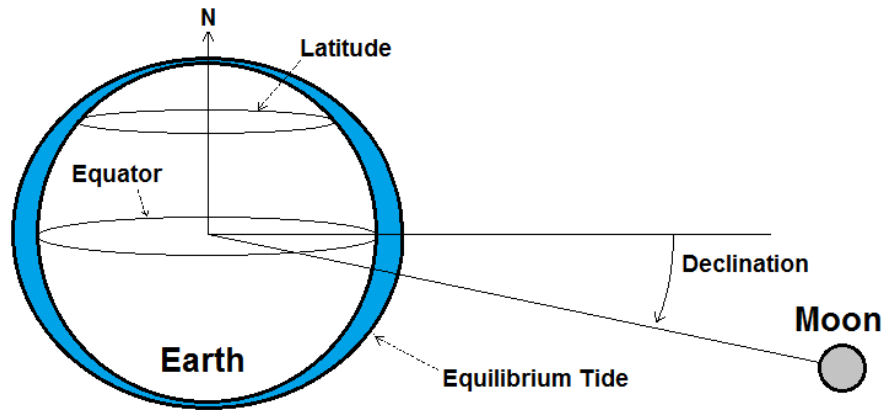


Figure 2.3: Lunar daily inequality, caused by north or south declination of lunar orbit with respect to the equator.

The previous paragraphs have of course been based entirely on a spherical Earth which is covered in water of uniform depth. This assumption is obviously incorrect as the various landmasses present on our planet interrupt the idealized tidal model. In fact, the only place on the planet in which true equilibrium tides can theoretically exist is in the southern hemisphere where an uninterrupted band of water circles the globe encompassing several different oceans. Even in this part of the world however, equilibrium tides take time to progress through the various water bodies. As such, actual water levels lag behind the theoretically determined constituents (Kamphuis, 2000).

Another large factor in the propagation of tides around the globe is the presence of seas, bays and estuaries. Certain tidal constituents tend to resonate in these partially closed water bodies, thus further complicating local tidal predictions. In certain locations some tidal constituents are amplified while others simply disappear. Often, the shape of a bay will have a funnelling effect on water levels resulting in spatial amplification. The Bay of Fundy is a prime example of this, where the local tidal constituents are amplified such that tides of up to 17 metres are observed, giving the Bay of Fundy the largest tidal range in the world (Sankaranarayanan and McCay, 2003).

Daily inequalities are often strongly influenced in resonating bays, with the difference between larger and smaller daily tides being increased by the local resonance to the point at which the smaller tide becomes non-existent. When this occurs, the typical semi-diurnal (twice per day) tide becomes what is known as a diurnal (once per day) tide (Kamphuis, 2000).

2.2.1.1 Tidal Analysis

The basic principal of performing a tidal analysis is to separate a measured water level record into as many of its constituents as possible. The tide can then be represented as a harmonic summation as shown in Equation 1 (Dronkers, 1964).

$$\eta_T(t) = \sum_{i=1}^I a_i \cos(\omega_i t + \alpha_i) \quad (1)$$

where $\eta_T(t)$ is the water level at time t , a_i and α_i are the respective amplitudes and phase angles of the tidal constituents, and ω_i is their angular frequencies. Since angular frequencies are already known, tidal analysis is therefore the task of determining values for a_i and α_i from measured data sets. Since there are many tidal constituents acting at a given location, separating and analysing the different tidal signals can be a very difficult task. As such, only the most significant constituents are typically considered.

Several numerical models have been developed to facilitate the analysis of tides. These models include (but are not limited to) MIKE 21 (DHI, 2011), ADCIRC (Luettich and Westerink, 2011) and an un-named tidal analysis package developed by Mike Forman at the University of British Columbia (DFO, 2009).

Once amplitudes and phase angles of the significant constituents have been deciphered for a given location, Equation 1 can be applied to predict future tidal levels at any given time. This process has been undertaken for most major ports and navigable waterways around the world, thus providing operators with reasonable estimates of navigable depths (Kamphuis, 2000). In fact, many previously analysed tidal constituents from around the globe have been documented and published and are publically available.

2.2.1.2 Tidal Currents

Naturally, substantial currents are induced by the perpetual raising and lowering of water levels around the globe. These currents can be considered extremely long water waves which are hundreds of kilometres in length. Since the depths in the nearshore zone (where water level fluctuations are important) are relatively shallow (order of metres), small amplitude wave theory (first discussed by Airy in 1845 and discussed further in Section 2.2.2 below) can be applied by use of the shallow water long-wave velocity equation (where $d/L < 0.05$). This equation is presented below as Equation 2.

$$C = \sqrt{gd} \quad (2)$$

where C is the velocity of propagation, g the gravitational acceleration and d the water depth. The wave length of an equilibrium tide is therefore:

$$L = CT \quad (3)$$

where T is the wave period, which in this case is 12.42 hours for the semi-diurnal lunar constituent and 12 hours for the semi-diurnal solar constituent. From Equations 2 and 3 it is seen that as a tidal current propagates into shallower water (decreasing depth), both the velocity of propagation and the wave length will decrease.

Given that tidal currents possess wave lengths of hundreds of kilometres, their propagation is dramatically influenced by the rotation of the Earth. As such tidal currents do not propagate in straight lines through the Earth's numerous connected water bodies. Instead, tides tend to propagate in circular patterns, forming extremely large eddies. One such eddy exists in the Gulf of St. Lawrence, where tides propagate in a counter-clockwise manner with the center of rotation (referred to as an amphidromic point) some 100 kilometres northeast of Prince Edward Island, near the Iles de la Madeleine (Forrester, 1983). At the center of rotation, water levels remain approximately constant. Figure 2.4 shows this rotation of tidal flow for the Gulf of St. Lawrence, where dashed lines represent lines of equal tidal range and solid lines show lines in which the tide has the same phase (arrives at the same time).

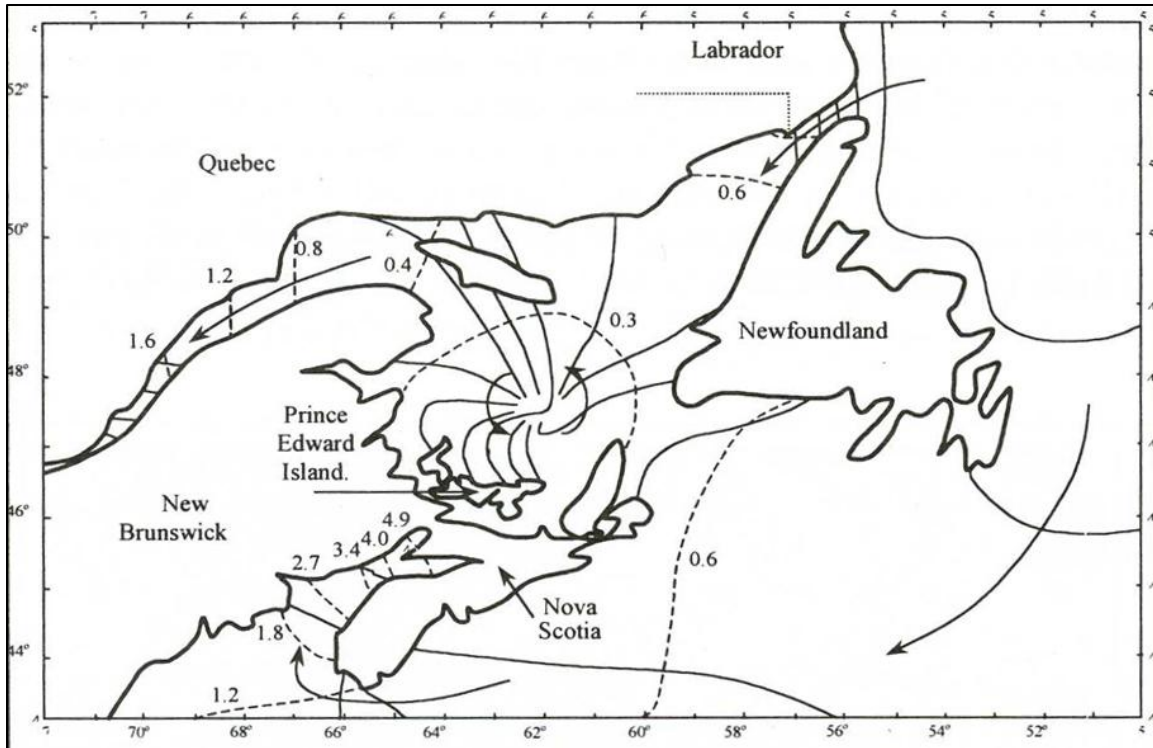


Figure 2.4: Semi-diurnal tide in the Gulf of St. Lawrence, illustrating lines of similar amplitude, lines of similar phase and the center of rotation, referred to as an amphidromic point (Kamphuis, 2000).

The presence of a counter-clockwise tidal circulation in the Gulf of St. Lawrence is important in the context of the present study, as it directly affects the propagation of currents near Shippagan Gully. Since the New Brunswick coastline runs approximately north-south and is located to the west of the center of tidal rotation in the Gulf of St. Lawrence, a weak tidal current is incurred from the north to the south along the Acadian Peninsula coast. Thus, in the absence of waves, a small longshore current is present at Shippagan Gully as water levels arrive at the inlet from the northeast and are propagated down shore.

In addition to the counter-clockwise tidal-circulation in the Gulf of St. Lawrence is the presence of a significant outflow jet from the St. Lawrence River. This outflow forms an estuarine plume which generates a weak southerly surface current due to the geometric properties of the estuary mouth and the relative buoyancy of the fresh water outflow (Sheng, 2001). This current follows closely the Gaspé Peninsula shoreline and subsequently passes the eastern shore of the Acadian

Peninsula in a southerly direction. This process strengthens the existing tide-induced current which is present in this region (Sheng, 2001).

2.2.1.3 Tidal Flow through an Inlet

As tides arrive at a shoreline, the water level along that shoreline will rise and fall in phase with the water level offshore (along lines of similar phase). However, in the presence of a partially closed bay such as those which are found behind barrier islands, the rising water levels in the sea will cause a current to flow into the bay, thus raising its water level. This inflow of water is referred to as the flood, while the subsequent outflow that comes with the falling water levels is called the ebb. If the bay is small enough and the entrance to the bay (tidal inlet) is large enough, the rise and fall of the water level in the bay will be precisely in phase with that of the ocean. As such, when both water bodies reach high tide, flow through the inlet will be zero (called the high water slack tide). Similarly, zero flow will be experienced at low tide (low water slack tide). Following this pattern, the maximum flows through a tidal inlet will theoretically occur at the halfway point between high and low tides, where the mean water level is reached.

When a tidal inlet is narrow, or where the tidal lagoon is large, the filling of the lagoon will be slowed during the flood tide. As such, the maximum water level in the tidal lagoon will occur later than in the ocean or sea, and thus flow will continue to progress into the bay for some time after high tide has been reached in the sea. Similarly, the maximum currents flowing through the inlet will no longer occur at the mean water level (half way between high and low tides). Instead, the currents will lag behind depending on the geometric characteristics of the tidal inlet and lagoon (Dean and Dalrymple, 2002).

A useful tool in analysing flow through tidal inlets is called the *tidal prism*. The tidal prism is the total volume of water that must flow in and out of a tidal lagoon with every tidal cycle (D'Alpaos *et al.*, 2010). It can be calculated by determining the differences between the high and low water levels throughout the tidal lagoon and multiplying them by the surface area. By calculating the tidal prism for a tidal inlet, average flow rates and current velocities through the inlet can be calculated based on the geometric properties of the inlet itself. Furthermore, tidal

prisms provide a means by which inlets can be classified and compared in a quantitative manner (D'Alpaos *et al.*, 2010).

For inlets with large tidal prisms, or for inlets in which the channel cross-section is relatively small, extremely large currents can be experienced both entering (flood) and exiting (ebb) the inlet. The erosive forces which accompany these high velocity flows can create a great deal of sediment transport in and around a tidal inlet. As such, tidal inlets are typically highly dynamic with respect to their long-term morphology. Sediment transport through tidal inlets will be discussed further in Section 2.3.

2.2.2 Nearshore Waves

2.2.2.1 Background on Wind Waves

In order to comprehend the complex nearshore hydrodynamics which occur in proximity to tidal inlets, some basic background on water waves is required. Furthermore, a basic understanding of the principal transformations which occur as waves propagate towards the shoreline must be discussed and understood.

Surface water waves are fluctuations in water levels which can be grouped into several different categories. In essence, any mechanism that creates a disturbance in a water body will ultimately create water waves. Water waves are commonly classified by their frequency (the inverse of wave period), with long period waves such as tides and tsunamis at one end of the frequency spectrum, and capillary waves (ripples) at the other. In the middle of the wave frequency spectrum lies a classification referred to as *gravity waves*, which are waves with periods typically between 1 and 30 seconds and wave heights which rarely exceed 10 m (Kamphuis, 2000). This wave classification receives its name from the restoring force which works to restore the still water level. Gravity waves are typically created by winds which blow over large expanses of open water, inducing shear stress at the air-water interface. This shear stress displaces the water in the direction in which the wind is blowing subsequently creating *wind waves*. The height of wind waves is directly related to three important variables: the wind speed, the

distance over which the wind blows (referred to as the fetch) and the duration for which the winds are sustained (Kamphuis, 2000).

Wind waves can be classified into two categories referred to as sea and swell. Sea waves are locally generated wind waves and are highly variable in both height and period. These waves tend to propagate more or less in the prominent wind direction and tend to create extremely rough seas where strong winds are present. Storm systems are typically accompanied by high local winds and subsequently large sea waves (Goda, 2000).

On large bodies of water, locally generated wind waves will continue to travel beyond the area in which they are generated. As they propagate across open water, their profile becomes more regular as individual waves separate and waves of similar frequency group together. The result of this process is well defined groupings of regular waves commonly referred to as *wave sets*, which continue to propagate over areas in which there is no-longer a generating wind. These waves are referred to as *swell waves* and are experienced on most ocean-bordering coastlines around the world (Goda, 2000).

Generally speaking, sea and swell waves occur simultaneously at a given location, as waves which have arrived from elsewhere are subjected to local winds. In these situations, the sea and swell wave components can be separated by period, with swell waves having a much lower frequency than sea waves (Goda, 2000).

There are numerous numerical models available to model the generation and propagation of water waves. These models can be loosely divided into two groups; phase-averaged models and phase-resolving models. Phase resolving models calculate the entire wave height profile (crest to trough) as the waves propagate across the model domain, thereby providing a realistic simulation of individual waves and wave trains in motion. Phase-averaged models on the other hand neglect changes in wave phase and instead calculate the wave action density for given wave spectra, from which characteristic wave parameters are inferred (Aquaveo, 2011). The most common wave parameter to be calculated in a phase-averaged wave model is referred to as the significant wave height, which is by definition, the average of the highest one third of the

wave heights. Significant wave height is an important parameter which is used in the majority of conventional coastal engineering practises due to the fact that it has been shown to closely approximate the highest wave height observed by a trained observer (USACE, 1984). As such, many classical formulations have been based on significant wave height, making it the basis of many coastal engineering calculations which are still used today.

Largely due to the significant reduction in computational time required by phase-averaged models over phase resolving models, phase-averaged models are the more practical model for simulating large domains or large temporal scales. SWAN (Ris, 1997), STWave (Smith *et al.*, 2001), MIKE21 (DHI, 2011) and CMS-Wave (Lin *et al.*, 2008) are all examples of widely used phase-averaged wave models.

In smaller domain applications, particularly where nearshore wave transformations (discussed in the following section) and wave reflection are of importance, phase-resolving wave models are likely to be the better suited model. Although computationally more demanding, these models can produce much better solutions for transient wave data. An example of a typical application in which a phase-resolving model will likely provide better results is studying the wave agitation inside a sheltered harbour or marina. Examples of widely used phase-resolving wave models are CG-Wave (Demirbilek and Panchang, 1998) and BOUSS-2D (Nwogu and Demirbilek, 2001).

2.2.2.2 Nearshore Wave Transformations

As both sea and swell waves propagate towards a shoreline and advance into shallower waters, they begin to feel the presence of the sea-floor. This transition occurs when the depth is approximately half of the wave length (Airy, 1845). At this point, several changes to the water surface profile begin to occur and are hence referred to as *nearshore wave transformations*. The four principal nearshore wave transformations are shoaling, refraction, diffraction, and breaking and are discussed in the following sections as they pertain to tidal inlets.

Once waves begin

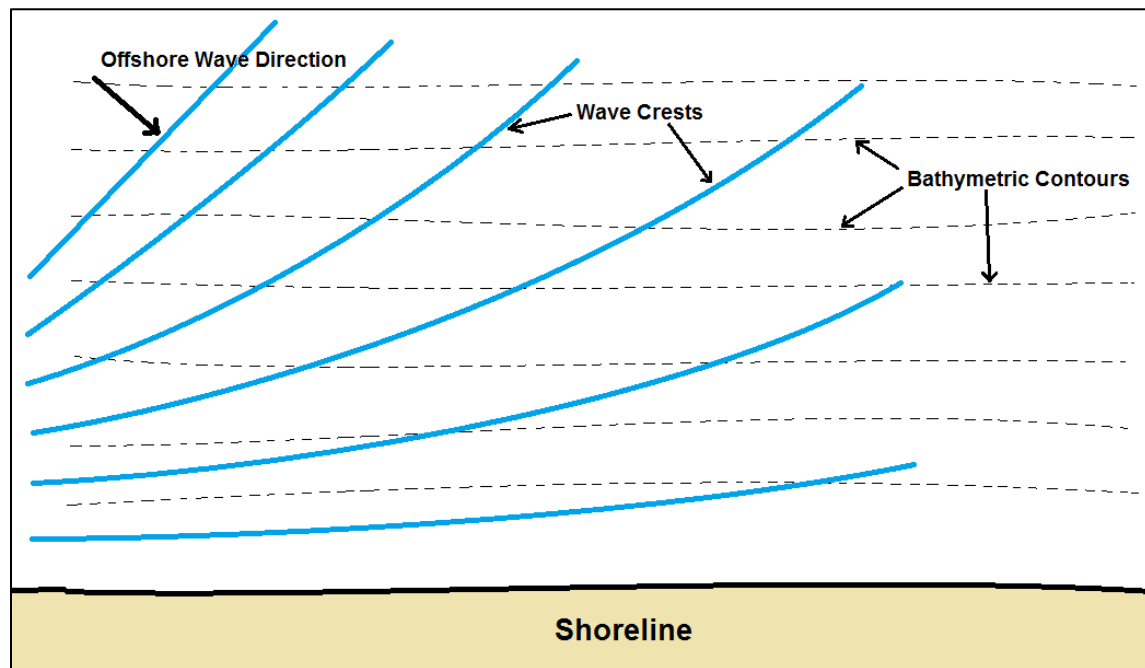
The result is a progressive steepening of the wave profile. This process continues as the wave propagates into shallower water until the wave profile becomes steep enough that it loses stability and breaks (discussed further below). Due to the process of wave shoaling, nearshore waves are typically much higher than they once were offshore. As such, waves arriving at a tidal inlet will have an increased height relative to the nearshore depths surrounding the inlet (assuming they remain un-broken). Similarly, their propagation through the tidal inlet (assuming the wave direction is perpendicular to the inlet mouth) will be determined largely by the available depths.

Analytically, wave shoaling is typically represented by a local shoaling coefficient (K_s), which is simply the ratio of the wave height at any given nearshore location to its offshore wave height. There are several theoretical and empirical methods available to predict K_s , which generally relate the offshore wave height (H_o) and period (T) to the local water depth (d). The most common of these methods is derived from small amplitude wave theory (first discussed by Airy in 1845) and is as follows:

$$K_s = \frac{H}{H_o} = \sqrt{\frac{1}{2n} \frac{1}{\tanh kd}} \quad (4)$$

Where H is the nearshore wave height at the location of interest, H_o is the offshore wave height, k is the wave number (2π divided by the wave length, L), d is the local water depth and:

shoaling will occur simultaneously along the entire length of the wave in the longitudinal direction (shore parallel). However, this is rarely the case as shorelines and bathymetric contours do not typically form straight lines and waves rarely approach the shoreline at exactly 90° . The result is that one end of the wave crest will reach shallower water first and begin to shoal prior to the opposite end of the wave crest. Since shoaling slows the wave speed and steepens the wave profile, the portion of the wave in shallower water will slow prior to the portion of the wave which is in deeper water, creating a bend in the wave crest (in plan-view). This process is referred to as refraction and ultimately results in wave crests bending as they propagate towards the shoreline in order to align themselves with the bathymetric contours. This process is illustrated in Figure 2.5 below.



$$K_r = \sqrt{\frac{b_0}{b}} = \sqrt{\frac{\cos \alpha_0}{\cos \alpha}} \quad (6)$$

where α is the wave crest angle relative to the shoreline and b is the distance between adjacent crests. In order to predict what K_r will be given the characteristics of an offshore wave, an estimation must therefore be made for the value of α at the location of interest. This is typically done using Snell's law, which is shown in Equation 7 below (Kamphuis, 2000).

$$\frac{\sin \alpha}{\sin \alpha_0} = \tanh \frac{2\pi d}{L} \quad (7)$$

where d is the local water depth and L the wave length at the given water depth. As such, d must first be known and L estimated from linear wave theory (or a higher order theory).

Once K_r and K_s have been approximated using the above formulations (or other), the offshore wave height can be multiplied by both coefficients to determine the nearshore wave height at the location of interest. This process is shown below as Equation 8:

$$H = K_s K_r H_$$

inlet or navigation channel. The concept of wave diffraction is shown in Figure 2.6 for both an idealized offshore breakwater and a typical engineered tidal inlet.

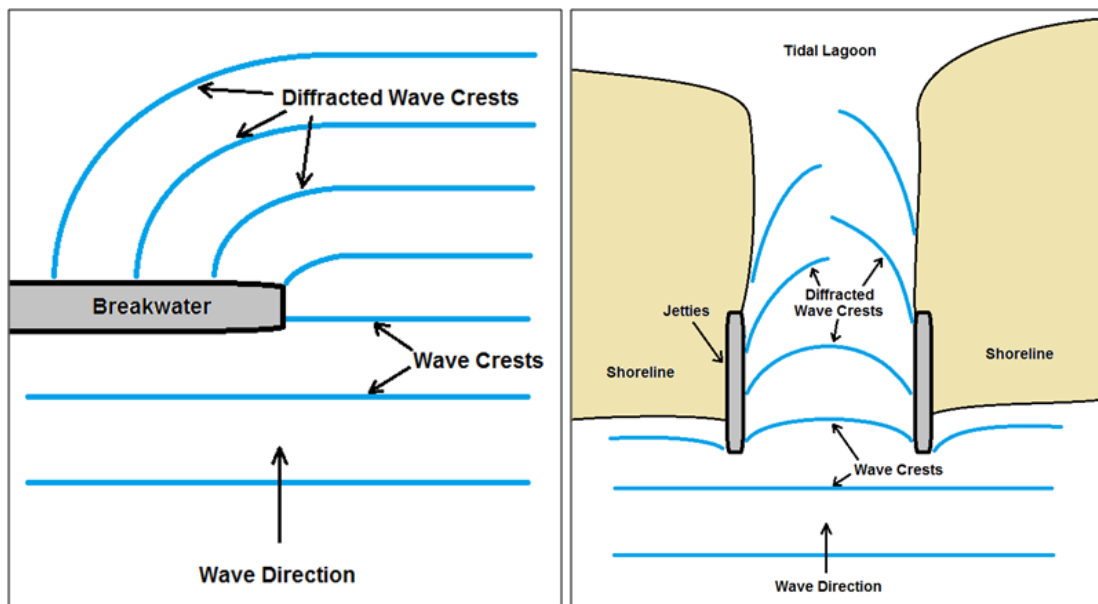
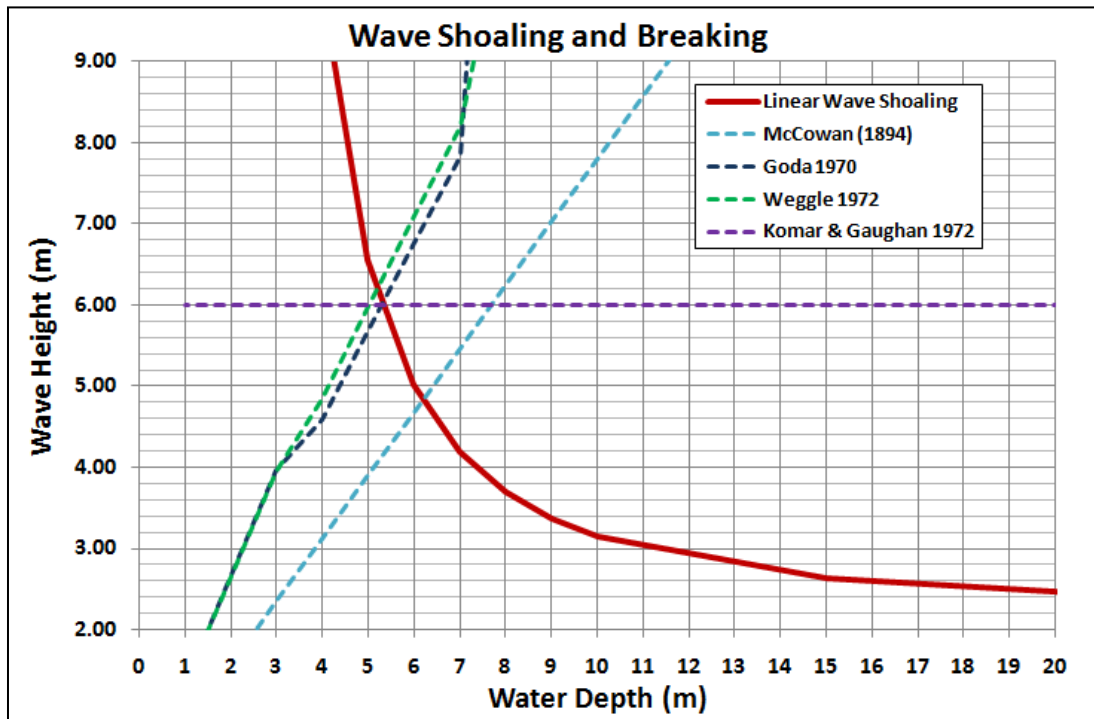


Figure 2.6: Wave diffraction around an offshore breakwater (left) and through a tidal inlet with engineered jetties (right).

There are dozens of accepted formulas to predict breaking wave heights and breaking wave depths given the offshore characteristics for a given wave. This has been the subject of numerous studies due to the fact that breaking waves create large forces on coastal structures (Rattanapitikon and Shibayama, 2000). Furthermore, the hydrodynamic effects of a breaking wave are a primary driving force behind coastal sediment transport. The simplest and oldest widely used breaking wave formula was first presented by McCowan in 1894 and relates breaking wave height to breaking wave depth as follows:

$$H_b = 0.78d_b \quad (9)$$

where H_b is the breaking wave height and d_b is the breaking wave depth. When plotted in conjunction with a shoaling formula such as the one which was presented in Equation 4 above, the line created by McCowan's formula will intersect the shoaling curve at the theoretical location of breaking. An example of this is shown in Figure 2.7, which also includes the breaking formulation of Goda (1970), Weggles (1972) and Komar & Gaughan (1972) for comparison.



The breaking wave criterion presented in Equation 9 is a relatively old and highly approximate method. Dozens of more recent researchers have shown breaking wave heights and depths to be influenced by other variables such as wave length and beach slope (Goda (1970), Weggel (1972), Komar and Gaughan (1972), Battjes (1974) and Smith and Krause (1990)). However, the method presented by McCowan still serves as a reasonable approximation in many cases (Rattanapitikon and Shibayama, 2000). More widely accepted and well validated formulas include the three additional formulations presented in Figure 2.7.

2.2.2.3 Wave Induced Currents

As waves propagate towards the shoreline, there is a mean transport of water in the direction of propagation. This phenomenon is referred to as *mass transport* and is not accounted for by linear wave theory (first discussed by Airy in 1845). Linear wave theory assumes that water particles subjected to surface waves will travel in *closed* elliptical orbits. This has been shown however to be untrue, as in reality the elliptical orbits move in the direction of wave propagation. As such, the elliptical orbits overlap and therefore do not close. This phenomenon is illustrated in Figure 2.8. The process is explained by the fact that the portion of the ellipse in which water particles move forwards occurs under the wave crest. Conversely, the backwards motion of water particles corresponds to the wave trough. Since the total depth at the wave crest is larger than at the trough, this portion of the elliptical orbit is also slightly larger, and thus results in the net forward movement of water particles. It follows that mass transport is larger for more energetic waves, or in other words waves with larger wave heights (

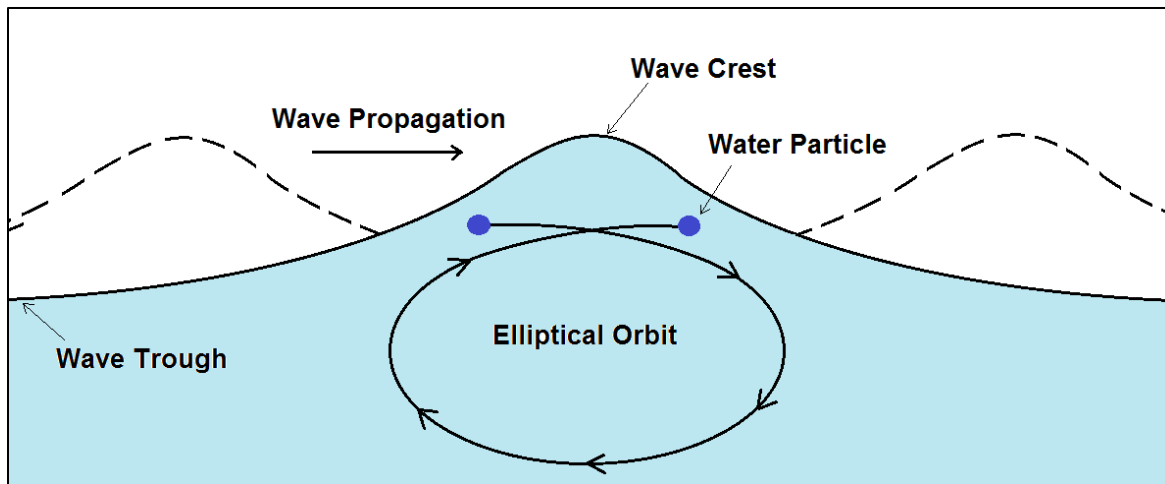
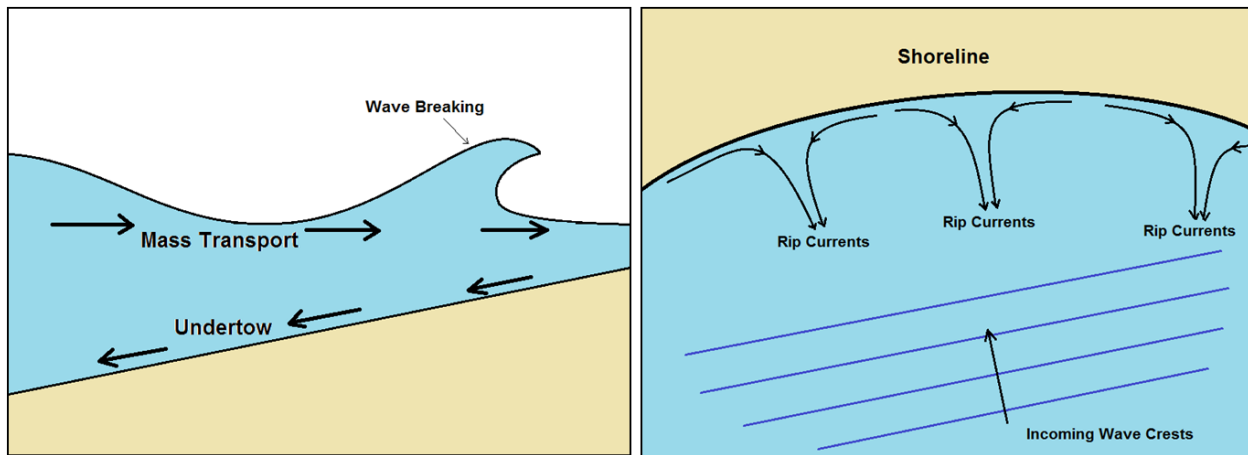


Figure 2.8: Overlapping elliptical orbit of water particles under surface waves.

As with any dynamic mechanism, mass transport has a momentum component. Following Newton's second law, forces must be generated whenever momentum changes in either magnitude or direction. In recognition of this, Longuet-Higgins and Stewart (1963) introduced the concept of wave momentum flux, designating the sum of the momentum flux and the mean pressure as the *radiation stress*. Radiation stress is a vector quantity which drives water currents in the nearshore zone in the general direction of wave propagation.

At the moment in which a wave breaks, the momentum flux significantly decreases as wave energy is dissipated. This change in momentum flux must be balanced, which is achieved by inducing large forces in the surf zone (nearshore zone in which waves break) in the direction of

direction, it can be considered a cross-shore process. The shoreward transport of water induced by mass transport must however be returned to the sea. On a long, uniform beach this process is referred to as *undertow* (refer to Figure 2.9). Undertow is a cross-shore return current (in the seaward direction) which propagates along the sea floor, thus counteracting mass transport and rendering the net cross-shore flow to zero (as must logically be the case). When the shoreline is non-uniform, *rip currents* can be formed in which converging currents exit together through the entire water column (as opposed to along the bottom as is the case for undertow). On most coastlines around the world, undertows and rip currents are found in combination. These processes return the shoreward transport of water to the sea and thus balance the theoretical cross-shore hydrodynamic model (refer to Figure 2.9).



a net longshore current in one direction along the coastline. Not surprisingly, this current is typically accompanied by sediment transport (discussed in Section 2.3.1). As such, its direction and magnitude can often be inferred from local morphological characteristics along the shoreline. This topic is discussed in the following sections.

2.3 Tidal Inlet Morphology

All of the coastal processes which were described in the previous sections have the ability to move sediment. This is due to the fact that they create currents which induce shear stresses capable of picking up sediment and moving it in the direction of flow through the tidal inlet or within the nearshore zone. Given that many of the above-described processes occur simultaneously at or near tidal inlets, the interaction between the various sediment transport mechanisms is extremely complex and is a partial subject of the present study. However, basic principals can be derived from previous research and case studies which are summarized in the following sections.

2.3.1 Sediment Transport Mechanisms

Coastal inlets are generally extremely dynamic in nature and are often inherently unstable. This is due to the complex cross shore and longshore sediment transport mechanisms which exist in these regions. Not surprisingly, cross shore and longshore sediment transport are directly related to their hydrodynamic counterparts which were discussed in the previous sections. It is useful to describe these processes independently, while recognizing that they in fact occur simultaneously at a coastal inlet.

The most important hydrodynamic processes at a tidal inlet are of course tides and waves. As was previously discussed, both tides and waves induce strong currents in the nearshore zone and into a tidal inlet/lagoon system. These processes generate shear stresses along the bed which are capable of inducing sediment motion. If the shear stress exerted by the passing current exceeds the critical shear stress of the bed material (sediment), the bed material will begin to move in the direction of the current. The classical numerical representation of the requirement for incipient sediment motion is known as the Shields parameter and was originally presented by A. Shields in 1936. The Shields parameter is a non-dimensionalization of a shear stress which represents the ratio of fluid force on a given particle to the weight of that particle (Shields, 1936). The Shields parameter is presented below as Equation 10.

$$\tau_* = \frac{\tau}{(\rho_s - \rho)gD} \quad (10)$$

where τ is a dimensional shear stress, ρ_s is the density of the sediment, ρ is the density of the fluid (water), g is gravitational acceleration and D is the characteristic particle diameter of the sediment. In order for sediment motion to begin, the following relationship must therefore be met:

2.3.1.1 Bed Load

Once incipient sediment motion has been reached, the mobile sediment will be transported in one of two possible ways. These are referred to as *bed load* and *suspended load*. *Bed load* is the term given to sediment which is moving along the bed, either by rolling, sliding and/or hopping. Numerous researchers have studied *bed load* resulting in the existence of multiple formulas for the calculation of *bed load rate*. The classical methods involve a term referred to as Einstein's Phi which is shown in Equation 14 (Einstein, 1942).

$$\phi = \frac{q_{sb} \rho^{1/2}}{\gamma_s^{1/2} D^{3/2}} \quad (14)$$

Where ϕ is Einstein's Phi, and q_{sb} is the bed load rate (in m^2/s). The formula shown in Equation 14 can be re-arranged to solve for q_{sb} , however Einstein's Phi must first be known.

Where Y and Y_{CR} are determined from the Y-axis component of Shields Curve (equation listed above), V is the fluid velocity, b is a coefficient depending on the Reynolds number ($b \approx 4.25$ for rough turbulent flow) and W is the density of the bed material divided by the density of water.

2.3.1.2 Suspended Load

The second form of sediment transport, referred to as *suspended load*, is the transport of sediment (or bed material) which is entrained in the water column (i.e. not moving along the bed). Calculating the *suspended load* rate is a more difficult endeavour than calculating *bed load* due to the fact that it requires integration through the entire water column. As such, *suspended load* formulations are generally computationally complex. Similar to *bed load*, multiple researchers have proposed empirical or semi-empirical formulas to calculate the *suspended load* throughout the water column. A classical example of one such formulation is the Rouse-Einstein (1937, 1954) formulation which is presented as Equations 20 and 21 below:

$$\frac{\partial(Cd)}{\partial t} + \frac{\partial(Cq_x)}{\partial x} + \frac{\partial(Cq_y)}{\partial y} = \frac{\partial}{\partial x} \left(K_x d \frac{\partial C}{\partial x} \right) + \frac{\partial}{\partial y} \left(K_y d \frac{\partial C}{\partial y} \right) + P - D \quad (22)$$

where:

- C = depth-averaged sediment concentration
- d = total depth ($h + \eta$)
- h = still water depth
- η = deviation of the water-surface elevation from the still-water level
- t = time
- q_x = flow per unit width parallel to the x-axis (ud)
- q_y = flow per unit width parallel to the y-axis (vd)
- u = depth-averaged current velocity parallel to the x-axis
- v = depth-averaged current velocity parallel to the y-axis
-

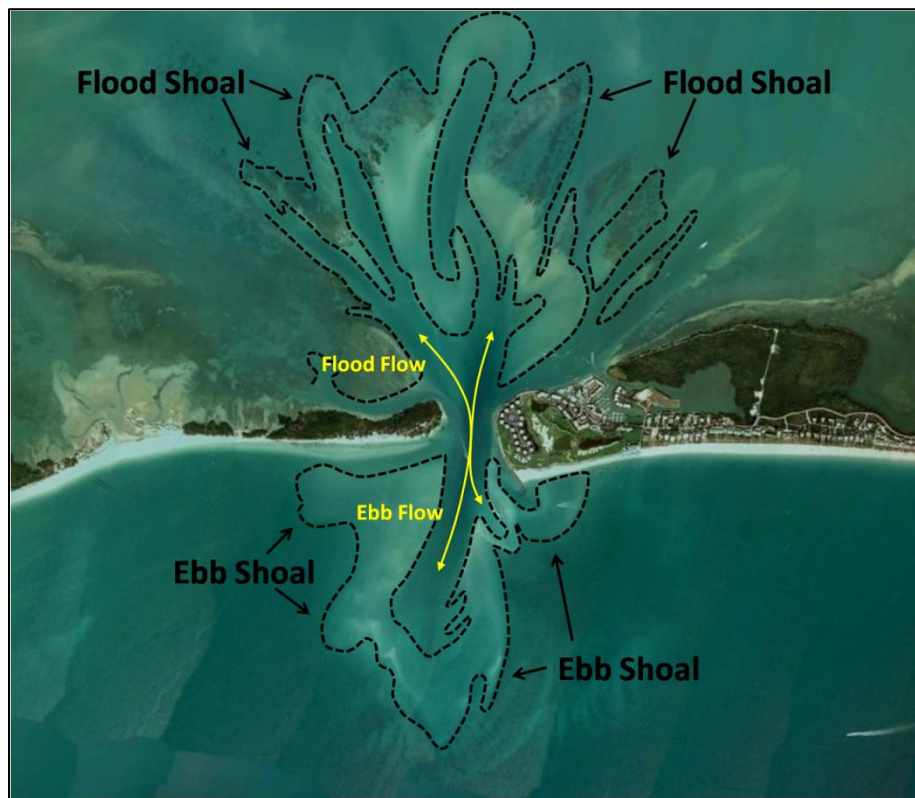
- ρ_w = density of water
- g = acceleration of gravity
- A = empirical coefficient typically ranging from 0.1 to 2

It can be seen from the above sections that numerous sediment transport models exist for the calculation of both bed load and suspended load, which can often be employed as stand-alone formulations, or in combination with other formulae. It can be seen that key parameters in all formulations are derived from the original work completed by Shields in 1936. As such, the concept of sediment transport can be more or less simplified to be a function of shear stress. To understand how sediment transport occurs at a tidal inlet, we must therefore look at the processes which produce shear stresses.

2.3.1.4 Cross-Shore Sediment Transport at Tidal Inlets

The most prominent cross-shore sediment transport mechanism at most tidal inlets is the hydrodynamics accompanied by tides. As was discussed in Section 2.2.1, a rising tide causes

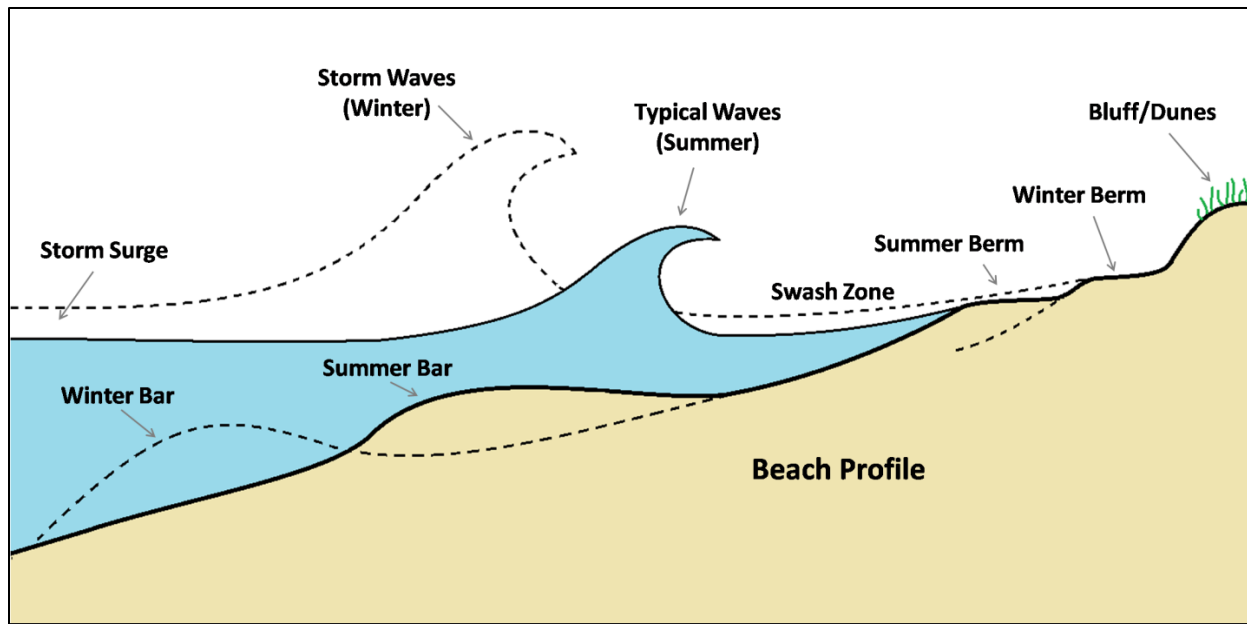
smaller. As such, the governing criteria for sediment motion may no longer be met and sediment will begin to be deposited. In the case of the ebb flow, sediment deposition will occur directly offshore where the ebb flow enters the ocean, thus creating what is commonly referred to as an ebb-shoal. Conversely the flood tide will transport sediment into the tidal lagoon, subsequently depositing the sediment in the confines of the sheltered back-bay system in the form of a flood-shoal. An example of this cross-shore morphology is illustrated in Figure 2.10 below.



producing shear stresses which exceed the critical shear stress for incipient sediment motion and thus induce sediment transport in the direction of wave propagation.

In the absence of any obstruction, waves propagating normal to the shoreline will break uniformly, eroding the bed at the breaking point (where the largest shear stresses induce the maximum sediment motion). The entrained sediment will then be deposited in one of two locations where the wave induced currents diminish; either in the swash zone (area immediately inshore of the breaking point), or immediately offshore and adjacent to the breaking point where it is transported by the undertow. The first of these scenarios results in the creation of a morphological beach feature referred to as a berm, while the second results in the creation of a feature which is called a bar (refer to Figure 2.11).

In the winter, waves are typically larger than they would be in the summer due to the increase in relative storm frequency and intensity. As such, more frequent and larger waves will break further offshore, thus shifting the primary breaking and subsequent erosion zone further from the shoreline. As is shown in Figure 2.11 below, this seasonal variance can create a multi-tiered beach profile with two berms and two bars, corresponding to the winter (high energy) and summer (low energy) seasons.



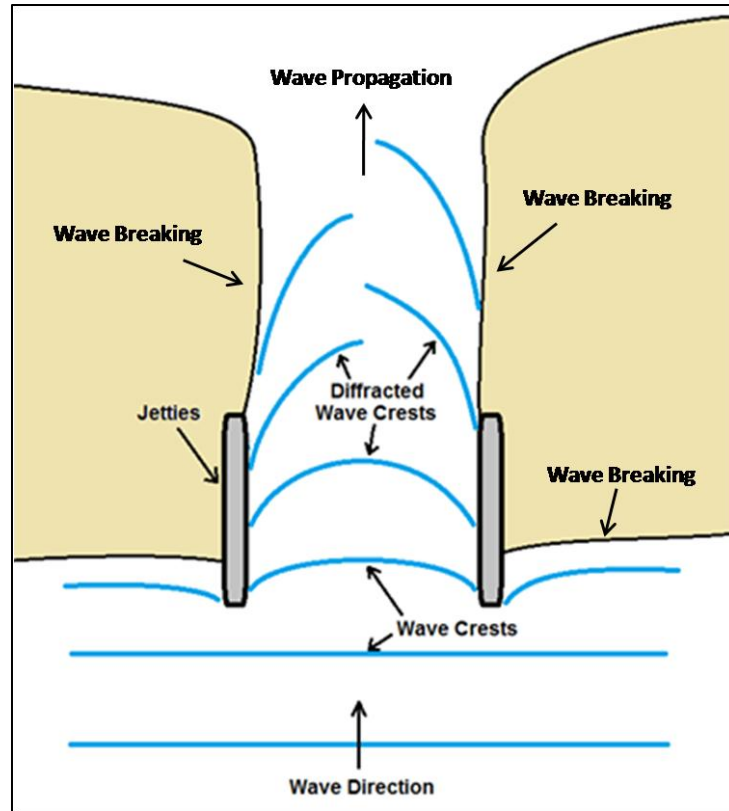


Figure 2.12: Wave propagation through an idealized, engineered inlet, illustrating wave breaking on inlet banks.

As is seen in Figure 2.12, the combination of refraction/diffraction inside the tidal inlet and the breaking process can induce large erosive forces along the inlet banks. This process can subsequently result in the widening of

induced currents (radiation stresses) which in this case have a longshore vector component. As such, sediment will be moved down the coastline.

There are two well known longshore transport formulae available to estimate potential longshore sediment transport rates. The first and best known equation for estimating bulk longshore sediment transport rate, Q_s , is published in CERC (1984), and is written as follows:

$$Q_s = \frac{I_s}{(\rho_s - \rho)(1 - n)g} \quad (24)$$

where:

$$I_s = 0.39P_{asb} \quad (25)$$

the case in coastal regions and thus a distinction must be made between the *potential* transport rate and the *actual* transport rate.

As was discussed previously, longshore currents are typically more frequent in one direction along a coastline, thus resulting in a net longshore current. The same can therefore be said of longshore sediment transport, where a net longshore transport direction is typically observed along a coastline. This direction can often be readily identified where abrupt changes in coastline are present, such as near coastal structures or natural headlands, as sediment will begin to build up on one side of the obstruction. An extremely important concept to understand with regards to longshore sediment transport is that when sediment is deposited on one side of an obstruction, the sediment supply is depleted in its lee (down-drift of the structure). As such, the presence of obstructions in the longshore sediment transport pathway will often lead to severe erosion further down the coastline. These processes are illustrated in Figure 2.13.

coastline; however the obstruction is dynamic in that it is itself a morphologic feature capable of movement and migration. Furthermore, the presence of the tidal inlet introduces cross-shore currents (ebbs and floods) which present a hydrodynamic obstruction to longshore sediment transport capable of altering its pathway.

In the most simplistic scenario, sediment will amass on the up drift side of a tidal inlet, as its passage across the inlet mouth is obstructed by deeper water and cross-shore currents. This in turn will result in erosion on the down drift side of the inlet, resulting in an offset inlet, where the up drift and down drift shorelines do not precisely align. However, at a coastal inlet with significant tidal flows, sediment arriving at the inlet will typically interact with the ebb and flood currents. The transported coastal sediment will enter the inlet from the up drift side where it is subsequently transported into the tidal lagoon with the flood current or offshore to the ebb-shoal with the ebb current. This represents an interaction between longshore and cross-shore processes, an immensely complex topic which is discussed in the following section.

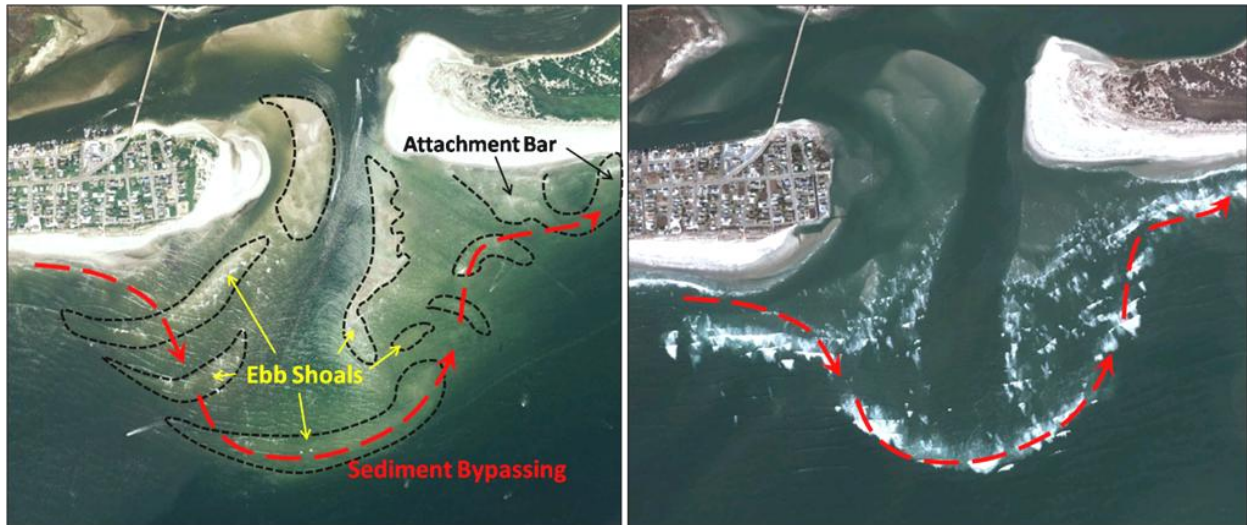


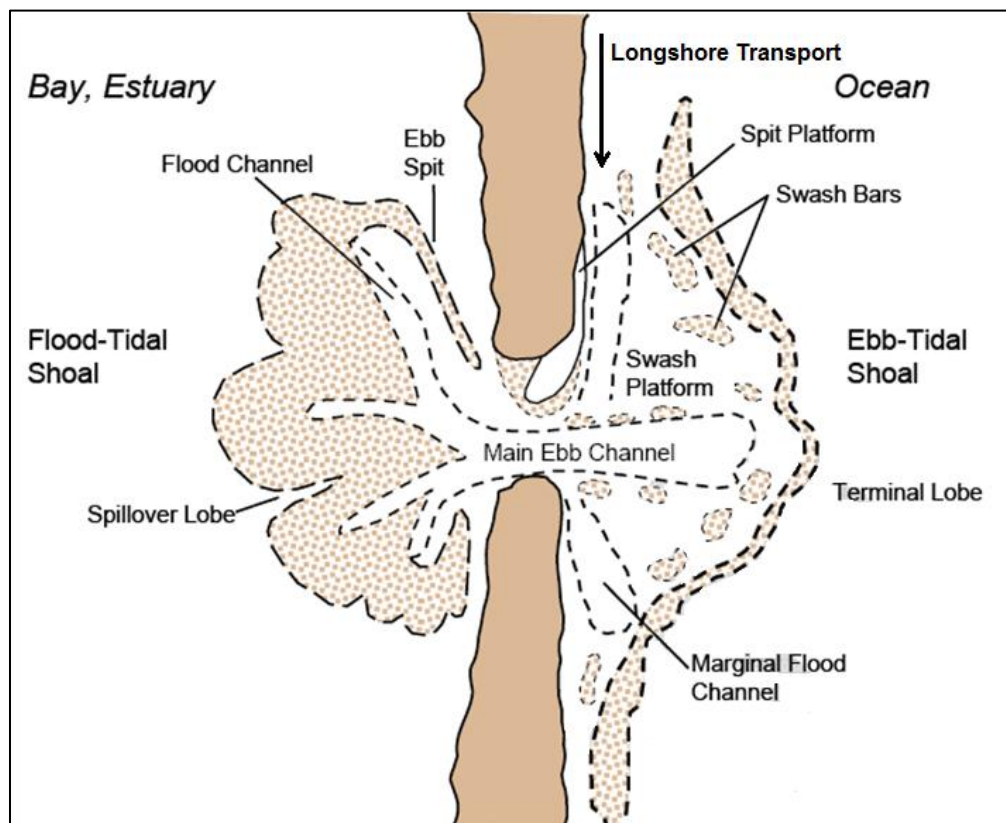
Figure 2.14: Corson Inlet, New Jersey, USA, showing extensive ebb-shoal complex (left) and waves breaking along the ebb-shoal (right). The sediment bypassing pathway is shown in red (modified from Google Earth, 2011).

project. As such, advances are being made through research (including the present study) in which engineered solutions to coastal inlet projects are tailored such that interruptions to naturally occurring sediment pathways are minimized.

As is pointed out in Seabergh and Kraus (2003), there are a number of techniques available to ensure that adequate sediment bypassing is obtained, through artificial means. These techniques generally include some form of controlled sediment trap, where accretion rates are predictable and can be artificially transferred to the down drift shoreline via dredging or pumping. A weir-jetty system is one such method in which sediment is contained ensuring an adequate beach on the up drift side of the inlet, while spill-over sediment is mechanically bypassed to the down drift shoreline (Seabergh and Kraus, 2003). This method controls and facilitates the artificial bypassing of sediment while maintaining a stabilized inlet mouth using coastal structures. Precise and well educated alignment of coastal jetties can also be effective in stabilizing

- The sediment supply,
- Local wave climate,
- The relative strength of wave action to tidal flow,
- The net and gross longshore sediment transport rates,
- Engineered activities in the channel, estuary and adjacent beaches (including dredging),
- The presence and configuration of jetties, and
- Other factors.

Figure 2.15 presents a definition sketch of typical tidal inlet morphology for a natural inlet in which there has been no human intervention.



of the swash platform, which are formed as incoming waves break over the ebb-shoal and carry sediment into the relative shelter in its lee.

The inlet depicted in Figure 2.15 is an example of an idealized inlet where morphological processes are fully developed, approximately symmetrical and well balanced. Most inlets however do not exhibit such well defined characteristics and are instead the morphologic result of the relative strength of the various processes at play. The relative strength of the longshore sediment transport rates along a coastline for example plays a large role in depicting the geometry of a tidal inlet. If there is a strong net transport direction (one direction largely outweighs the other), the inlet will take a much different geometric shape than for the case of an inlet with relatively balanced longshore transport (similar transport rates in both directions). The effects of the relative strength in longshore sediment transport on inlet morphology are presented in Figure 2.16 below, after Galvin (1971).

